

PION PRODUCTION FROM FEW NUCLEON SYSTEMS  
IN THE REGION OF THE  $\Delta_{1232}$  RESONANCE.

by

KEITH MICHAEL FURUTANI

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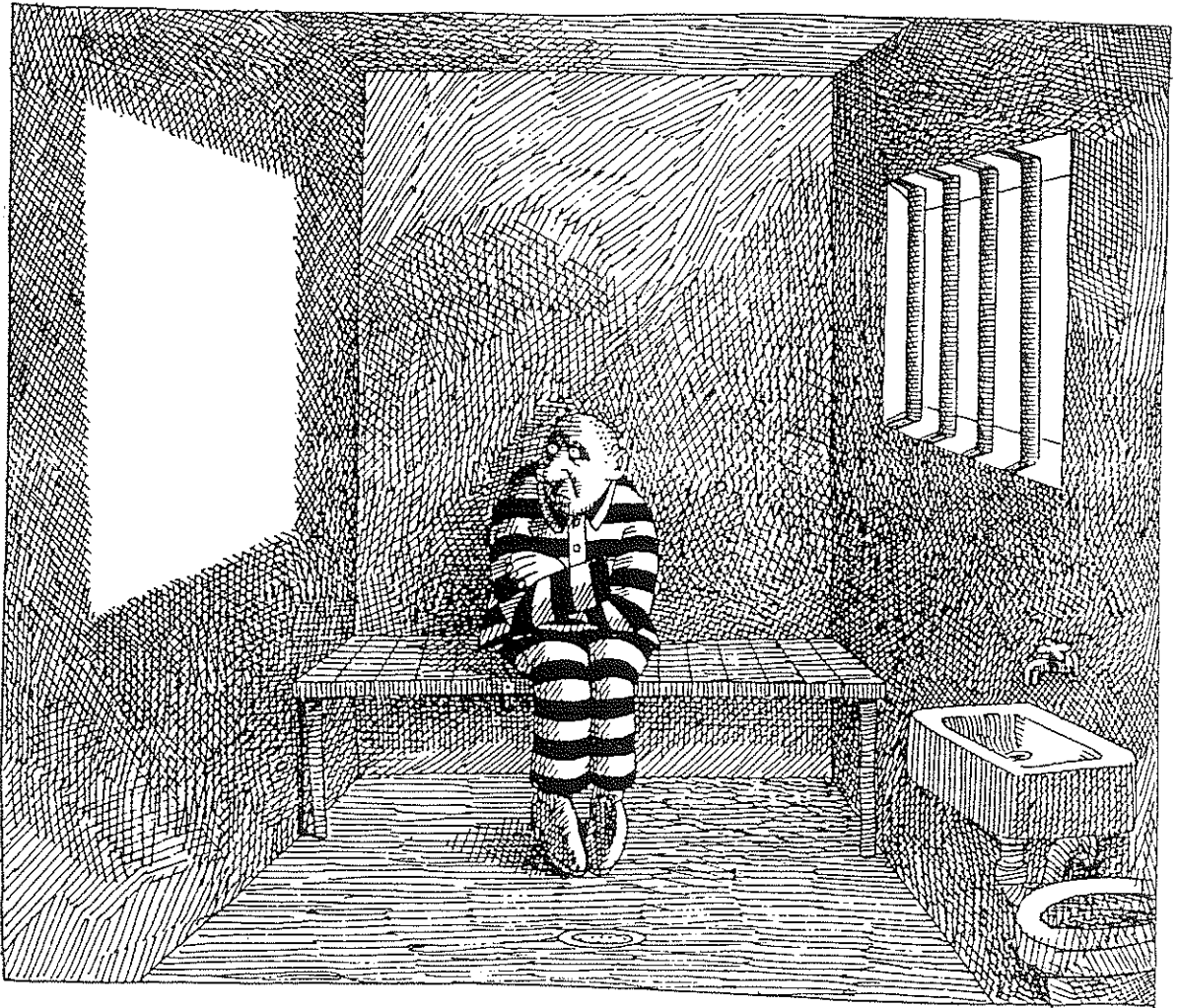
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DOCTOR OF PHILOSOPHY

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## Abstract

Angular distributions of the differential cross sections and analyzing powers have been measured for the  ${}^3\text{He}(\bar{p},\pi^+){}^4\text{He}$  and  ${}^4\text{He}(\bar{p},\pi^+){}^5\text{He}$  reactions at proton bombarding energies between 240 and 507 MeV providing a comprehensive set of data spanning the region of the  $\Delta_{1232}$  resonance.

For these experiments, the University of Manitoba/TRIUMF liquid  ${}^3\text{He}$  target was extensively refurbished and modified for provision of stable and reliable operation. A Monte Carlo study of the TRIUMF Medium Resolution Spectrometer for detecting pions was performed to investigate the acceptance, solid angle, and pion decay length of the spectrometer. Furthermore, for the  ${}^5\text{He}$  spectra, a final state interaction calculation was carried out which uses information from the  $n + \alpha$  elastic phase shifts.

The results obtained are of very high quality and are compared with a recent microscopic  $(p,\pi^+)$  calculation and a phenomenological study based upon the  $pp \rightarrow d\pi^+$  amplitudes.

## Dedication

*To my family  
for giving me the courage to strive for such a goal  
and  
to Sandra  
for the love and support to achieve it.*

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# Chapter 1

## Introduction.

It has been fifty six years since Yukawa proposed the existence of a nuclear force quantum, the pion [Yuk35], and forty four years since its discovery [Lat+47]. In the intervening period to the present day, our knowledge of the Nucleon-Nucleon (NN) interaction has evolved to a point where extremely sophisticated potentials describe a vast amount of data. The NN data are conventionally parameterized in terms of partial waves, and it is at this point that experiment is compared to theory. The NN interaction is contemporarily described by boson exchange potentials (in which the pion is just one possible boson) which quantitatively describe the NN elastic phase shifts [Mac89] very well.

While these potentials describe very well the on-shell amplitudes (elastic scattering phase shifts), it is the off-shell scattering which is the test of the validity of the potential. The success of the description of NN elastic scattering is therefore naturally extended to a description of the inelastic channel,  $NN \rightarrow NN\pi$ . These reactions are less understood and tested, primarily because now the pion is allowed to carry away some of the momentum and energy and hence one is probing a much more extensive region of phase space than in the elastic case. Furthermore, there are now three bodies in

Isospin of two nucleons		Total cross section
$T_i$	$T_f$	
1	0	$\sigma_{10}$
0	1	$\sigma_{01}$
1	1	$\sigma_{11}$
0	0	forbidden

Table 1.1: Isospin cross section notation.

the final state so that the theoretical computations are more difficult than for the two body case. The contemporary  $NN \rightarrow NN\pi$  theories are still in the spirit of the boson exchange potentials, but with diagrams which allow an inelastic pion production channel [DKS87]. Such theories are still in the early stages of testing and development.

## 1.1 Elementary Pion Production

The elementary  $NN \rightarrow NN\pi$  process is comprised of several possible reactions (for example:  $pp \rightarrow np\pi^+$ ,  $pp \rightarrow pp\pi^0$  and  $np \rightarrow pp\pi^-$ ). A simple classification results by considering the isospin basis. In the isospin basis there are only four reactions. These are denoted by the isospin of the two nucleons in the incident channel,  $T_i$ , and the outgoing channel,  $T_f$  [GW54] (see Table 1.1). Note that of the four only three are allowed, since the isospin of the pion is one.

The possible elementary pion production reactions and their relation to the isospin cross sections are shown in Table 1.2. Note that for the case of  $pp \rightarrow np\pi^+$  ( $\sigma_{10}$ ) there is the possibility of the  $np$  system forming a bound

Reaction	Cross section
$pp \rightarrow d\pi^+$	$\sigma_{10}^d$
$pp \rightarrow np\pi^+$	$\sigma_{10}^{np}$
$pp \rightarrow pp\pi^0$	$\sigma_{11}$
$np \rightarrow d\pi^0$	$\frac{1}{2}\sigma_{10}^d$
$np \rightarrow np\pi^0$	$\frac{1}{2}\sigma_{10}^{np} + \frac{1}{2}\sigma_{01}$
$np \rightarrow nn\pi^+$	$\frac{1}{2}\sigma_{11} + \frac{1}{2}\sigma_{01}$
$np \rightarrow pp\pi^-$	$\frac{1}{2}\sigma_{11} + \frac{1}{2}\sigma_{01}$

Table 1.2: Isospin cross section reactions.

state deuteron. Therefore, one usually makes the further distinction between the bound and unbound case,  $\sigma_{10}^d$  and  $\sigma_{10}^{np}$  respectively.

Using the available data the isospin cross sections may be determined and these are shown in Fig. 1.1 (from Ref. [VA82]). Note that the low energy pion production cross section is described mostly by  $\sigma_{10}^d$  (that is the  $pp \rightarrow d\pi^+$  reaction).

The description of elastic and inelastic processes from nuclei (both nucleon-nucleon scattering and pion production) is often formulated from the point of view of the elementary NN processes. However, one now must consider the field of the nucleus and this is most often described in an averaged way by an optical potential. The technique is usually performed by fitting the optical potential parameters to the case of elastic scattering at or near an equivalent energy for the process of interest. As an example, in the case of a stripping reaction,  $A(d,p)A+1$ , the analysis is usually done in a simple Born Approximation (with plane waves, a PWBA, where the field of the nucleus is ignored, or with distorted waves from an optical potential, a DWBA). A

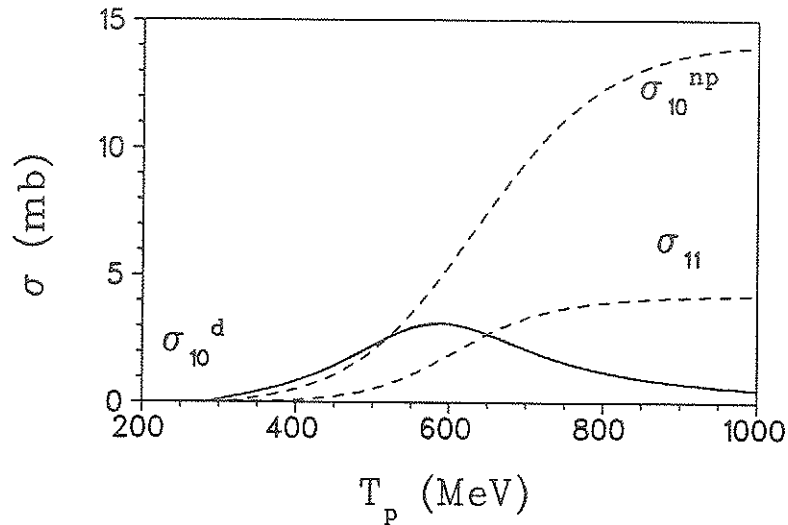


Figure 1.1: Elementary pion production cross sections.

method which is commonly used in the analysis of nucleon-nucleus elastic scattering is to include the information from the elementary processes (for example  $NN \rightarrow NN$ ) often obtained from the phase shifts. Such calculations are called Impulse Approximations (PWIA for plane waves and DWIA for distorted waves). A less common and much more difficult approach is to calculate the reaction mechanism using, for example, a boson exchange potential. The motivation for such calculations is to include the higher order processes explicitly rather than in an averaged way from an optical potential and in so doing investigate detailed points of the reaction mechanism. Such microscopic calculations are usually difficult and therefore are not as prevalent as the DWBA or DWIA calculations.

The reactions studied in this thesis are for few body nuclei ( $A=3,4$ ) of the general  $A(p, \pi^+)A+1$  reaction. These reactions may be studied theoretically in the context of a DWBA analysis (similar to those done for, say,  $(d,p)$



stripping reactions) or with a full microscopic calculation. Examining these reactions from the point of view of the fundamental  $NN \rightarrow NN\pi$  processes is also a useful approach, for example in a PWIA or DWIA calculation. These facets of  $A(p,\pi^+)A+1$  reactions and the experiments performed to investigate them will now be discussed.

## 1.2 Nuclear Pion Production

Exclusive proton induced nuclear pion production is the study of reactions like  $A(p,\pi)A+1$ , where the residual  $A+1$  nucleus is left in a well defined nuclear state. Due to the large mass difference between the incident proton and outgoing pion, there is a large momentum transfer imparted to the recoil nucleus. The momentum transfer (specifically the three momentum transfer) is defined in the following way.

$$\vec{q} = \vec{p}_f - \vec{p}_i \quad (1.1)$$

$$Q_{\text{cm}} = |\vec{q}| = \sqrt{p_i^2 + p_f^2 - 2p_i p_f \cos \theta_{\text{cm}}} \quad (1.2)$$

Even at threshold, the momentum transfer is large compared to the nuclear Fermi momentum (about 270 MeV/c). As an illustration, the momentum transfer versus the proton bombarding energy for the  ${}^3\text{He}(p,\pi^){}^4\text{He}$  and the  ${}^{40}\text{Ca}(p,\pi^){}^{41}\text{Ca}$  reactions is shown in Fig. 1.2.

It was this large momentum transfer which incited activity in the research of this reaction. It was hoped that using a simple reaction mechanism, similar to that for  $(d,p)$  stripping reactions studies, much could be learned about the high momentum components of nuclear wave functions. It was found, however, that such a simple picture did not explain the data adequately. Such ambiguities motivated detailed investigations of the reaction mechanism and also the limits of pion optical potentials. Continued uncertainty inspired

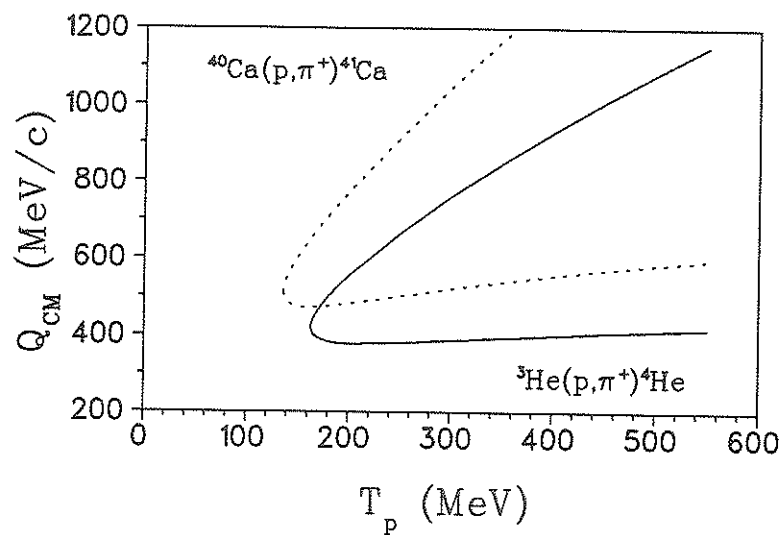


Figure 1.2: Momentum transfer for  $(p, \pi^+)$  reactions. The lower arm corresponds to  $Q_{\text{cm}}$  for  $0^\circ$  and the upper arm for  $180^\circ$ . (the solid line is for  ${}^3\text{He}(p, \pi^+){}^4\text{He}$  and the dotted line for  ${}^{40}\text{Ca}(p, \pi^+){}^{41}\text{Ca}$ ).

more effort in both theoretical and experimental investigations. While the  $A(p,\pi^+)A+1$  reaction has been a very difficult reaction to study theoretically, it has also shown to be a very sensitive proving ground for current nuclear reaction theories.

### 1.2.1 Experimental

The early work of nuclear pion production was carried out at only selected energies and angles. An example is the study of the  $^{12,13}\text{C}(p,\pi^+)^{13,14}\text{C}$  and  $^{14}\text{N}(p,\pi^+)^{15}\text{N}$  reactions at  $T_p = 600$  MeV and  $\theta = 0^\circ$  done at CERN [Dom+70] over 20 years ago. This work was followed shortly by work at Uppsala [DHG71] where both positive and negative pion production on a large variety of targets was extensively studied (see for example Refs. [Dah+73a], [Dah+73b] and [Dah+73c]). All the work was carried out near threshold (185 MeV) and with non-polarized beams but is generally considered to be the pioneering work on nuclear pion production. As an example of the quality of these early works, an energy spectrum for  $^{12}\text{C}(p,\pi^+)^{13}\text{C}$  from Ref. [DHG71] is shown in Fig. 1.3.

The earliest work of nuclear pion production involving measurements of polarization as well as differential cross section is that of Heer *et al.* [HRT58]. They investigated pion production from C and Al at 209 MeV, but at only a few angles. The polarized beam was generated in the standard manner of that time, namely via a double scattering experiment. The authors noted that such a measurement is extremely time consuming and therefore similar measurements were not performed. With the advent of the intense polarized proton beams, spin observables (notably the analyzing powers) then became additional observables of investigation. In the Ann Arbor [KS78] convention

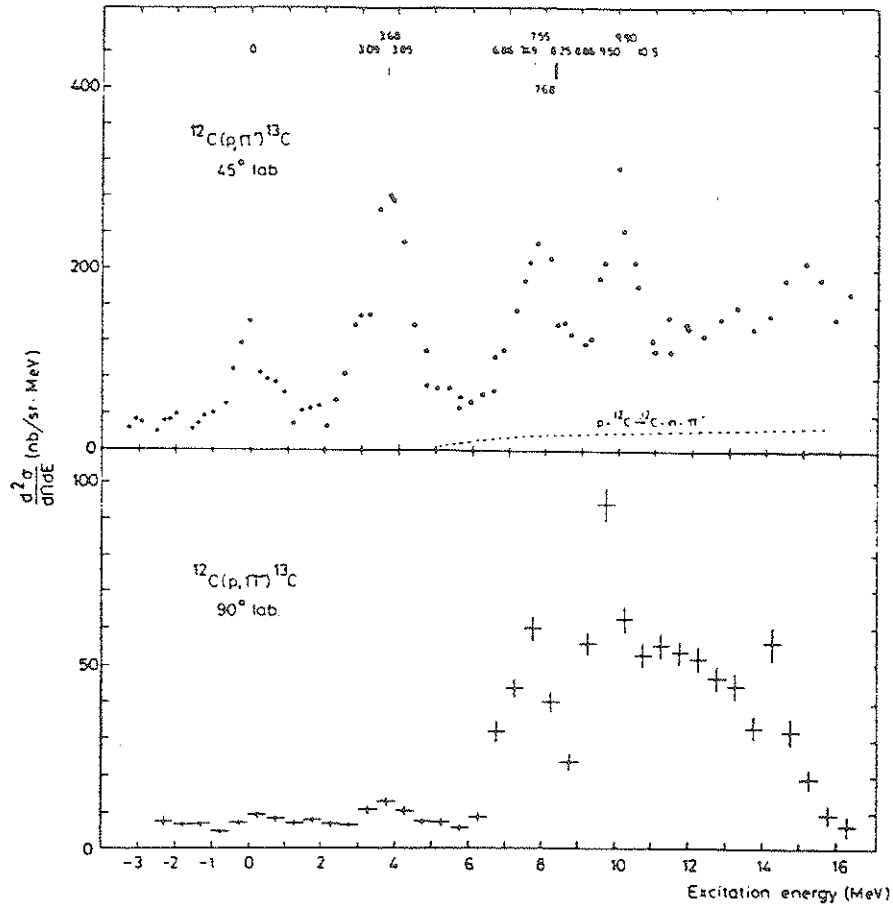


Figure 1.3: Energy spectrum from Dahlgren *et al.* showing quality of early pion production investigations.

the analyzing power is defined as

$$A_{N0} = \frac{1}{P} \frac{d\sigma(\uparrow)/d\Omega - d\sigma(\downarrow)/d\Omega}{d\sigma(\uparrow)/d\Omega + d\sigma(\downarrow)/d\Omega} \quad (1.3)$$

where  $P$  is the proton beam polarization. The subscript  $N$  indicates that the beam is polarized normal to the scattering plane and the subscript  $0$  indicates the target is unpolarized. The arrows indicate the direction of the polarization of the incident beam, i.e. up or down, with respect to the scattering plane (for example  $d\sigma(\uparrow)/d\Omega$  is then the differential cross section for polarized-up protons). In the older Madison convention [BH71] the analyzing power is notated  $A_y$ . One should note that, by definition, the range of the analyzing power is between  $+1$  and  $-1$  and that it must be zero, by symmetry, for scattering angles of  $0^\circ$  and  $180^\circ$ .

Measurements of the analyzing power for a variety of nuclei near threshold were performed at the Indiana University Cyclotron Facility (IUCF) [Sjo+81] and at TRIUMF in Vancouver, Canada [Aul+78]. The analyzing power was found to be rich in structure and has, since then proven to be the point of greatest difficulty when comparing theory with data. In some nuclei, where the differential cross section may display only little variation with energy, the analyzing power may exhibit strong variation with energy. The analyzing power thus becomes a very important observable when studying  $(p, \pi^+)$  reactions.

As an indication of the amount of interest in these reactions generated by the research of the 1970's is the number of review articles of this field of research (see Refs. [Hoi79], [MM79] and [Fea81]). These reviews are excellent surveys of the data of that period and many of the conclusions and questions raised are still applicable today. Since the  $(p, \pi^+)$  conference in Indiana in 1981 [Ben82], many experiments have been performed and are usually of very high quality stressing the importance of polarization measurements.

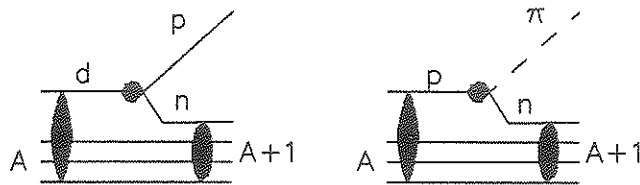


Figure 1.4: Comparison of  $(d,p)$  and  $(p,\pi^+)$  stripping reactions.

Such high quality data have inspired a great deal of theoretical activity as well, which will be presented in the following section.

### 1.2.2 Theory

In the initial investigations it was hoped that  $A(p,\pi^+)A+1$  reactions would provide information about the high momentum components of nuclear wave functions. In the context of Direct Reaction theory one can make the obvious comparison between  $(d,p)$  and  $(p,\pi^+)$  reactions, as shown in Fig 1.4. It is therefore not difficult to see why these reactions were originally studied in the so called One Nucleon Model (ONM) or pionic stripping process analogous to a DWBA formalism used for  $(d,p)$  reaction studies.

As one might expect, the calculations were found to be very sensitive to optical model wave functions. Examples of DWBA calculations can be found in the work of Tsangarides [Tsa79]. His calculations were applied to many nuclei near threshold. In light of the comparison between  $(d,p)$  and  $(p,\pi^+)$  reactions, investigations of  $A(d,p)A+1$  and  $A(p,\pi^+)A+1$  were done at similar momentum transfers for several nuclei and the DWBA analyses per-