

COMPUTATION OF BEARING CAPACITY AND PASSIVE PRESSURE  
COEFFICIENTS  
IN SAND USING STRESS-CHARACTERISTICS AND CRITICAL STATE

by

HOVAN Jean-Michel

A thesis  
presented to the University of Manitoba  
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## ABSTRACT

Bearing capacity coefficients for footings and passive pressure coefficients for walls in sand are computed using stress-characteristic solutions with the Cambridge 'Critical State' strength model. The basic equilibrium equations are combined with the Coulomb criterion to produce a set of so-called 'basic differential equations'. These equations are solved using a finite differences method and stress-characteristics are computer-drawn. The solution is rigid-plastic and does not take into account volume strains prior to failure. In the Critical State model, and hence in the analysis, the angle of shearing resistance is not considered to be constant in the sand, and is allowed to vary in the domain under stress. This takes account of the well-known curvature Coulomb-Mohr envelope, in contrast with most solutions which assume it is straight.

Bearing capacity coefficients for footings and passive pressure coefficients for walls are plotted as functions of the size of the structure (breadth of the footing, or height of the wall) and of the compressibility of the sand. These theoretical results agree well with experimental data and with previous analysis.

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# LIST OF SYMBOLS

a	-	$\sin(\psi+\mu)/((2\sigma\sin\phi \cos(\psi-\mu))$
b	-	$-\sin(\psi-\mu)/((2\sigma\sin\phi \cos(\psi+\mu))$
B	-	footing breadth
c	-	soil cohesion
Cc	-	compression index
CSL	-	critical state line
CVR	-	critical void ratio
e	-	void ratio
G <sub>s</sub>	-	specific gravity
H	-	wall height
I <sub>D</sub>	-	relative density
k	-	real variable
K <sub>p</sub>	-	passive pressure coefficient
l	-	scale parameter
m	-	number of radial lines
M	-	slope of the critical state line when it is projected on to a constant volume plane
n	-	number of spiral lines
N <sub>γ</sub>	-	bearing capacity coefficient
p	-	pressure (mean principal stress)
q	-	deviator stress
Q	-	ultimate bearing capacity
S	-	$S = \sigma_n$
T	-	$T = \tau_f$

$V$  - specific volume  
 $V_{\lambda}$  - specific volume on reference section  
 $X$  - coordinates axis  
 $Y$  - idem  
 $Z$  - idem  
 $\xi$  -  $\chi + \psi$   
 $\delta$  - contact friction angle mobilized between sand and footing base or wall surface  
 $\epsilon$  - strain  
 $\phi$  - angle of shearing resistance  
 $\gamma$  - bulk density of soil  
 $\eta$  -  $\chi - \psi$   
 $\lambda$  - slope of the critical state line  
 $\mu$  -  $\pi/4 - \phi/2$ , the angle between the slip lines and the major principal stress direction  
 $\sigma$  - normal stress  
 $\tau$  - shear stress  
 $\omega$  -  $\pi/2 - \mu$   
 $\chi$  -  $(\ln \sigma)/2 \tan \phi$   
 $\psi$  - inclination of major principal stress to the vertical axis (counterclockwise +ve)

#### Subscripts

ave - average  
 f - at failure  
 H - horizontal  
 i - initial

max - maximum  
min - minimum  
n - normal (to the plane of application)  
ps - plane strain  
tx - triaxial  
v - vertical  
1,2,3 - principal (stresses)

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## Chapter I

### INTRODUCTION

Bearing capacity factors for footings and passive pressure coefficients for walls are invariably reported in the literature as functions of the angle of shearing resistance of the soil. However, it is well acknowledged that identifying the angle of shearing resistance in a given problem is difficult. This thesis starts from the understanding that sand behaviour cannot be adequately described by a straight Coulomb-Mohr-Strength envelope as is done in traditional analysis

To get better correlation between theoretical results and experimental tests, it appears necessary to improve the methodology with respect to the sand parameters involved. One such parameter is the angle of shearing resistance that is commonly considered to be constant for reasons of simplicity. However, experimental evidence (Schofield and Wroth, 1968) shows that the Coulomb envelope flattens with increasing stress. An understanding of this evidence is mainly based on the reasonable assumption that while the intergranular friction component of the angle of shearing resistance may be essentially independent of stress level, the dilatancy component is not. That is, the angle of shearing

resistance is not a soil constant but more a variable indicating the response of the soil structure under stress.

One attempt was made by Graham and Pollock (1972) to develop a non-constant- $\phi$  solution for the failure of walls and footings. To compute the ultimate stresses under a footing base or behind a retaining wall, Graham (1974) used the so-called 'slip-line' or 'stress-characteristic' method and introduced some developments of the Sokolovski' (1960) solutions. The varying- $\phi$  relationship was composed of three distinct linear parts as a function of stress level, and had been proposed earlier by De Beer (1963). It is important to note that the work required the use of computers that fortunately became available at that time (late 1960's and early 1970's). Although the consideration of a non-constant  $\phi$  represented a more realistic approach to computer modelling of sand behaviour, some limitations could be put forward due the simplicity of the De Beer model.

With the appearance of a more reliable soil model, namely the critical state model (Roscoe, Schofield, Wroth, 1958) in conjunction with great improvement of the computer software, it was felt that the time had come to address the variable  $\phi$  question in a new 'critical state' framework. In view of these developments and in order to arrive at an adequate assessment of the influence of the angle of shearing resistance, it was thus necessary to reprogram the stress charac-

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<sup>1</sup> Also written "Sokolovskii"

teristic solution for a varying angle of shearing resistance based on the critical state model. This is the main objective of this thesis.

Firstly, in chapter 2, in order to comprehend and predict the behaviour of a sand, a conceptual framework is established : this is the critical state model. In chapter 3, a numerical procedure adapted by the author to include a new form of  $\phi$ -variation is developed and used to address the two different cases of a rough footing and of a rough passive wall. Chapter 4 describes the parametric study that has been conducted by means of a revised computer program and also the choice of the required input parameters. Results and their discussion follow in Chapter 5 and 6. Topics for further research and conclusions are listed in Chapter 7. Appendix A deals with the question of how the theoretical and parametric results could be applied in construction problems of engineering interest. Appendix B contains the computer listing and a typical output. Finally, A sample of a computer plotting program is presented in Appendix C.

## Chapter II

### SAND MODELLING

#### 2.1 INTRODUCTION

Since Coulomb introduced the simple equation  $\tau = c + \sigma \tan(\phi)$  ( $\tau$  shear stress at failure,  $\sigma$  normal stress at failure) to describe how the normal stress and the shear stress relate in a soil failure, soil mechanics has had to be developed in order to cope with the increasing complexity of industrialized society. In this process, researchers have had to approach various difficult situations such as the liquefaction of sands, slope stability problems, earthquakes, etc. These studies led them to acknowledge the complexity of the soil's response and induced them to conceive a 'unified' theory capable of handling various possible states of soil media.

#### 2.2 THE CRITICAL STATE MODEL

Roscoe, Schofield and Wroth (1958) following ideas suggested earlier by Taylor (1948) proposed that the basic parameters ' $\tau, \sigma, e$ ' were necessary and sufficient to describe the behaviour of a soil element. They conducted a series of drained and undrained tests on sand to obtain a unique yield surface composed of two major subsurfaces (Fig 2.1):

1. The Hvorslev surface
2. The Roscoe surface

The link between the two subsurfaces is called the critical void ratio line (CVR) or critical state line (CSL) and represents the ultimate state of a soil sample. Once at this 'critical state', the soil can be sheared without any further change of the three basic parameters.

Technically in material science we are dealing with yield surfaces because plastic-hardening is involved. However, in most soil mechanics work using the Coulomb-Mohr criterion, this should be called a 'rupture' or 'failure' envelope.

Fig 2.1 represents a schematic diagram of the model. This three dimensional plot can be expressed in two ways :

1. a  $(\tau, \sigma, e)$  diagram
2. a  $(q, p', V)$  diagram
  - a) with  $p' = (\sigma_1' + \sigma_2' + \sigma_3')/3$  , the mean principal stress.
  - b)  $q = \sigma_1' - \sigma_3'$  , the deviator stress
  - c)  $V = 1 + e$  , the specific volume

Although plot(1) is easier to visualize, plot(2) has the merit of taking into account the intermediate principal stress. As acknowledged by the position and the shape of the yield surface the shearing resistance of a soil depends not only on the normal stress but also on the value of the

initial void ratio. However, it is important to note that the ultimate state represented by the CVR line is associated with independence from stress history effects.

In a  $\ln(p'):V$  space, the normal consolidation line (NCL) and the critical state line (CSL) are experimentally found to be parallel (Wroth and Basset, 1965, Fig 2.2). For various reasons associated with the low compressibility  $\lambda$  of sands<sup>2</sup> it has proved helpful to modify the original presentation of the critical state model by projecting all the points onto a reference section set arbitrarily as  $p'=1$ . This has permitted the three dimensional model to be represented in a two dimensional space, namely  $q/p' : V_\lambda$  (Atkinson and Bransby, 1978, Fig 2.3), where

1.  $V_\lambda = V + \lambda \text{Log}(p')$ , and
2.  $\lambda$  is the slope of the CSL.

The present study has used these normalized plots as a representation of the critical state model for sand. This has enabled the author to develop a simple iteration process which can be used in calculating failure loads of structures in sand (section 3.6)

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<sup>2</sup> The term 'compressibility' is retained for  $\lambda$  ( $\lambda=C_c/2.3$ )

### 2.3 MODELLING

In 1937, Casagrande stated simply that a dense sand exhibits a high friction angle whereas a loose sand tends to exhibit a lower friction angle (Rowe, 1967). In the context of the sophistication and the simplicity of the critical state model and since the angle of shearing resistance is the characteristic that determines the response of the soil structure to external stresses, it has become necessary to understand the dependency of the angle of shearing resistance with respect to the 'basic parameters'  $e$  and  $p'$  (or  $e$  and  $\sigma$ ).

Rowe (1967) separated the angle of shearing resistance  $\phi$  into three components :

1. strength developed by frictional resistance and depending on the mineralogical content of the sand.
2. strength developed by the energy required to cause expansion or dilation of the material.
3. strength developed by energy required to rearrange and reorient soils particles.

Fig 2.4 shows the distribution of each component as a function of the relative density  $I_D$ . As might be expected, the dilatancy effect plays an increasing role as the relative density increases. This explains the higher friction angle of denser sands.

Although Rowe's diagram (Fig 2.2) is of considerable interest in understanding the behaviour of  $\phi$  as a function of

the void ratio (represented by  $I_D$ ) it does not address the influence of the confining pressure, that is, the second basic parameter of the critical state model.

Studying the behaviour of sands at extremely low pressures, Ponce and Bell (1971) found themselves in the position where they had to take into account the influence of the confining pressure and proposed a three-dimensional diagram  $\phi = f(e, P)$ ; (Fig 2.5).

With regards to this representation, it is important to note that at extremely low pressures, shear produces almost no rearranging of particles since most of the energy is dissipated in expansion and true sliding; in other terms, it can be stated that low confining pressures cannot sufficiently restrain the sand to permit an internal reorganisation. Another consideration is the rapid decrease of  $\phi$  with increasing pressure and the greater influence of the rearranging component over the dilatancy one for loose sands.

Even if the detailed study of each component of the angle of shearing resistance does not come directly into use in geotechnical calculations, that is, only the overall knowledge of  $\phi$  is used, the parallel between the critical state model ( $e, \sigma, \tau$ ) in Fig 2.1a, and the three-dimensional  $\phi$  diagram ( $e, \sigma, \phi$ ) in Fig 2.5 permits the establishment of a sound 'foundation' to explain the systematic variation of the angle of shearing resistance. More particularly should



be noticed the conceptual parallels between the  $\tau$ -axis which represents the real physical shearing resistance of soil in engineering applications (macro-behaviour) and the  $\phi$ -axis representing a measure of strength arising from engineering analysis (micro-behaviour).

#### 2.4 CHOICE OF A $\phi_{ps} \equiv f(\phi_{tx})$ RELATIONSHIP

This thesis addresses the computation of bearing capacity coefficients for strip footings and passive pressure coefficients for retaining walls. Both cases are two-dimensional and require the knowledge of a plane strain friction angle. Since triaxial tests are easiest to perform in the laboratory, they are used to establish the critical state model and they lead to angles of shearing resistance for three dimensional (axisymmetric) stress states. It was therefore necessary to find a relationship that would transform triaxial  $\phi$ -values ( $\phi_{tx}$ ) to  $\phi$ -plane strain values ( $\phi_{ps}$ ).

Various researchers have proposed such relationships:

1. Wroth (1984,pers. comm.)  $\phi_{ps} = 9/8 \phi_{tx}$
2. Lade (1976)
  - a)  $\phi_{ps} = 1.5 \phi_{tx} - 17^\circ$  for  $\phi \geq 34^\circ$ .
  - b)  $\phi_{ps} = \phi_{tx}$   $\phi_{tx} \leq 34^\circ$
3. Bishop (1966) : Fig 2.6 in the Sixth Rankine Lecture, 1966 (after Cornforth, 1964).

These various proposals are presented in Fig 2.7.

In this thesis, relationship (3) by Bishop has been used in the revised form shown in Fig 2.8. For use in the computer, the relationship was plotted in logarithmic terms and was subdivided into three linear parts (Fig 2.8) :

1.  $\phi_{tx} < 33^\circ$        $\phi_{ps} = \phi_{tx}$
2.  $33^\circ \leq \phi_{tx} < 36^\circ$        $\ln \phi_{ps} = 1.666 \ln \phi_{tx} - 2.336$
3.  $\phi_{tx} \geq 36^\circ$        $\ln \phi_{ps} = 1.293 \ln \phi_{tx} - 1.002$

## Chapter III

### STRESS-CHARACTERISTIC METHOD AND METHODOLOGY

#### 3.1 INTRODUCTION

At the state of plastic limit equilibrium, the soil beneath a footing or behind a passive retaining wall is stressed to its limiting or yield condition. In loose and medium dense sands, failure is close to the critical state, with :  $\frac{\partial p}{\partial \epsilon} = \frac{\partial q}{\partial \epsilon} = \frac{\partial v}{\partial \epsilon} = 0$ . In dense sands, further straining is required after failure before critical state is reached. Sokolov (1960) developed the stress-characteristic method to compute the stresses at failure. Graham (1968) pursued this idea and improved the numerical accuracy of the solution.

#### 3.2 CHOICE OF A YIELD CRITERION

Coulomb (1776) suggested the equation :

$$\tau_f = c + \sigma_n \tan(\phi)$$

This equation states that plastic flow will occur (plastic failure) when the shear stress  $\tau$  reaches a value that depends on two material parameters and on the stress level. These are the cohesion of the soil  $c$ , the angle of shear resistance  $\phi$  and the normal stress to the plane of application  $\sigma_n$ .

Although this equation does not take into account the influence of the intermediate stress, Fig 3.1 provides some experimental evidence that the use of the Coulomb criterion is more accurate than the corresponding Tresca or the Von Mises criteria. The Coulomb criterion has been adopted in this thesis to define a 'yield' criterion.

### 3.3 SOKOLOVSKI METHOD

The Coulomb criterion is combined with the static equilibrium equations to provide a set of so-called 'basic equations' (Sokolovski, 1960). One method of solving these equations consists of using logarithmic transformations to generate curvilinear coordinates whose directions coincide with the direction of the failure planes. This is commonly known as the method of stress characteristics. In turn, the solved system in logarithmic stress space provides a set of slip lines or a slip-line field whose positions are known in the physical (x,z) plane.

Assuming that the soil could be modelled by a constant value of the angle of shearing resistance, Sokolovski (1960) developed a numerical procedure based on approximations of the hyperbolic differential equations to obtain the slip-line field for cohesionless as well as cohesive soils. Since this thesis aims at analysing localized  $\phi$ -variations in a granular (cohesionless) material, a more general application of Sokolovski's procedure has had to be developed in this thesis project.

### 3.4 LOCAL $\phi$ -VARIATION ASSOCIATED WITH SOKOLOVSKI'S METHOD

#### 3.4.1 Local system of references

The proposed model is based on the representation of Fig 3.2. Locally, the curvature of the Coulomb-Mohr envelope can be treated as presenting a 'cohesive' component in a general system of references. This entails using the general differential equations developed for a  $(c, \phi)$  medium (Sokolovski, 1960). This is justified by the more general aspect of a cohesive frictional medium over a non-cohesive one. That is, a cohesionless soil is a more restricted case where the envelope has no intercept with the  $r$ -axis.

#### 3.4.2 General system of references

a) In a two-dimensional representation (Fig 3.2), the equilibrium equations are :

$$1. \quad \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} = X$$

$$2. \quad \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} = Y \quad \text{Eqs. 3.1}$$

b) The Coulomb-Mohr criterion as a yield condition is described by (Fig 3.3) :

$$1. \quad \max(|\tau_n| - f(\sigma_n)) = 0$$

2. or  $T = f(S)$  where

a)  $T = \tau_n$

b)  $S = \sigma_n$

This yield condition can be written as (Fig 3.3) :

$$\frac{d}{d\omega} (|\tau_n| - f(\sigma_n)) = 0 \quad \text{Eq. 3.2}$$

where

1.  $\sigma_n = 1/2(\sigma_1 + \sigma_3) + 1/2(\sigma_1 - \sigma_3)\cos 2\omega$

2.  $\tau_n = 1/2(\sigma_1 - \sigma_3)\sin 2\omega$

Substituting  $\sigma_n$  and  $\tau_n$  into Eq. 3.2 :

$$\frac{d}{d\omega} (|\tau_n| - f(\sigma_n)) = (\sigma_1 - \sigma_3)\sin 2\omega (\cos 2\omega + f'(\sigma_n))$$

and therefore  $\frac{d}{d\omega} (|\tau_n| - f(\sigma_n)) = 0$  provides the general condition  $\cot 2\mu = f'(\sigma_n)$  where  $2\mu = \pi - 2\omega$  (Fig 3.4). This angle determines the position of the slip lines inclined to the principal axis at the angle  $\pm \mu$ . At failure the stress conditions are :

1.  $\sigma_n = 1/2(\sigma_1 + \sigma_3) - 1/2(\sigma_1 - \sigma_3)\cos 2\mu$

2.  $\tau_n = 1/2(\sigma_1 - \sigma_3)\sin 2\mu$

In a chosen system of references (Fig 3.5), for any point A

1.  $\begin{matrix} \sigma_x \\ \sigma_y \end{matrix} = 1/2(\sigma_1 + \sigma_3) \pm 1/2(\sigma_1 - \sigma_3)\cos 2\mu$

2.  $\tau_{xy} = 1/2(\sigma_1 - \sigma_3)\sin 2\mu$

substituting  $\sigma_n$  and  $\tau_n$  into  $\sigma_x, \sigma_y, \tau_{xy}$  we get :

$$1. \quad \begin{matrix} \sigma_x \\ \sigma_y \end{matrix} = \sigma_n + \tau_n \frac{\cos 2\mu \pm \cos 2\psi}{\sin 2\mu}$$

$$2. \quad \tau_{xy} = \tau \sin 2\psi / \sin 2\mu$$

Substituting  $\sigma_x$ ,  $\sigma_y$ ,  $\tau_{xy}$  into the equilibrium equations 3.1 produces a new set of equilibrium equations 3.3 :

$$1. \quad (1 + \cos 2\mu \cos 2\psi) \frac{\partial \chi}{\partial x} + \cos 2\mu \sin 2\psi \frac{\partial \chi}{\partial y} - \sin 2\mu (\sin 2\psi \frac{\partial \psi}{\partial x} - \cos 2\psi \frac{\partial \psi}{\partial y}) = \frac{\sin^2 2\psi}{2\tau_n}$$

$$2. \quad \cos 2\mu \sin 2\psi \frac{\partial \chi}{\partial x} + (1 - \cos 2\mu \cos 2\psi) \frac{\partial \chi}{\partial y} + \sin 2\mu (\cos 2\psi \frac{\partial \psi}{\partial x} + \sin 2\psi \frac{\partial \psi}{\partial y}) = \frac{\sin^2 2\mu}{2\tau_n}$$

where the Mandel function  $\chi$  is defined by :

$$2\chi = \frac{d\sigma_n}{\tau_n} - 2d\psi$$

It is now convenient to introduce new variables defined by :

$$1. \quad \eta = \chi(\mu) - \psi$$

$$2. \quad \xi = \chi(\mu) + \psi$$

Using these variables, Eqs. 3.3 become :

$$1. \quad \frac{\partial \eta}{\partial x} + \tan(\psi - \mu) \frac{\partial \eta}{\partial y} = a$$

$$2. \quad \frac{\partial \xi}{\partial x} + \tan(\psi + \mu) \frac{\partial \xi}{\partial y} = b \quad \text{Eqs. 3.4}$$

where

$$\frac{a}{b} = \pm \frac{\sin 2\mu (X \sin(\psi \pm \mu) - Y \cos(\psi \pm \mu))}{2\tau_n \cos(\psi \pm \mu)}$$

If the present problem is redefined in dimensionless terms with gravitational acceleration in the X-direction, the body forces are (Fig 3.5) :

$$x = 1 \quad \text{and} \quad y = 0$$

The right-hand part of Eqs. 3.4 becomes :

$$\frac{a}{b} = \pm \frac{\sin 2\mu \sin(\psi \pm \mu)}{2\tau_n \cos(\psi \pm \mu)}$$

Since at failure  $\tau_n = 1/2(\sigma_1 - \sigma_3)\sin 2\mu$  (Fig 3.4), we obtain :

$$\frac{a}{b} = \pm \frac{\sin(\psi \pm \mu)}{(\sigma_1 - \sigma_3) \cos(\psi \pm \mu)}$$

For cohesionless materials, the angle of shearing resistance  $\phi$  is defined as the angle between the tangent to the Coulomb-Mohr circle passing through the origin and the  $\sigma$ -axis. In principle, in the general case, the locus of the intersection of the tangent and of the Coulomb-Mohr stress circle could describe any curve (Fig 3.6). In order to match the definition of  $\phi$  with the theoretical general case, an assumption is made in the context of the finite differences method :

Locally ,  $\phi$  can be determined as :  $\sin \phi = (\sigma_1 - \sigma_2)/2\sigma$



where  $\sigma = (\sigma_1 + \sigma_3)/2$

In view of this assumption,  $a$  and  $b$  in Eqs. 3.4 can be rewritten as :

$$\frac{a}{b} = \pm \frac{\sin(\psi \pm \mu)}{2\sigma \sin\phi \cos(\psi \pm \mu)}$$

The system of equations to be solved becomes, in the adopted system of coordinates (Fig 3.7) where now the  $z$ -axis is in the gravity direction, and the  $x$ -axis is horizontal :

$$\begin{aligned} 1. \quad \frac{\partial \xi}{\partial Z} + \tan(\psi + \mu) \frac{\partial \eta}{\partial X} &= b = \frac{-\sin(\psi - \mu)}{2\sigma \sin\phi \cos(\psi + \mu)} \\ 2. \quad \frac{\partial \eta}{\partial Z} + \tan(\psi - \mu) \frac{\partial \xi}{\partial X} &= a = \frac{\sin(\psi + \mu)}{2 \sin \cos(\psi - \mu)} \quad \text{Eqs. 3.5} \end{aligned}$$

$$, \text{ as before } \quad \eta = x - \psi \quad \xi = x + \psi$$

$$\text{and } 2d\chi = d\sigma_n / \tau_n - 2d\mu$$

Now  $d\psi$  can be integrated after substitution of  $\sigma_n$

and  $\tau_n$ , bearing in mind that  $\sigma_n = \sigma - \tau_n$

$$2d\chi = \frac{2\sin^2 2\mu + 2\tau_n d\mu}{\tau_n} - 2d\mu$$

$\tau$  can be found from :  $\sin\phi = \tau / \sigma \sin 2\mu$ .

$$\text{Since } 2\mu = \pi/2 - \phi ,$$

$$\text{we have } \sin\phi = \tau / \sigma \cos\phi ,$$

$$\text{and thus } \tau = \sigma \sin\phi \cos\phi .$$

$$\text{Finally, } 2d\chi = d\sigma / \sigma \tan\phi$$

With the assumption that  $\phi$  is non-variable within a limited finite range of stresses, the last equation can be integrated to give :

$$x = \log \sigma / 2 \tan \phi$$

The set of differential equations 3.5 can now be easily solved by the method of the characteristics.

Use the differential relations :

$$d\xi = \frac{\partial \xi}{\partial Z} dZ + \frac{\partial \xi}{\partial X} dX \quad ; \quad d\eta = \frac{\partial \eta}{\partial Z} dZ + \frac{\partial \eta}{\partial X} dX$$

Solving Eq. 3.5 with these relations we obtain :

$$1. \quad \frac{\partial \xi}{\partial Z} = \frac{bdX - \tan(\psi + \mu)d\xi}{dX - \tan(\psi + \xi)dZ}$$

$$2. \quad \frac{\partial \eta}{\partial Z} = \frac{adX - \tan(\psi - \mu)d\eta}{dX - \tan(\psi - \mu)dZ}$$

Eqs. 3.6

The stress characteristics correspond to cases where the numerator and the denominator are simultaneously equal to zero. This determines two real separate families of stress characteristics :

$$1. \quad \frac{d\eta}{dZ} = a \quad \text{along} \quad \frac{dX}{dZ} = \tan(\psi - \mu)$$

$$2. \quad \frac{d\xi}{dZ} = b \quad \text{along} \quad \frac{dX}{dZ} = \tan(\psi + \mu)$$

Eqs. 3.7

Note that the slopes  $(\psi \pm \mu)$  of the characteristics are in the direction of limiting shear stress, and are therefore the local inclinations of failure planes.

Rewriting the above set in finite-difference terms

$$1. (X-X_1) = (Z-Z_1)\tan(\psi_1 - \mu) \quad ; \quad \eta - \eta_1 = (Z-Z_1)a_1$$

$$2. (X-X_2) = (Z-Z_2)\tan(\psi_2 + \mu) \quad ; \quad \xi - \xi_2 = (Z-Z_2)b_2$$

Eqs. 3.8

The unknowns are the variables  $X, Z, \eta, \xi$  which can be evaluated from the previously known variables  $X_1, X_2, Z_1, Z_2, \eta_1, \xi_2$ . The solution will then proceed from two previously determined points say  $P_1, P_2$  to a new point  $P$  (Fig 3.8).

Solving Eqs. 3.8 directly for  $X, Z, \eta, \xi$ .

$$Z = \frac{Z_1 \tan(\psi_1 - \mu) - X_1 - Z_2 \tan(\psi_2 + \mu) + X_2}{\tan(\psi_1 - \mu) - \tan(\psi_2 + \mu)}$$

$$X = X_1 + (Z-Z_1)\tan(\psi_1 - \mu) = X_2 + (Z-Z_2)\tan(\psi_2 + \mu)$$

$$\eta = \eta_1 + (Z-Z_1)a_1 \quad ; \quad \xi = \xi_2 + (Z-Z_2)b_2$$

On reversing the logarithmic transformation :

$$\sigma = \exp((\xi + \eta)\tan\phi) \quad \psi = 1/2(\xi - \eta)$$

This simple finite-difference method yields only approximate results since no account of the curvature of the slip lines is considered in the use of  $\psi_1$  and  $\psi_2$ . Previous researchers have used iteration pro-

cesses on  $\psi$  (De Jong, 1959, Sokolovski, 1960), and on both  $\sigma$  and  $\psi$  (Graham, 1968) to take account of the curvature of generalized slip-line fields.

The computer program used in this thesis added one additional iteration process.

1. iteration on  $\sigma$  (Sokolovski, 1960; De Jong, 1959)
2. iteration on  $\psi$  (Graham, 1968)
3. iteration on  $\phi$  (Author)

This is a new development of earlier procedures suggested by Graham and Pollock (1972). The additional iteration is used to model the dependency of  $\phi$  on pressure, that is, the curvature of the Coulomb-Mohr envelope.

### 3.5 THE $\phi$ -ITERATION PROCESS

The process has been developed specifically to improve computational accuracy in a domain where  $\phi$  varies with stress level. That is, the  $\phi$ -value used in any small section of the solution arises from the solution itself, and is not predetermined by the analyst. It is based on the normalized plot of the critical state model outlined in chapter 2. Fig 2.2 represented this model on a plot of  $q/p'$  vs  $V+\lambda \ln(p')$ . For every state of stress and every specific volume there corresponds a value  $q/p'$  on the overconsolidated Hvorslev surface represented by the straight line in the normalized plot.

Since  $q/p' = 6 \sin\phi / (3 - \sin\phi)$ ,

$$\phi_{tx} = \text{Arcsin}(3M/(6 + M)) \text{ where } M = q/p'$$

In the  $\phi$ -iteration procedure, this value of  $\phi_{tx}$  is compared with the initial value of  $\phi$  previously used to calculate  $p'$ . If  $|\phi_{tx} - \phi_i| / \phi_i > 10^{-3}$ , the initial value of  $\phi$  is reset as  $\phi_i = \phi_{tx}$ . The process is repeated until the convergence criterion set as  $10^{-3}$  in the program is satisfied.

### 3.6 DETERMINATION OF THE SLIP-LINE FIELD

Proceeding from a known boundary, that in this work is commonly taken as the edge of a passive zone beneath the free surface where  $\psi = \pi/2$ , the program computes point by point the slip-line field towards an end boundary where failure stresses are to be evaluated. The type of problem being examined (that is wall or footing, for  $K_p$  or  $N\gamma$  values respectively) determines the data input to the program.

If point O is taken as the point of origin of the adopted system of physical coordinates, set as either the bottom corner of a footing or the top of a passive wall, the slip-line field consists of two families of characteristics :

1. a set of curved radial lines originating from point O.
2. a set of spiral lines intersecting the radial lines in turn at an angle of  $2\mu = \pi/2 - \phi$

Since point O is a singular point of the slip-line field, it represents a point of discontinuity in the mathematical solution of the basic equations. At stress equals zero, the logarithmic transformations that are involved in the solutions tend to infinity. In order to handle this problem, Graham (1968) introduced a surcharge term in the computation of the slip-line field that allows the logarithmic stress range to remain finite. The effect of the surcharge term is then reduced by 'shrinking' the field by a factor of  $10^n$ , n being a number of scale reductions necessary to almost eliminate the effect of the surcharge. Graham (1968) showed that the surcharge component is negligible after four or five scale reductions depending on the required accuracy. It is important to note that this scale reduction process does not introduce any scale effects itself since all computations are carried out in dimensionless terms. The influence of footing size on the results is handled separately in chapter 6 (Discussion of Results).

The present study has retained this procedure for handling the singularity at point O.

### 3.7 EDGE OF THE PASSIVE ZONE :THE INITIAL BOUNDARY CONDITION

This boundary is determined prior to any computation of the characteristic field and is defined as :

x : horizontal distance from O

$z = x \tan (\pi/4 - \phi/2)$  depth of the soil.

$\sigma = (1.0 + z) / (1 - \sin\phi)$  where 1.0 is the sur-charge term

$\psi = \pi/2$  direction of the major principal stress

Since the solution addresses a non-constant  $\phi$  problem, the edge of the passive zone has to be developed as a boundary using the  $\phi$  iteration process to compute a coupled set of compatible  $\phi$ - and  $\sigma$ - values. As a result, the shape of this boundary is not a straight line but exhibits a slight convex curvature with increasing depth because stress levels (and therefore  $\phi$ -values) change with depth.

### 3.8 THE END BOUNDARY : SURFACE WHERE LOADINGS ARE NEEDED

Two distinct cases can be identified.

#### 3.8.1 Passive wall

The solution addresses the case of a fully rough wall where the angle of wall friction which is mobilized corresponds to the local angle of shearing resistance of the sand (Graham, 1971; Shields and Tolunay, 1972)). Since no discontinuities exist at the interaction between the sand and the wall, the wall is in this case a slip surface where normal stresses computed by the solution can be integrated as a force boundary. The computer program includes a subroutine that performs this operation and integrates the pressures into the dimensionless coefficient  $K = ( \int_0^H \sigma_H dx ) / 0.5 \gamma H^2$  where  $H$  is the wall height, and  $\sigma_H$  is horizontal pressure.

In this case,  $K_p$  is a function of the height of the wall, even though it is a dimensionless coefficient. That is, the program permits the examination of scale effect.

### 3.8.2 Bearing capacity for footings

Differing from the passive wall failure, a foundation failure presents two symmetrical zones of failure where the soil flows outward from the centre-line.

Graham (1971) showed that a solution that involves the extension of the slip lines up to the base of the footing should also address the problem of the discontinuity of the base friction angle  $\delta$  at the centre-line of the footing base. He proposed a linear- $\delta$  variation to handle this situation where the footing base is considered to be a slip surface.

An alternative considers that an elastic wedge is trapped below the footing preventing the extension of the slip lines. This concept seems now acknowledged by various researchers. Vesic (1973) concluded that the stress and the deformation patterns under a compressed area in an actual footing is such that it always leads to the formation of a wedge.

The present solution is based on the assumption of an elastic trapped wedge. In this case, the end boundary in the stress field computation is the lower edge of the elas-



tic wedge and is also a slip-line. This boundary is inclined to an angle  $\phi$  with the footing base. Since the angle of shearing resistance has been allowed to vary in this thesis, the edge of the elastic zone must be determined using the  $\phi$  iteration procedure. This leads to a slightly curved edge which is concave upwards.

Vertical stresses on the two symmetrical lower boundaries of the wedge are calculated for different-sized footings and then expressed as the dimensionless (but scale dependent) parameter  $N\gamma$ , where for surface footings :

$$Q = 0.5 \gamma B^2 N\gamma$$

### 3.9 THE RIGID-PLASTIC ASSUMPTION

The solution which has been developed above is based on :

1. a yield criterion : the Coulomb-Mohr criterion incorporated into a critical state model,
2. static equilibrium equations, expressing rigid-plastic behaviour (Fig 3.9).

The solution does not take into account any strain-stress relationship in the sand and assumes negligible volume strains prior to failure. Clearly, the validity of such an assumption must be associated with the mode of failure observed. Even if the soil density alone does not determine the mode of failure patterns, Fig 3.10 shows that a loose sand (compressible) tends to fail in a 'punching shear fail-

ure' mode whereas a dense sand (incompressible) tends to fail in a 'local' or 'general shear failure' mode. These various failure modes were reported by Vesic (1963) and De Beer (1970).

It therefore seems that the idea of associating a relatively incompressible material such as a medium dense sand with a rigid-plastic behaviour does not alter the basic validity of the theoretical stress-characteristic solution. In the absence of an exact solution, this assumption is retained, especially for the range of densities considered (medium dense sand). By careful consideration of boundary conditions that changed with displacement of the structure, Graham (1974) was able to compute approximate load-displacement interactions.

## Chapter IV

### PARAMETRIC STUDY

#### 4.1 INTRODUCTION

In order to address a great number of possible variables , a broad range of input parameters has been tested.

#### 4.2 DEFINITION OF THE CRITICAL STATE MODEL

The use of normalized plots of the sand behaviour (Fig 4.1) requires the input of a series of parameters necessary to define the  $q/p'$  and the  $V_\lambda$  axes.

##### 4.2.1 $V_\lambda$ axis

$V_\lambda$  is defined by  $V_\lambda = V + \lambda \ln(p')$  where  $V$  is the specific volume ( $V=1+e$ ) and  $\lambda$  is the slope of the critical state line.  $V_\lambda$  has been allowed to vary within a range of values referred to as  $(V_\lambda)_{\min}$  and  $(V_\lambda)_{\max}$  where :

$$(V_\lambda)_{\min} = (1+e_{\min}) + \lambda \ln(p'_{\min})$$

$$(V_\lambda)_{\max} = (1+e_{\max}) + \lambda \ln(p'_{\max})$$

Two ranges of values were tested. They are hereafter referred to the 'limited' range case, and the 'broad' range case. For the limited range , the necessary values were fixed to be :

$$e_{\min} = 0.5, \quad p'_{\min} = 10 \text{ kPa}$$

$$e_{\max} = 0.7, \quad p'_{\max} = 10^4 \text{ kPa}$$

For the broad range case, these values were :

$$e_{\min} = 0.4, \quad p'_{\min} = 10 \text{ kPa}$$

$$e_{\max} = 0.8, \quad p'_{\max} = 10^4 \text{ kPa}$$

It is important to note that the values of the void ratio and pressure were chosen to be representative of extreme values that might be encountered in real sands, and to permit a parametric study. Bearing in mind that this study is directed particularly at medium dense sands (section 3.9), values of the void ratio should be associated with a 'medium' range. This question is addressed in Appendix A. The literature records few cases where the compressibility  $\lambda$  (or  $C_c$ ) has been measured. Mostly, the behaviour of sands is related to its placement void ratio and unit weight. Values which have been found include :

Ref.	Sand type	$\lambda$
Lee and Farhoomand, 1967	subrounded	0.150
Lee and Farfoomand, 1967	angular	0.250
Atkinson and Bransby, 1978	"typical"	0.100
Vesic and Clough, 1968	Chatahoochee river	0.175

The work in this thesis has been performed on the basis of  $\lambda$ -values chosen to cover the range described by this limited number of references. In order to study the influence of the compressibility  $\lambda$ , a range of values of  $\lambda$  was tested for both the limited and the broad range cases. These values are :

$$\lambda = 0.05, 0.10, 0.175, 0.250$$

#### 4.2.2 q/p'-axis

The limits of the chosen range were fixed as  $(q/p')_{\min}$  and  $(q/p')_{\max}$  corresponding respectively to  $(v_{\lambda})_{\max}$  and  $(v_{\lambda})_{\min}$ . A minimum and a maximum angle of shearing resistance were given as input parameters and the use of the relationship  $q/p' = 6\sin(\phi)/(3-\sin\phi)$  provided the two limits  $(q/p')_{\min}$  and  $(q/p')_{\max}$ . For the limited range case, the values of the limits of the angle of shearing resistance were chosen as :

$$\begin{aligned}\phi'_{\min} &= 32^{\circ} \\ \phi'_{\max} &= 40^{\circ}\end{aligned}$$

For the broad range case, these values were :

$$\begin{aligned}\phi'_{\min} &= 28^{\circ} \\ \phi'_{\max} &= 45^{\circ}\end{aligned}$$

#### 4.3 OTHER INPUT PARAMETERS

##### 4.3.1 Density $\gamma$

An average dry density was considered for each case and was defined by :

$$\gamma_{ave} = \frac{G_s}{1 + e_{ave}} \gamma_w$$

where  $e_{ave}$  =

$$G_s = 2.65 \text{ specific gravity}$$

$$\gamma_w = 9.81 \text{ kN/m}^3 \text{ water unit density}$$

##### 4.3.2 Scale parameter : $l$

All the variables throughout the computation of the stress-characteristics were expressed in dimensionless terms. That is  $\sigma = \sigma_r / \gamma l$ ,  $X = X_r / l$ ,  $Z = Z_r / l$  where  $l$  is a scale parameter used to convert real physical plane dimensional parameters  $\sigma$ ,  $X$ ,  $Z$  into dimensionless ones, and vice-versa. The scale parameter  $l$  was chosen as the horizontal length of the edge of the passive zone (Fig 4.1 and 4.2).

##### 4.3.3 Number of spiral and radial lines

In order to compare results from each case studied, the numbers of spiral and of radial lines used in the computation were kept identical throughout the parametric study. Ten spiral lines and twenty radial lines were considered. It is important to note that the radial lines are not straight lines but curved under gravity forces. This in effect re-

quires the computation of extra radial lines after each scale reduction of the domain (section 3.5). To input a higher number of spiral lines and of radial lines would slightly improve the accuracy of the solution. However, it would also lengthen the computation time. It was felt that a compromise has been reached with the chosen figures (10 spirals, 20 radials). This question was explored by Graham (1968).

#### 4.3.4 Edge of the passive zone

The number of spirals having being chosen, the abscissas of the intersection of the spirals and of the edge the passive zone are input in dimensionless terms. Corresponding  $z$ ,  $\sigma$ ,  $\psi$ -values are then computed by the computer program (section 3.7)

#### 4.3.5 Summary of all input parameters

For each case computed, the input parameters are:

1.  $e_{\min}, e_{\max}$  : minimum, maximum void ratio
2.  $p'_{\min}, p'_{\max}$  : minimum, maximum pressure
3.  $\phi'_{\min}, \phi'_{\max}$  : minimum, maximum angle of shearing resistance
4.  $\lambda$  : compressibility ( $\lambda = C_c/2.3$ )
5.  $l$  : scale parameter
6.  $n$  : number of spiral lines
7.  $m$  : number of radial lines

8.  $x(n)$  : abscissa of the intersections of the spirals and of the edge of the passive zone

Further discussion of factors influencing the choice of input parameters in a given problem is given in Appendix A.

#### 4.4 VALIDITY OF THE COMPUTER PROGRAM

The validity of the numerical results produced by the computer has been established by testing the program with a constant- $\phi$  solution. That is, the  $\phi$ -angle was not allowed to vary throughout the computation but was fixed at a constant value. Various values of the  $\phi$ -angle were tested and results are summarized in Table 4.1 together with results previously obtained by Graham and Stuart (1971). Both sets of numerical results agree within 1% representing slight differences in computer rounding. This confirms the basic validity of the computer program for constant- $\phi$  modelling. This step is needed because the original program was in ALGOL, and had to be translated into WATFIV for use at the University of Manitoba. This comparison does not however confirm the validity of the variable- $\phi$  subroutine calculations which form the new contribution by the author.



## Chapter V

### RESULTS

#### 5.1 BEARING CAPACITY COEFFICIENT FOR FOOTINGS

Since the determination of the slip-line field proceeds from an assumed "known" boundary (edge of the passive zone) towards the initially unknown boundary (edge of the elastic wedge), the coordinates of this end boundary are determined at the very end of the computation and depend on the input parameters. Fig 4.1 showed the position of the end boundary (E), the edge of the passive zone (P) and the direction of the computational process. Point A which is the intersection of the extreme characteristic and of the elastic wedge lies also on the centre-line of the considered footing. Therefore, the actual breadth of the footing is not predetermined but computed after the development of the stress-characteristics from the known passive zone boundary. Due to the symmetry of the problem, the actual breadth of the footing is twice the abscissa of point A. For every set of input parameters  $e_{\min}, e_{\max}, \gamma, \lambda, p'_{\min}, p'_{\max}, \phi_{\min}, \phi_{\max}$  and  $l$ , the computer program calculates a set of coupled results :

1. dimensionless bearing capacity coefficient  $N\gamma$ , where

$$a) N\gamma = 2 \int_0^{8/2} \sigma_v dx / 0.5\gamma B^2, \text{ and}$$

b) where the vertical stress  $\sigma_v = \sigma(1 + \sin\phi \cos^2\psi)$

2. breadth of the footing B (metres)

## 5.2 PASSIVE PRESSURE COEFFICIENT FOR WALLS

Similar to the computation of the bearing capacity coefficients for footings, the determination of the slip-line field behind a retaining passive wall proceeds from the edge of a passive zone towards the end boundary. In this case, the end boundary represents the fully rough back-surface of the wall. Fig 4.2 showed the position of the intersection of the extreme characteristic with the wall. As in the footing case where the breadth is determined a posteriori, the actual height of the wall is determined at the end of the computation process and is represented by point A. Therefore, the ordinate  $z(A)$  of point A is the computed actual height of the wall. As a result of the parametric study outlined in the preceding chapter, every set of input parameters produces :

1. the passive pressure coefficient for a retaining wall

$K_p$  (dimensionless).

a)  $K_p = \int_0^H \sigma_H dz / 0.5\gamma H^2$

b) where  $\sigma_H = \sigma(1 - \sin\phi \cos^2\psi)$

2. the height of the wall H (metres).

### 5.3 SLIP-LINE FIELDS

Due to the number of characteristics, plots of the resulting slip-line fields are unacceptably "dense". For clarity, all the spirals but only half of the radial lines have been computer-drawn. Fig 5.1 shows the slip-line field for the footing case, only half of the total field is drawn, the other part being symmetric about the centre-line of the base of the footing. Fig 5.2 shows the slip-line field behind a failing passive wall.

### 5.4 RESULTS

#### 5.4.1 Bearing capacity coefficient for footings

Chapter 4 described the parametric study that was conducted. Results of the computations are plotted in :

1. semi-log scale<sup>3</sup> : Figs 5.3a,b
2. log-log scale : Figs 5.4a,b where the abscissa ranges from 0.1 m to 10 m.
3. Compressed log-log scale : Figs 5.5a,b where the abscissa ranges from 0.01 m to 10 m.

Figs 5.6a,b present the plot of  $N_\gamma$  versus the compressibility  $\lambda$ .

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<sup>3</sup> a and b refer to the "limited" and to the "broad" range case respectively.

#### 5.4.2 Passive pressure coefficient for walls

Results are plotted in :

1. semi-log scale : Figs 5.7a,b
2. log-log scale : Figs 5.8a,b
3. compressed log-log scale : Figs 5.9a,b

Figs 5.10a,b represent the plots of the passive coefficients for wall as function of the compressibility  $\lambda$  of the sand.

These results are discussed in more detail in Chapter 6.

## Chapter VI

### DISCUSSION OF RESULTS

#### 6.1 BEARING CAPACITY COEFFICIENTS FOR FOOTINGS

##### 6.1.1 Influence of the breadth of the footing

Figs 5.3a,b showed semi-logarithmic plots of the bearing capacity coefficients versus breadth of footing. The bearing capacity coefficients increase rapidly as breadth decreases and flatten almost asymptotically as the breadth increases. This behaviour is observed for both the broad range case and the limited range case. However, it is important to note that the broad range case produces higher values at small sizes ( $N\gamma = 303$  for  $B=0.08$  m,  $\lambda = 0.250$ ) than the limited range case ( $N\gamma = 170$  for  $B=0.10$  m,  $\lambda = 0.250$ ). At large footing sizes, a reversed behaviour is observed: the broad range case gives lower values ( $N\gamma = 44$  for  $B=8$  m,  $\lambda = 0.250$ ) than the limited range case ( $N\gamma = 55$  for  $B=7.8$  m,  $\lambda = 0.250$ ). An explanation of this phenomenon is due to the curvature of the Coulomb-Mohr envelope. In a given range of stresses, the broad range model is associated with higher and lower angles of shearing resistance than the limited range model. Therefore, the stresses at which the soil reaches its minimum  $\phi$ -value under large footings on sands permitting angles of shearing resistance as low as  $28^\circ$  (the

broad range case) are greater than the equivalent stresses under similar footings resting on sands that exhibit angles of shearing resistance as low as  $32^\circ$  (the limited range case). This in turn leads to lower values of the bearing capacity coefficients for the broad range case for large sizes as compared with those obtained for the limited range case.

A similar phenomenon occurs for small-size footings where stresses at which the soil first reaches its maximum value in the broad range case ( $\phi'_{\max} = 45^\circ$ ) are lower than those in the limited range case ( $\phi'_{\max} = 40^\circ$ ). In turn, higher  $N_\gamma$  values are observed for the broader range case.

Another factor attracts attention for both ranges, although it is more acute for the broad range case. Lines representing the behaviour of the sand in Figs 5.4a,b and Figs 5.8a,b for various compressibilities intersect at a breadth around 0.8 m. Thus, for a large footing, say breadth over 0.8 m, the  $N_\gamma$  value associated with a more compressible sand (for instance  $\lambda = 0.250$ ) is lower than the  $N_\gamma$  values associated with a less compressible sand (for instance  $\lambda = 0.050$ ). The reverse behaviour is observed for small footings with breadth say, less than 0.8 m. That is, the  $N_\gamma$  values associated with a more compressible sand ( $\lambda = 0.250$  for instance) are greater than those associated with a less compressible sand ( $\lambda = 0.050$  for instance).

It appears necessary to set upper and lower limits for the values of the  $N\gamma$  coefficients since the mode of failure changes with the size of the footing. It was remarked earlier that large footings fail by 'punching' and small footings along localized shear surfaces (Vesic, 1963). The notion of punching failure is understandable when thinking of the amount of energy required to develop a general shear failure mode in compressible soils, that is to move the soil outside the footing in an upward direction. The upper and lower limits of  $N\gamma$  correspond to cases where only part of the foundation soil has reached limiting equilibrium. In narrow foundations, this part is restricted to the narrow zones in which continuous shear failure has developed. In very large foundations, or on compressible soil, the plastic region is bounded to the side of the footing by soil whose straining has not taken it into plastic behaviour. Controversy may arise in determining the criterion to fix the boundaries between these different failure mechanisms. It has not been possible in the past to approach this question quantitatively. As the changes of behaviours are probably subtle and progressive, experimental evidence is needed. This remains outside the scope of this thesis.

#### 6.1.2 Influence of the compressibility $\lambda$

Both the limited and the broad range cases show similar patterns of behaviour (Fig 5.6). At small sizes  $N\gamma$  increases with compressibility  $\lambda$ . However, for large footings, the

trend is reversed, and  $N_\gamma$  decreases slightly with increasing compressibility. The value is found to remain constant (non-varying with  $\lambda$ ) for a size of 1.3 m in the case of a limited range and for a size of 1.21 m in the broad range case.

The explanation of this behaviour resides in the counteracting influences of the size of the footing and that of the compressibility  $\lambda$  of the sand. For small footings, say  $B=0.5$  m, the likelihood of failing in a rigid-plastic mode diminishes as the soil becomes more compressible. The sand compresses but does not fail and postpones failure at higher stresses. This explains the increases of the  $N_\gamma$  values. However, the same pattern of change is not observed in the case of large footings, say  $B=8$  m. Here, the influence of the compressibility still does exist (leading to a higher  $N_\gamma$ ) but this influence is overridden by the influence of the scale effect and failure envelope curvature. As shown in the preceeding paragraph, for constant  $\lambda$ , the  $N_\gamma$  value decreases as the footing size increases. The combined effect of the influences of both the scale effect and compressibility leads to the observed behaviour, that is an overall decrease of bearing capacity coefficients for large footings. This behaviour is readily apparent in footings at large field scale. It contributes to an understanding of the common perception that small laboratory footings behave differently from field-scale footings.



It is important to note that the behaviour described above in Fig 5.6 is compatible with that observed in the preceding paragraph (Fig 5.4). As mentioned previously in section 6.1, the  $N_\gamma$  values for large footings decrease as the compressibility of the sand increases while the  $N_\gamma$  values of small footings increase with compressibility (lines 1,2,3,4 of Figs 5.4a,b).

Although it is difficult to separate the influence of the compressibility from that of pressure-dependent scale effects, this study shows that the two separate influences should not be neglected when an accurate evaluation of the  $N_\gamma$  coefficients is required.

## 6.2 COEFFICIENTS OF PASSIVE PRESSURE FOR WALLS

Computation of the coefficients of passive pressure for walls are very similar to those carried out for the estimation of the bearing capacity coefficients for footings. The patterns of the results are also similar. Consequently, a detailed discussion of these results and an explanation of the observed behaviour will not be repeated. The main points will be highlighted here and reference made to the previous detailed discussion of the bearing capacity coefficients.

### 6.2.1 Influence of the height of the wall

Both the limited range and the broad range exhibit the same general characteristics:

1. For a given  $\lambda$ , the coefficient of passive pressure for walls  $K_p$  decreases with the height of the wall (Figs 5.7a,b).
2. Curves corresponding to various compressibilities intersect for a wall height of 1.8 m. This is apparent on Figs 5.8a,b.
3. For large wall heights, say over 1.8 m, the  $K_p$  coefficients associated with more compressible sand are smaller than those associated with less compressible sands (lines 1,2,3,4 on Figs 5.7;5,8).
4. For small wall heights, say less than 1.8 m, the  $K_p$  coefficients associated with more compressible sand are larger than those associated with less compressible sands (lines 4,3,2,1 on Figs 5.6;5.8).

Similar to the case of the bearing capacity problem, the general tendency is for the  $K_p$  coefficient to decrease with increasing wall size. This decrease is affected by the compressibility of the sand. Differences between the various lines representing the compressibility of the sand are also more marked for the broad range case.

### 6.2.2 Influence of the compressibility

The general characteristics for both the limited range and the broad range case (although more marked for the broad one) are:

1. For small heights, less than 1.8 m, the  $K_p$  coefficient increases with increasing compressibility. This is due to a lower likelihood of the sand of failing in a rigid-plastic mode. The sand compresses and postpones failure to higher stresses.
2. For large wall heights, over 1.8 m, the  $K_p$  coefficient decreases with increasing compressibility. This is due to the counteracting influences of pressure dependent size increases leading to lower values and of increased compressibilities leading to larger values. The overall observed behaviour is a decrease of the  $K_p$  coefficient with increasing compressibility.

## 6.3 CONTOURS OF THE ANGLE OF SHEARING RESISTANCE

### 6.3.1 Distribution of the $\phi$ -angle under a footing

Figs 6.1a,b show the contours of the angle of shearing resistance under 0.1 m wide and 2 m wide footings. In both cases,  $\lambda=0.250$  and the broad range values are considered. The contours are represented by the dash lines. For both footing sizes, a decrease of roughly  $10^\circ$  is observed between the edge of the passive zone and the centre line. As expected, values of the  $\phi$ -angle are greater for a small footing

(38°-48°) than those for a large footing (32°-42° for the 2 m footing); this is due to the larger pressures under the large footing.

It is interesting to compare the  $N\gamma$  values obtained from a constant  $\phi$ -solution (Table 6.1) where  $\phi$  is now taken as the average mobilized shearing resistance throughout the domain under stress, and corresponding  $N\gamma$  values obtained from the present analysis based on the size of the footing and on the compressibility of the sand. The importance of this step in selecting appropriate  $\phi$ -values was introduced earlier by Graham and Stuart (1971).

For small footings, the  $N\gamma$  value corresponding to a constant- $\phi$  analysis is 224 for  $\phi=44^\circ$  and 330 for  $\phi=46^\circ$ . The  $N\gamma$  value obtained from Fig 5.6b (broad range case,  $\lambda=0.250$ ,  $B=0.1$  m) is 300. For the large footing, the respective figures are 78 for the constant- $\phi$  solution with  $\phi=38^\circ$  and 75 from Fig 5.6b. Depending on the average value of  $\phi$  which is taken to model the failure domain, the  $N\gamma$  values from the two methods are clearly similar. It therefore appears that a conventional constant- $\phi$  analysis may produce accurate results as long as the 'correct'  $\phi$  is chosen for the analysis. Since in practice it is extremely difficult to estimate a 'correct' angle of shearing resistance for a domain under stress, the present analysis based on the size of the footing and on the compressibility of the sand is inherently sound.

### 6.3.2 Distribution of $\phi$ behind a passive wall

Figs 6.2a,b show the contours of the angle of shearing resistance behind passive walls for heights of 0.1 m and 2.0 m. In both cases, the variation of  $\phi$  in the failure zone is of the order of  $8^\circ$ . However, the mobilized  $\phi$  angles for a 2 m wall are much smaller ( $36^\circ$ – $44^\circ$ ) than those obtained for the 0.1 m wall ( $41^\circ$ – $49^\circ$ ). This is due to the larger pressures involved in the case of the larger wall.

As was shown for the bearing capacity problem, it is interesting to compare the  $K_p$  values that arise from the present analysis<sup>4</sup> with those obtained for a constant  $\phi$ -solution where  $\phi$  is taken as the average value of the angle of shearing resistance throughout the failing domain. Table 6.2 summarizes these values. The results of the two analyses are reasonably close but depend strongly on the average value of the  $\phi$ -angle which is taken. Like the bearing capacity problem, it is apparent that estimation of the average  $\phi$ -angle of the failing domain remains difficult. The critical state analysis presents a potentially easier alternative.

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<sup>4</sup> referred as a 'critical state' analysis.

#### 6.4 COMPARISON WITH OTHER THEORETICAL RESULTS

Fig 6.3 and 6.4 show the comparisons for walls and for footings with results obtained by Graham and Stuart (1971) who used the De Beer (1963)  $\phi$ -pressure relationship. Only the broader- and the narrower-range results of the present analysis are plotted. Line (3) represents the results obtained for the broader range case with  $\lambda=0.250$ . Line (2) corresponds to the results obtained for the limited range case with  $\lambda=0.050$ . Lines (1) and (4) represent results provided by the Graham and Stuart study for a loose sand and for a dense sand respectively.

It is important to notice that for both the wall problem and the footing one, the curves corresponding to the narrower range (line 2) remain in the Graham and Stuart range and almost parallel to the loose sand. Ideally, if the variation of the  $\phi$ -angle did not have any influence, curves corresponding to a medium sand (lines 2) should approximatively be situated in the middle of the Graham and Stuart range. Therefore, the fact that lines (2) are offset towards the loose sand brings some evidence that the variation of the  $\phi$ -angle does alter the computation of the ultimate stability of the domain under stress.

Lines (3) representing the broader range where the  $\phi$ -angle is allowed to vary from  $28^\circ$  to  $45^\circ$  must be understood as an accentuation of what lines (2) represent (narrow range,  $\phi = 32^\circ - 40^\circ$ ). In view of this fact, it is therefore expected

that lines (3) be closer to the Graham and Stuart dense state for small sizes and closer to the Graham loose state for large sizes. This is what Figs 6.3 and 6.4 reveal. However, since the minimum  $\phi$ -angle in the broader range in the present study is as low as  $28^\circ$ , compared with the Graham and Stuart minimum of  $32^\circ$ , the range of behaviour of the sand represented by lines (3) is broader than the Graham and Stuart range. Consequently, bearing capacity coefficients for large footings and passive pressure coefficients for large walls are lower than those obtained by Graham and Stuart.

## 6.5 COMPARISON WITH EXPERIMENTAL RESULTS

### 6.5.1 Footing case

Experimental results provided by Vesic (1973), Graham and Stuart (1971) and Meyerhof (1951) have been plotted in Fig 6.5. In order to compare these results with the range of values produced by the critical state analysis, the upper limit of the set of  $N_\gamma$  values produced (broad range,  $\lambda=0.05$ ) as well as the equivalent lower limit (limited range,  $\lambda=0.250$ ) have also been plotted in Fig 6.5. Clearly, the general trend of experimental values correlate reasonably well with the form of the solution arising from the critical state analysis. In particular, the range of the theoretical solutions which have been primarily carried out for medium dense sands covers closely the range of experimental values obtained for medium relative density sands.

However, it is clear that it is still not possible without extensive additional testing to predict from this work, the bearing capacity of a footing on a particular sand. Since this thesis was prepared, the author has been made aware of a recent literature review in the paper: Amar, S., Bague-  
lin, F., and Canepa, Y., 1984 : "Etude experimentale du com-  
portement des fondations superficielles", Annales de l'insti-  
tut technique du batiment et des travaux publics, sols et  
fondations 189 # 427, Paris, France.

#### 6.5.2 Wall case

Similar to the footing problem, the broadest range of  $K_p$  values obtained from the critical state analysis have been plotted against experimental data (Fig 6.6). The upper limit of the range is described as the set of values yielded by the broad range case analysis ( $\lambda=0.05$ ), whereas the lower limit correspond to the limited range ( $\lambda=0.250$ ). As for the experimental data, results obtained by Kerisel (1972), Horn (1967) and Graham (1971) have been considered. Even if the experimental values seem quite dispersed, a general trend still appears, that is, a decrease of the  $K_p$  values with in-  
creasing wall height. Like the bearing capacity problem, the range of theoretical values which is primarily valid for me-  
dium dense sands broadly covers the range of experimental data.



In both cases ( $N\gamma$  and  $K_p$  values), the theoretical solutions and field data encompass scale effect influences whereas small scale laboratory tests do not in general. It therefore appears that the scattering of the experimental values outside the theoretical range is not an abnormality of the sand behaviour but represents tests on sands which may be very angular, rounded, or well-graded, and which lie outside the ranges of behaviour shown in Figs. 4.3a,b.

#### 6.6 SINGULARITY OF SIZE 1m

Figs 5.5a,b and Fig 5.9a,b showed respectively the plots of the bearing capacity coefficients against the breadth of the footing, and the passive pressure coefficients against the height of the wall. These are plots where abscissas represent a broader range (0.01,0.1,1.0,10 m) compared with the preceding plots (Figs 5.4 and 5.8) where the range was (0.1,1.0,10.0 m). This has the effect of compressing the diagrams in the X-direction and thus accentuating all the possible singularities in the curves. Such a singularity is situated for a size of 1 m (height of the wall or breadth of the footing). As is apparent in Figs 5.5 and 5.9, the  $N\gamma$  values and the  $K_p$  values exhibit localized slope discontinuities at size 1 m. These discontinuities can be explained by the way the  $N\gamma$  values and the  $K_p$  values are calculated.

For a passive wall,  $K_p$  is evaluated from the equation :

$$P_p = 0.5 \gamma H^2 K_p$$

For a footing,  $N\gamma$  is calculated from :

$$Q = 0.5 \gamma B^2 N \gamma$$

Both equations have the same form :

$$y = 0.5 k \gamma x^2$$

The function  $x^2$  shows a singularity when the variable  $x=1$ , that is  $x^2 = x (1^2 = 1)$ . Therefore it is important to consider the singularity which is introduced by squaring the variable  $x$  which as explained above represents the breadth of the footing or the height of the wall.

Consider :  $x^2 - x = x(x-1)$  and notice that  $K_p$  and  $N\gamma$  are evaluated as a function of the form :  $k / |x(x-1)|$ .

The function  $y=1 / x(x-1)$  is plotted in Fig 6.7. Three intervals can be distinguished:

1.  $0 < x < 0.5$  asymptote at  $x=0$  with  $y$  decreasing.
2.  $0.5 < x < 1$  asymptote at  $x=1$  with  $y$  increasing.
3.  $1 < x$  asymptote at  $y=0$  with  $y$  decreasing

These patterns of slope changes are similar in some respects to those observed in Figs 5.5 and 5.9 :

1.  $0 < \text{size} < 0.5 \text{ m}$   $N\gamma$  and  $K_p$  decrease rapidly
2.  $0.5 < \text{size} < 1\text{m}$   $N\gamma$  and  $K_p$  decrease slowly and a local slope discontinuity obtained at size 1 m.
3.  $1 < \text{size}$   $N\gamma$  and  $K_p$  decrease first rapidly and seem to tend to an asymptotic value.

The mathematical function  $y=1/|x(x-1)|$  explains the behaviour observed for  $K_p$  and  $N\gamma$ . The local slope discontinuity at size 1 m is therefore not due to some particular feature of the sand response but is introduced artificially by the normalizing equations which serve the purpose of defining the bearing capacity and passive pressure coefficients in dimensionless terms.

These results have shown that an estimation of bearing capacity and passive pressure coefficients based solely on a constant angle of shearing resistance seriously underestimate the complexity of the problem. The scale effect factor plays notably such an important role that practicing engineers use some empirical rules to take it into account. Alternatively, they neglect the usefulness of calculating failure loads in sands for full size structures. Bjerrum (1973) mentions such an empirical rule by stating that calculation of large footings for gravity structures in the North Sea is done by applying a reduction of  $\phi$  of the order  $n.1^\circ$  to  $n.4^\circ$  for each  $10^n$  reductions of the footing size ( $n$  is an integer). This rule arises from the qualitative work presented by Graham and Pollock (1972). It thus appears that current practice acknowledges the importance of known but poorly determined factors such as those described in this thesis.

## Chapter VII

### CONCLUSIONS AND FURTHER RESEARCH

#### 7.1 CONCLUSIONS

Conclusions drawn from this study are :

1. The advent of more powerful and more convenient computers and peripherals makes it possible to investigate the development of the static equilibrium equations combined with a new sand model (Critical State model).
2. The angle of shearing resistance is only one component among others that are necessary to get accurate computations of bearing capacity coefficients for footings and passive pressure coefficients for walls.
3. The two most important additional parameters that are taken into account in this study are :
  - a) the footing breadth or the wall height.
  - b) the compressibility of the sand ( $\lambda$  or  $C_c$ ).
4. Bearing capacity coefficients for footings and passive pressure coefficients for walls have been expressed in dimensionless terms. For a given sand, they decrease with increasing size (breadth of the footing or height of the wall). This is due to the effect of higher pressures associated with larger

walls or footings which decrease the mobilized angle of shearing resistance.

5. The variation of the angle of shearing resistance has been found to be approximately  $10^\circ$  in the rigid-plastic domain beneath a footing (Fig. 6.1). The equivalent variation of the  $\phi$ -angle behind a passive wall is  $8^\circ$  (Fig. 6.2).
6. The produced solutions correlate well with other theoretical and experimental results (Figs. 6.3 to 6.6).

## 7.2 SUGGESTIONS FOR FURTHER RESEARCH

1. The behaviour of rigid-plastic zones has been described using stress-characteristics. It would be of considerable interest to model with a Finite Element mesh the trapped elastic zone beneath a footing and connect it to the stress-characteristic solution. Similarly, a major further development would be to connect a stress-characteristic solution to the compressed non-failing field to the side of a footing on strain-hardening soil when "punching" is occurring.
2. Contours of the mobilized angle of shearing resistance have been drawn. The critical state analysis could also permit the computation of the variation of the void ratio. This could be investigated to draw the contours of the void ratio or volume strain be-

neath a footing or behind a passive wall. In turn, the variation of the density could be addressed.

3. The rigid-plastic assumption does not take into account volume strains prior to failure. Experimental tests could determine some index of "compressibility" that could be used to factor the  $N\gamma$  and  $K_p$  values for compressible sands. Alternatively, approximate load-displacement relationships can be derived by close attention to boundary conditions.
4. Further laboratory tests should be performed to further confirm the validity of the critical state model. This, in turn, will reinforce confidence in the critical state analysis.

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$\phi$ -angle	$N_Y$ -value from Graham and Stuart	$N_Y$ -value from the computer program
$30^\circ$	23	23.17
$32^\circ$	31	31.25
$34^\circ$	45	45.60
$36^\circ$	68	69.20
$38^\circ$	98	98.90
$40^\circ$	146	147.90
$42^\circ$	240	242.97
$\phi$ -angle	$K_p$ -value from Graham and Stuart	$K_p$ -value from the computer program
$30^\circ$	5.6	5.65
$34^\circ$	7.8	7.85
$38^\circ$	11.3	11.60

TABLE 4 .1 - CHECK OF NUMERICAL ACCURACY OF THE PROGRAM,  
COMPARISON OF CONSTANT-  $\phi$  SPECIFICATION WITH  
EARLIER RESULTS FROM GRAHAM AND STUART (1971) .

	Average- $\phi$ in the domain under stress	$N_Y$ value from Fig 5.4 b	$N_Y$ value for a constant- $\phi$
0.10 m footing	44 <sup>0</sup> 46 <sup>0</sup>	300	224 330
2.0 m footing	38 <sup>0</sup>	75	78

FIG 6.1 - COMPARISON OF  $N_Y$ -VALUES FROM THE CRITICAL STATE MODEL  
(VARIABLE  $-\phi$ ) WITH EQUIVALENT VALUES USING A CONSTANT  
AVERAGE MOBILIZED  $\phi$  IN FAILURE DOMAIN.

	Average- $\phi$ in the domain under stress	Critical state analysis Fig 5.8 b	$K_p$ value for a constant- $\phi$ Caquot, Kerisel (1956)
0.10 m wall	$45^0$	25.5	25
2.0 m wall	$40^0$	12.5	14

FIG 6.2 - COMPARISON OF  $K_p$ -VALUES FROM THE CRITICAL STATE MODEL  
(VARIABLE- $\phi$ ) WITH EQUIVALENT VALUES USING A CONSTANT  
AVERAGE MOBILIZED  $\phi$  IN FAILURE DOMAIN.

Classification	$\gamma_{min}$	$\gamma_{max}$	$e_{min}$	$e_{max}$	Max. size	$D_{10}$	$C_u$	$C_c$	F
SP-SM	90	108	0.54	0.84	#16	.058	6.0	2.2	.555
SM	75	97	0.83	1.36	3/4"	.0065	31	5.5	.638
SP	92	112	0.48	0.80	#4	.15	3.0	.93	.667
SP	93	113	0.46	0.77	1 1/2"	.16	2.4	.92	.674
SP	95	116	0.43	0.74	#4	.30	3.7	1.0	.721
SP-SM	92	113	0.46	0.80	3/4"	.08	3.0	.88	.739
SP	85	107	0.54	0.94	#30	.10	2.3	1.3	.740
SP	97	118	0.40	0.70	1 1/2"	.11	3.2	1.2	.750
SP	99	120	0.38	0.67	1 1/2"	1.8	4.4	.76	.763
SM-ML	83	108	0.62	1.11	#4	.012	8.3	1.5	.790
SP-SM	79	103	0.60	1.08	#30	.09	2.4	1.5	.800
SP	103	124	0.33	0.60	3/4"	.17	5.0	.75	.818
SM	105	126	0.31	0.57	5"	.02	350	.30	.838
SP-SM	87	112	0.48	0.90	#4	.08	3.0	1.3	.875
SM	82	108	0.54	1.02	#16	.023	6.5	1.4	.889
SW-SM	95	119	0.39	0.74	3"	.05	10	1.4	.897
SP	98	122	0.36	0.69	#4	.37	5.1	1.2	.917
SW-SM	98	125	0.34	0.71	3"	.07	6.8	1.0	1.088
SP-SM	97	124	0.33	0.70	3/4"	.10	5.0	1.4	1.121
SP-SM	84	115	0.44	0.97	1 1/2"	.085	4.7	1.4	1.205
SP-SM	94	123	0.34	0.76	1 1/2"	.12	4.4	1.3	1.235
SM	99	128	0.31	0.70	3"	.02	240	1.8	1.258
SP-SM	80	114	0.44	1.06	#16	.07	3.7	1.6	1.409
SW-SM	80	116	0.42	1.07	1 1/2"	.074	6.6	2.4	1.547
SM	83	120	0.38	0.99	#4	.015	26	6.1	1.605
SM	102	134	0.23	0.62	3/4"	.01	120	1.9	1.695
GN-GM	113	127	0.31	0.47	3"	.14	86	1.2	.517
GP-GM	112	129	0.32	0.52	3"	.03	200	.50	.625
GW-GM	116	133	0.26	0.44	5"	.17	171	2.2	.692
GP-GM	110	128	0.30	0.51	3"	.11	191	15	.700
GP-GM	117	133	0.24	0.41	5"	.125	160	4.0	.708
GW-GP	111	130	0.27	0.49	3"	.20	105	7.5	.815
GP	116	134	0.23	0.43	5"	.27	111	6.2	.870
GW	119	139	0.24	0.45	3"	.51	45	2.2	.875
GW	120	139	0.20	0.39	3"	.45	51	1.6	.950
GW	119	139	0.21	0.41	3"	.18	94	1.1	.952
GW	111	132	0.25	0.49	3"	2.9	9.7	1.8	.960
GP	115	136	0.22	0.44	5"	.38	29	.81	1.000
GP	114	135	0.22	0.45	3"	2.0	11	.77	1.045
GW-GM	121	141	0.19	0.39	3"	.30	77	2.3	1.052
GM	122	141	0.17	0.36	1 1/2"	.025	381	3.0	1.118
GW-GM	114	137	0.21	0.45	3"	.60	16	1.2	1.143
GW	112	138	0.20	0.48	3"	2.0	12	1.3	1.400
GW	109	137	0.21	0.52	3"	2.0	14	2.6	1.476
GP	114	140	0.18	0.45	3"	1.7	10	.76	1.500
GM	101	132	0.25	0.64	1 1/2"	.03	260	12	1.560
GW-GM	111	139	0.19	0.49	3"	1.8	13	2.3	1.578
GP	115	142	0.17	0.44	3"	.31	87	8.2	1.588
GW	123	146	0.13	0.34	3"	.21	124	1.1	1.615
GW-GM	110	139	0.19	0.50	5"	.42	43	2.1	1.631
GW-GM	115	142	0.17	0.45	3"	.15	133	1.1	1.647
GP-GM	112	140	0.18	0.48	3"	.42	26	4.2	1.667
GW-GM	112	140	0.18	0.48	5"	.25	56	1.0	1.667
GW-GM	114	142	0.16	0.45	3"	1.2	15	1.7	1.812
GP	112	141	0.17	0.48	3"	1.4	7.1	.73	1.823
GW-GM	118	147	0.12	0.40	3"	1.3	19	1.1	2.333

TABLE A.1 - COMPACTIBILITY(F) OF COHESIONLESS

SOILS (WHERE  $F = (e_{max} - e_{min}) / e_{min}$ ),  
AFTER WINTERKORN AND FANG (1975).



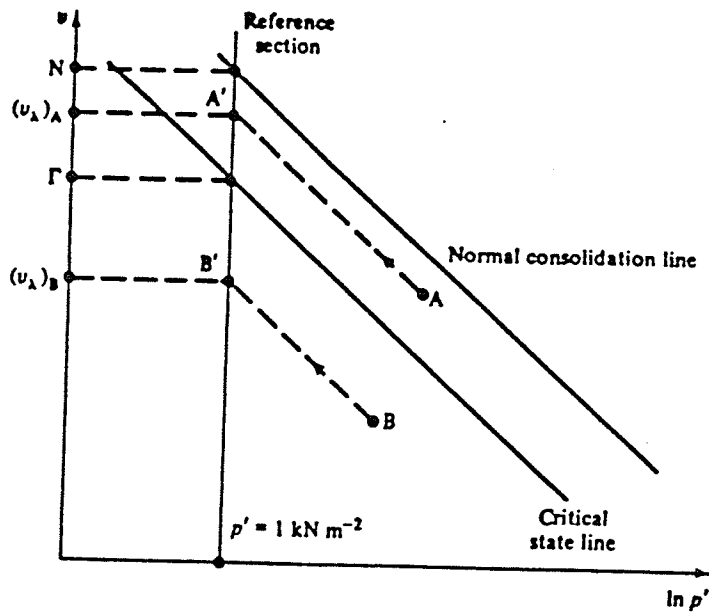


FIG 2.2 - THE NORMAL CONSOLIDATION LINE (NCL), CRITICAL STATE LINE (CSL) AND THE "REFERENCE SECTION" AFTER ATKINSON AND BRANSBY (1978)

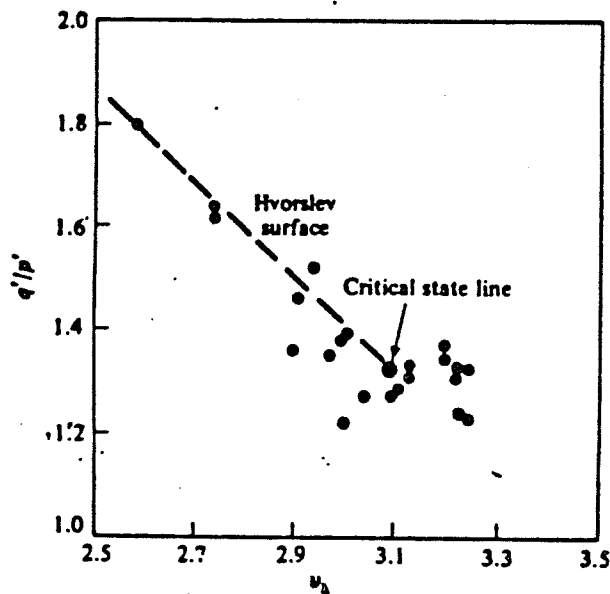


FIG 2.3 - NORMALIZED PLOT OF THE CRITICAL STATE MODEL FOR SAND, AFTER ATKINSON AND BRANSBY (1978)

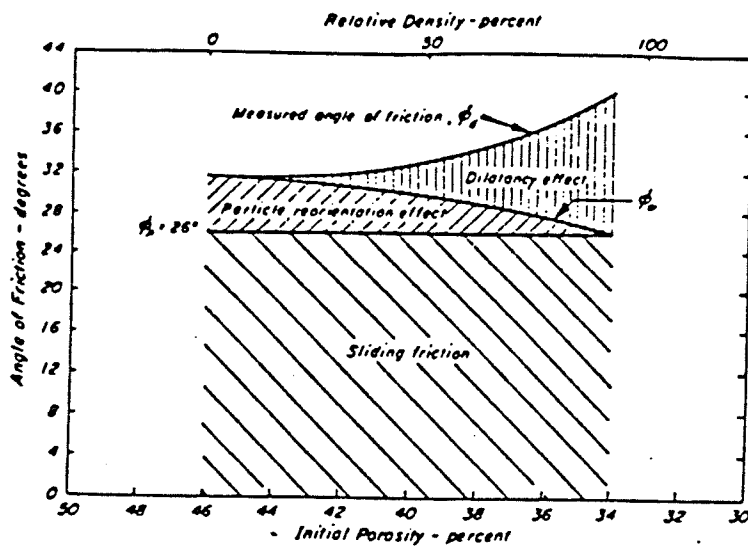


FIG 2.4 - COMPONENTS OF STRENGTH OF SAND, AFTER ROWE (1967).

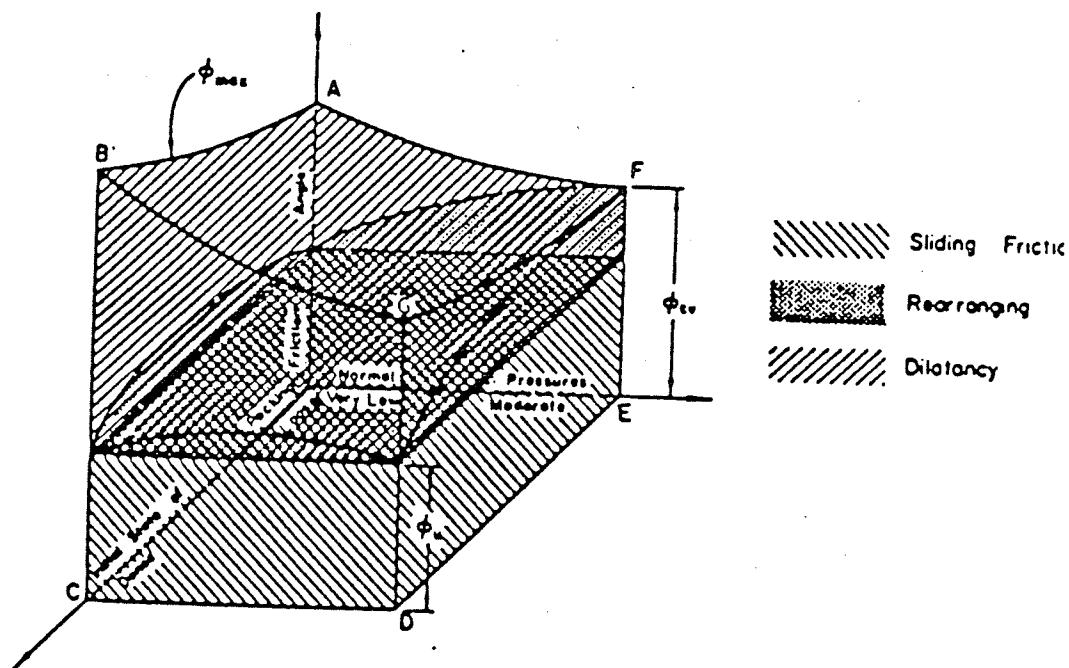


FIG 2.5 - THREE-DIMENSIONAL REPRESENTATION OF ANGLE OF SHEARING RESISTANCE, AFTER PONCE AND BELL (1971).



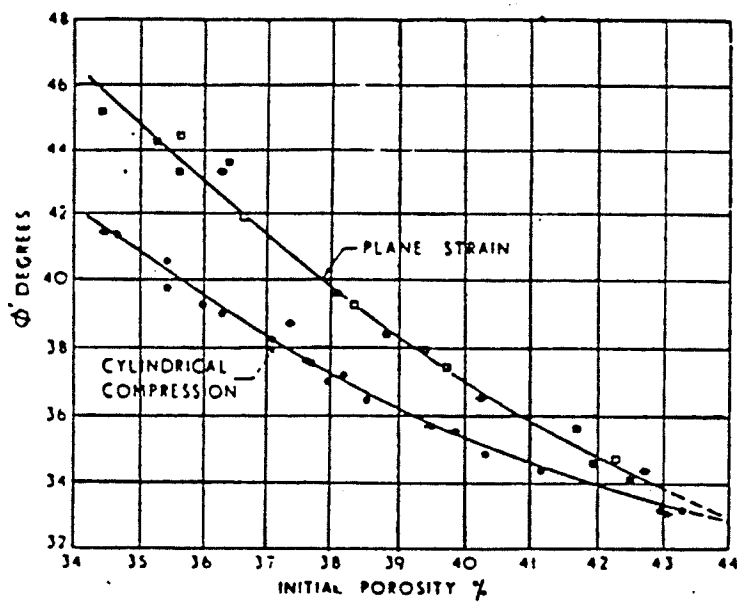
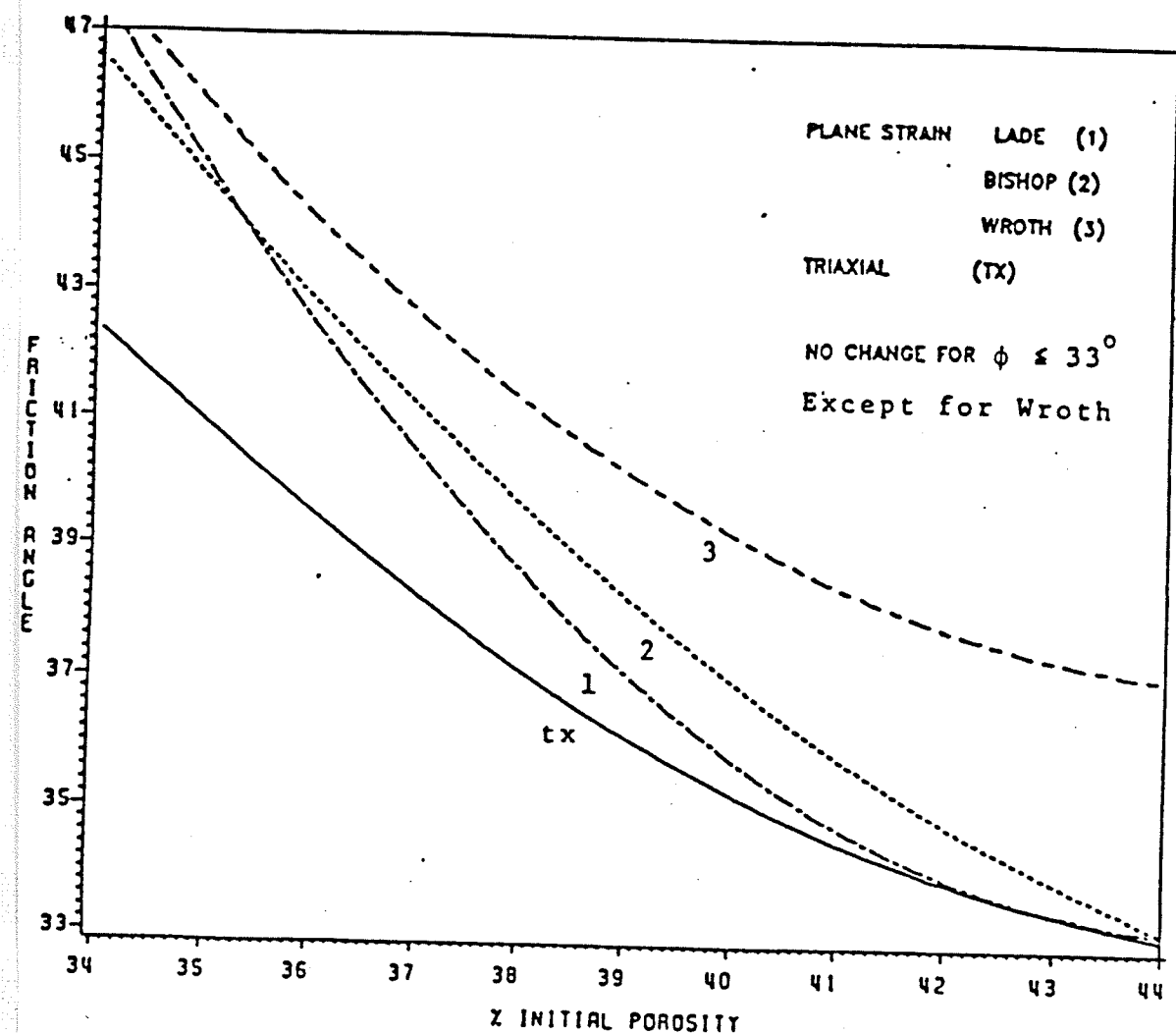
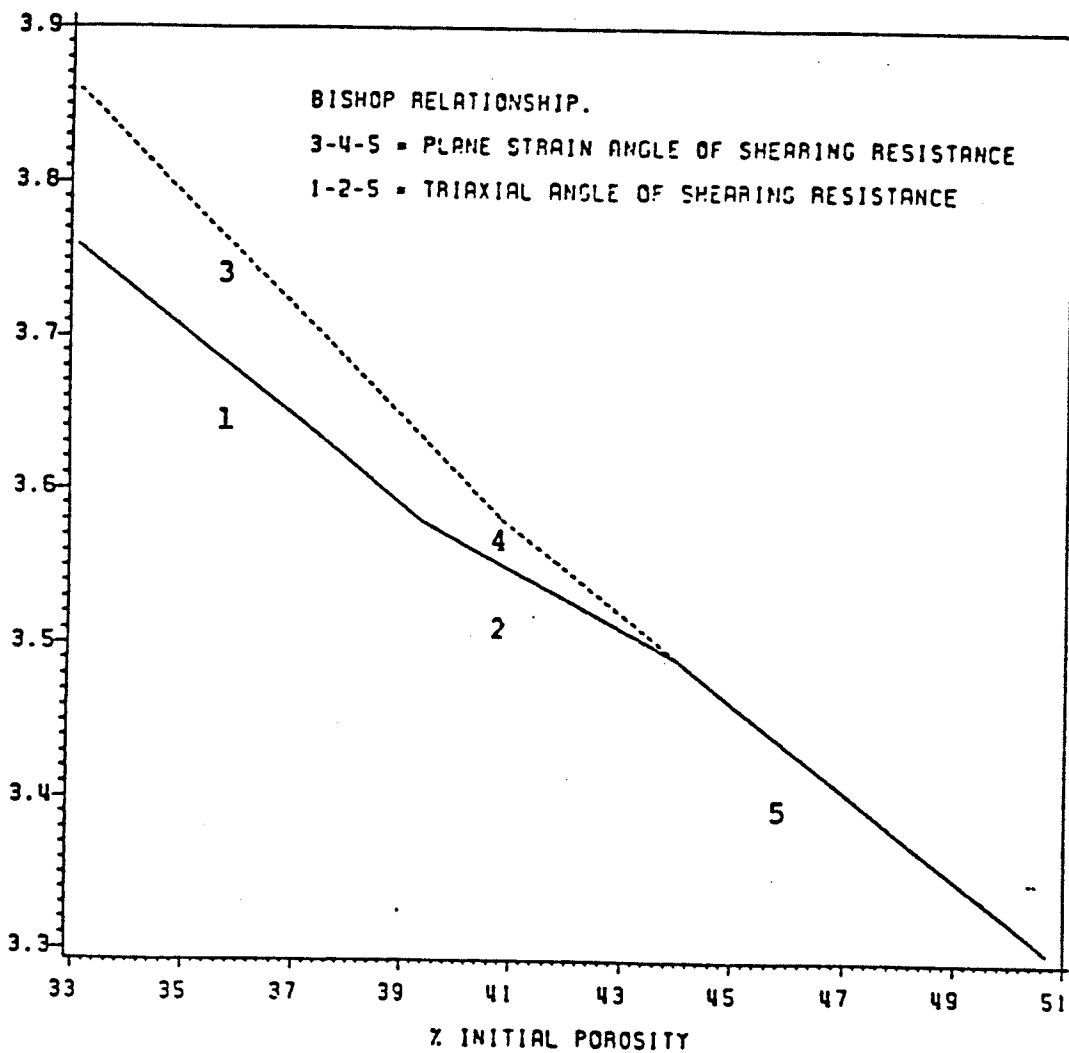


FIG 2.6 - COMPARISON OF RESULTS OF DRAINED PLANE STRAIN AND CYLINDRICAL COMPRESSION TESTS ON BRASTED SAND, AFTER CORNFORTH (1961), PRESENTED BY BISHOP (1966)



G 2.7 - COMPARISON OF VARIOUS PLANE STRAIN  $\phi$ -ANGLES  
 VERSUS TRIAXIAL  $\phi$ -ANGLE RELATIONSHIPS



G 2.8 - THE BISHOP (1966)  $\phi_{ps} - \phi_{tx}$  RELATIONSHIP, LOG-LOG SCALES.

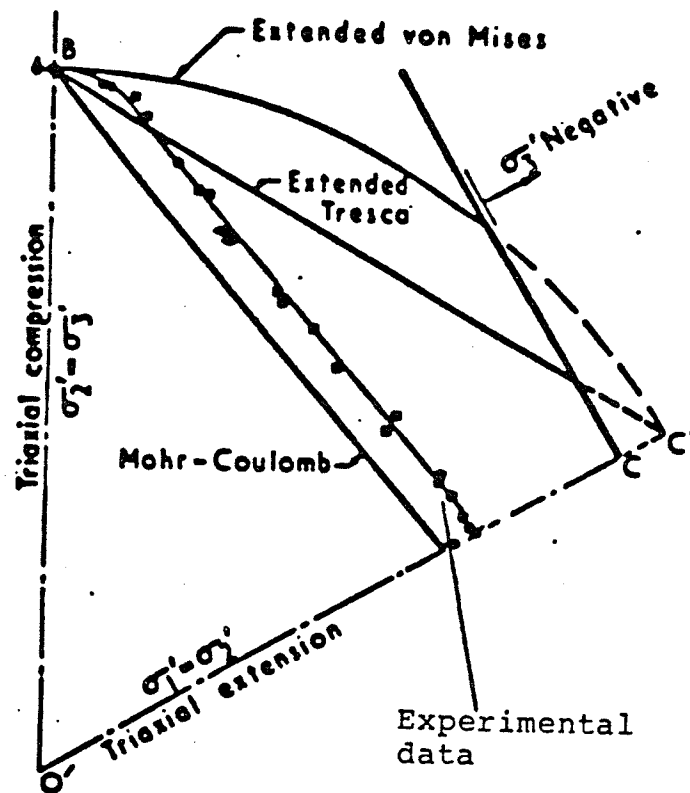


FIG 3.1 - COMPARISON OF VARIOUS YIELD CRITERIA,  
After Green and Bishop (1969)

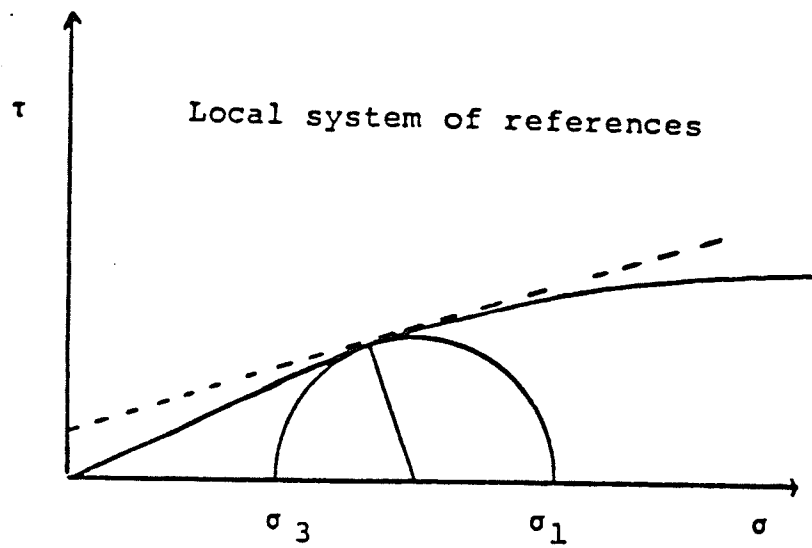


FIG 3.2 - LOCAL REFERENCE SYSTEM OF STRESS AND STRENGTH COORDINATES

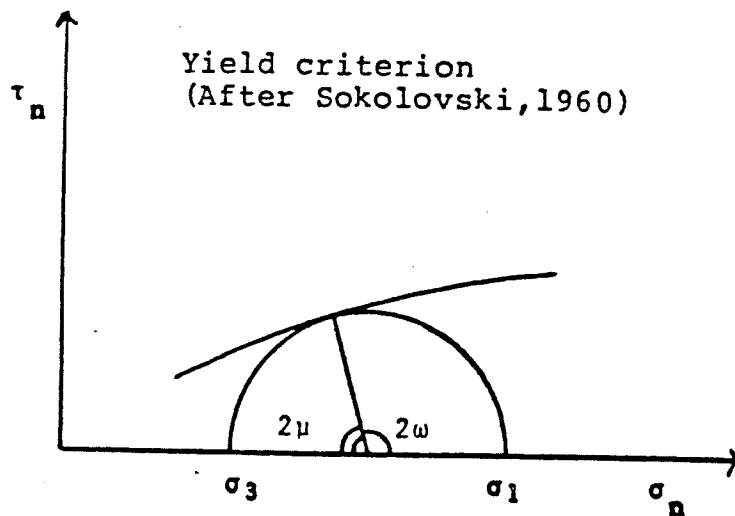


FIG 3.3 - YIELD CRITERION, AFTER SOKOLOVSKI (1960)

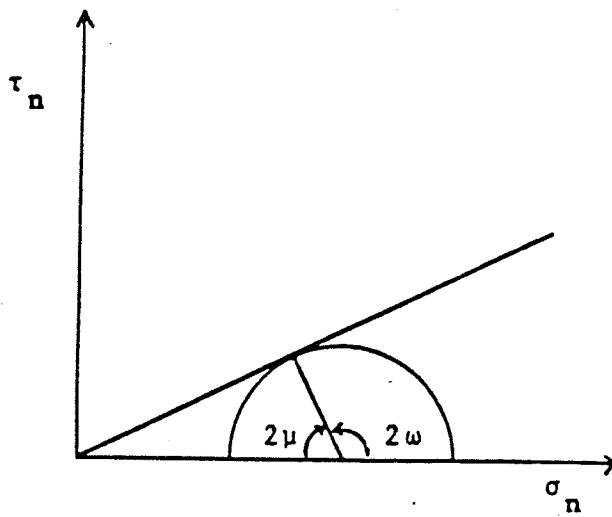


FIG 3.4 - GENERAL RELATIONSHIP BETWEEN  $\mu$  AND  $\omega$

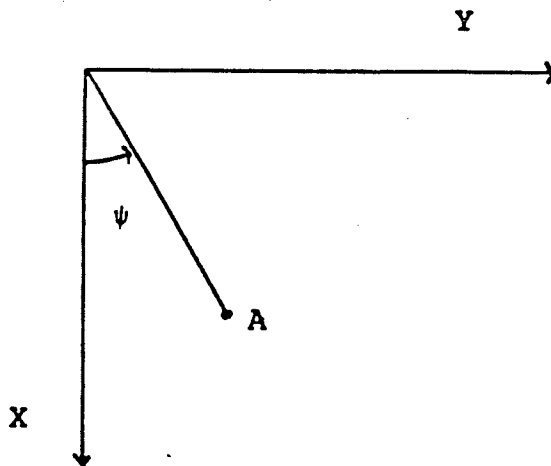


FIG 3.5 - GENERAL SYSTEM OF COORDINATES

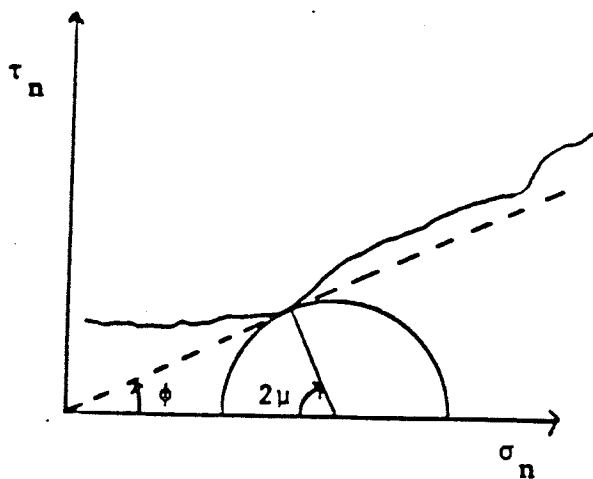


FIG 3.6 - DEFINITION OF  $\phi$

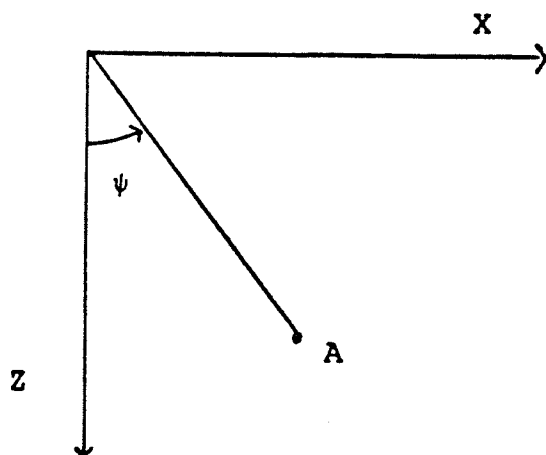


FIG 3.7 - AUTHOR'S SYSTEM OF COORDINATES

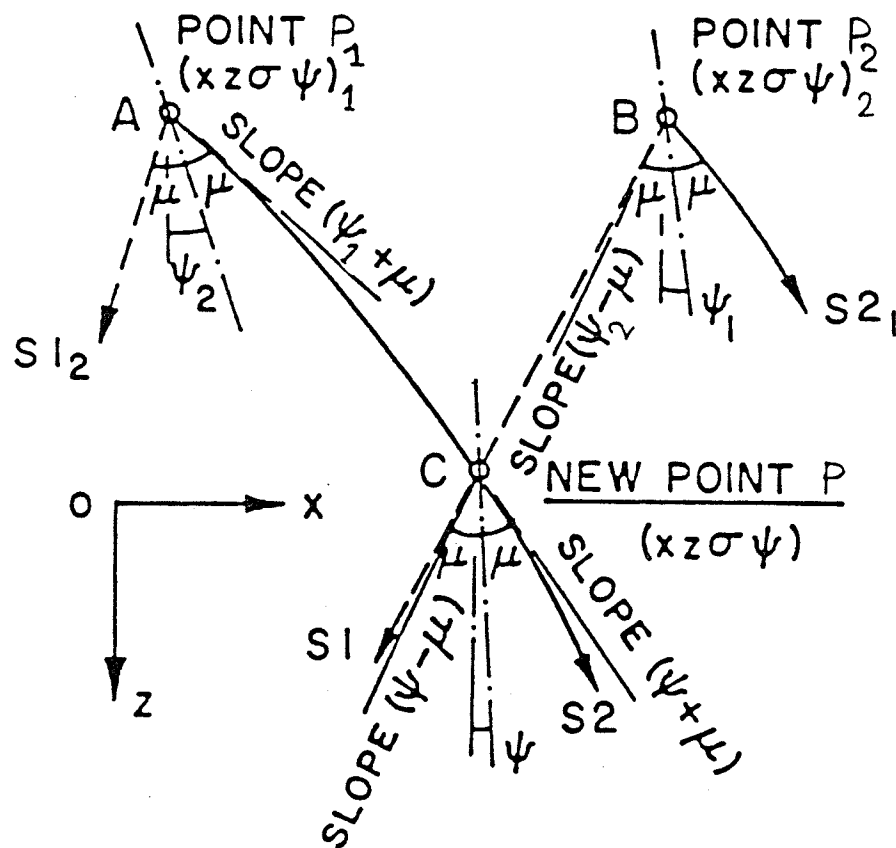


FIG 3.8 - COMPUTATION OF NEW POINT  $P$  FROM  
 KNOWN POINTS  $P_1, P_2$ ,  
 (AFTER GRAHAM, 1968)



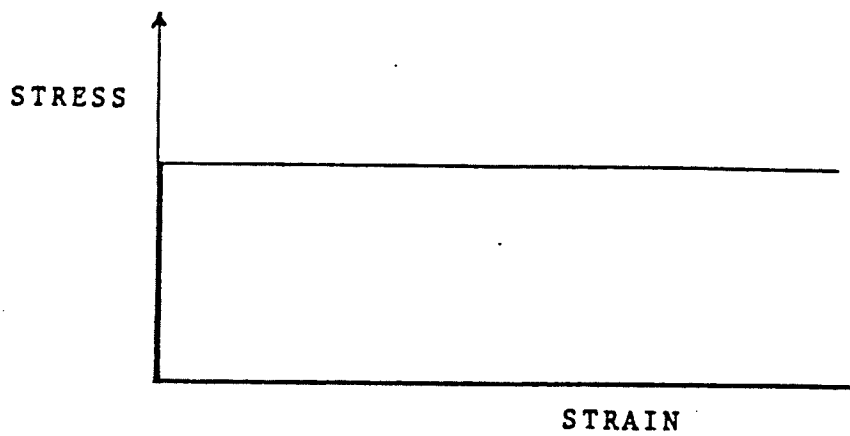


FIG 3.9 - THE RIGID-PLASTIC SOIL MODEL

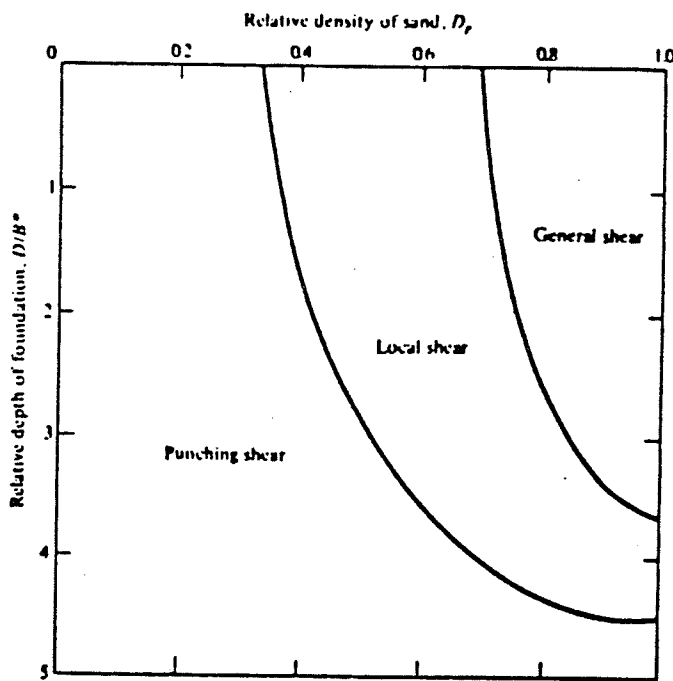


FIG 3.10 - MODES OF FAILURE OF MODEL FOOTINGS IN SAND, AFTER VESIC (1963), AS MODIFIED BY DE BEER (1970)

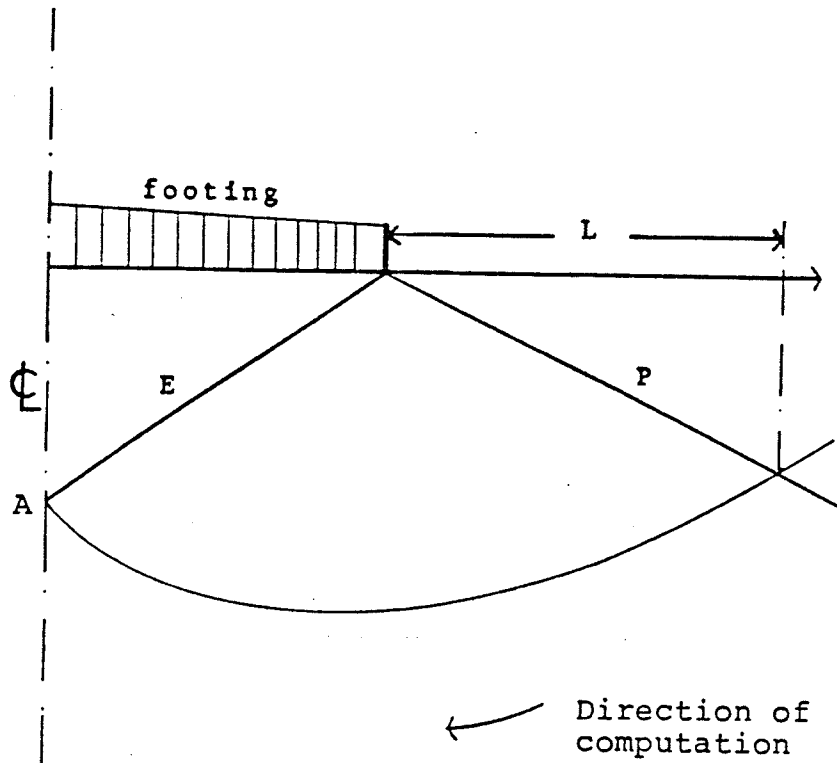


FIG 4.1 - BASIC DEFINITION OF TERMS,  
THE FOOTING PROBLEM

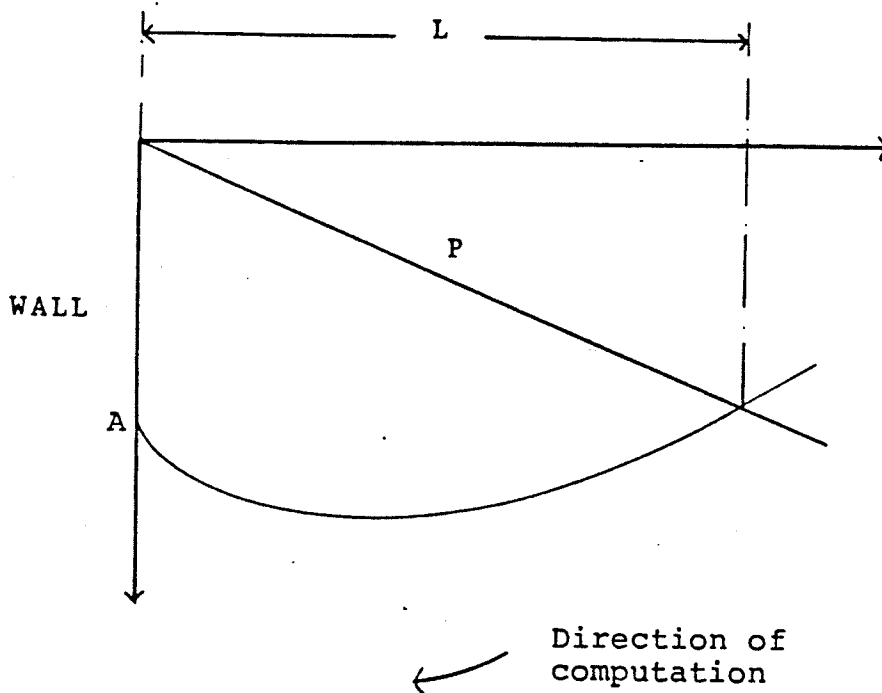


FIG 4.2 - BASIC DEFINITION OF TERMS,  
THE PASSIVE WALL PROBLEM

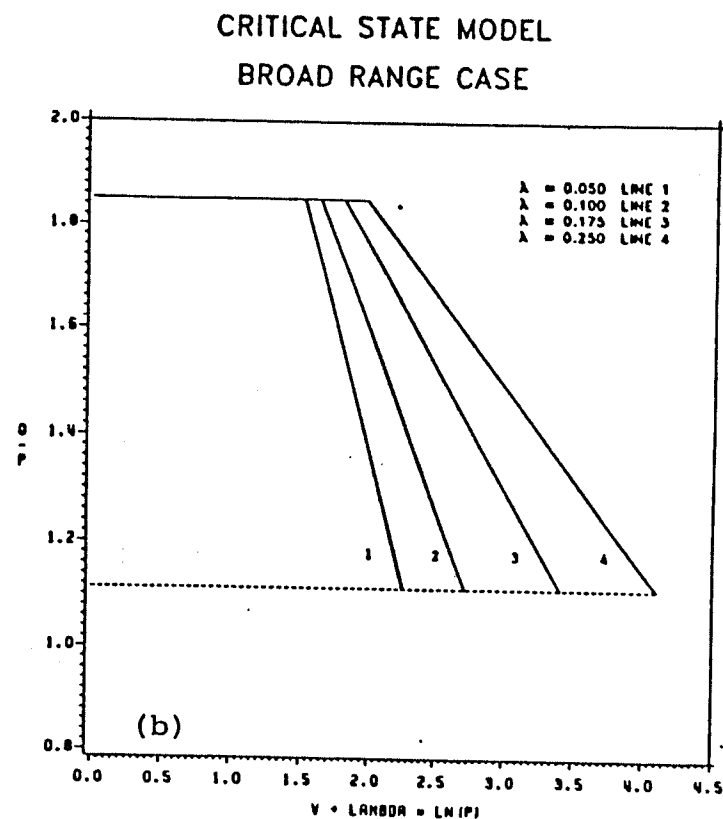
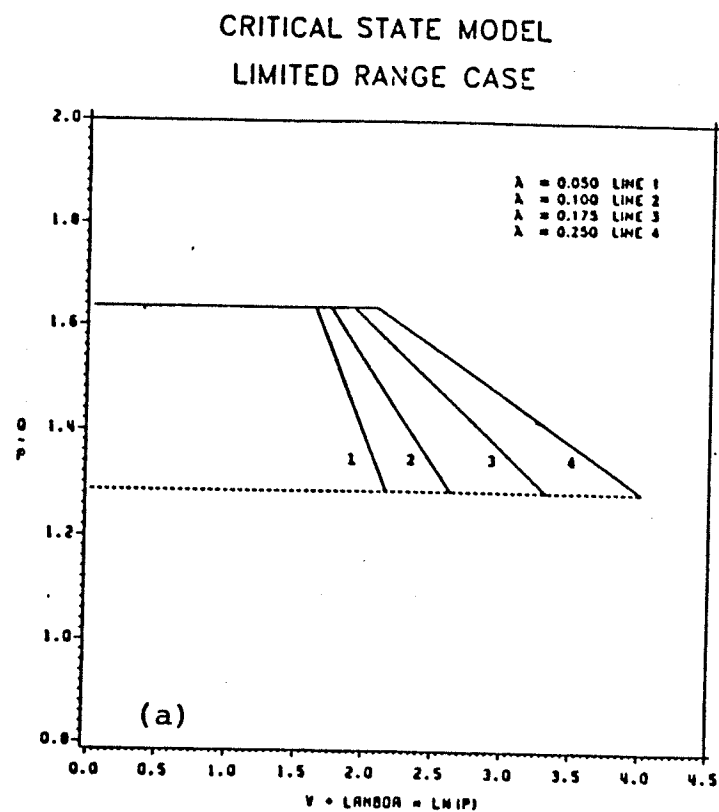


FIG 4.3 - PARAMETRIC DEFINITION OF THE ADOPTED CRITICAL  
CRITICAL STATE MODEL. (a) LIMITED RANGE CASE,  
(b) BROAD RANGE CASE.

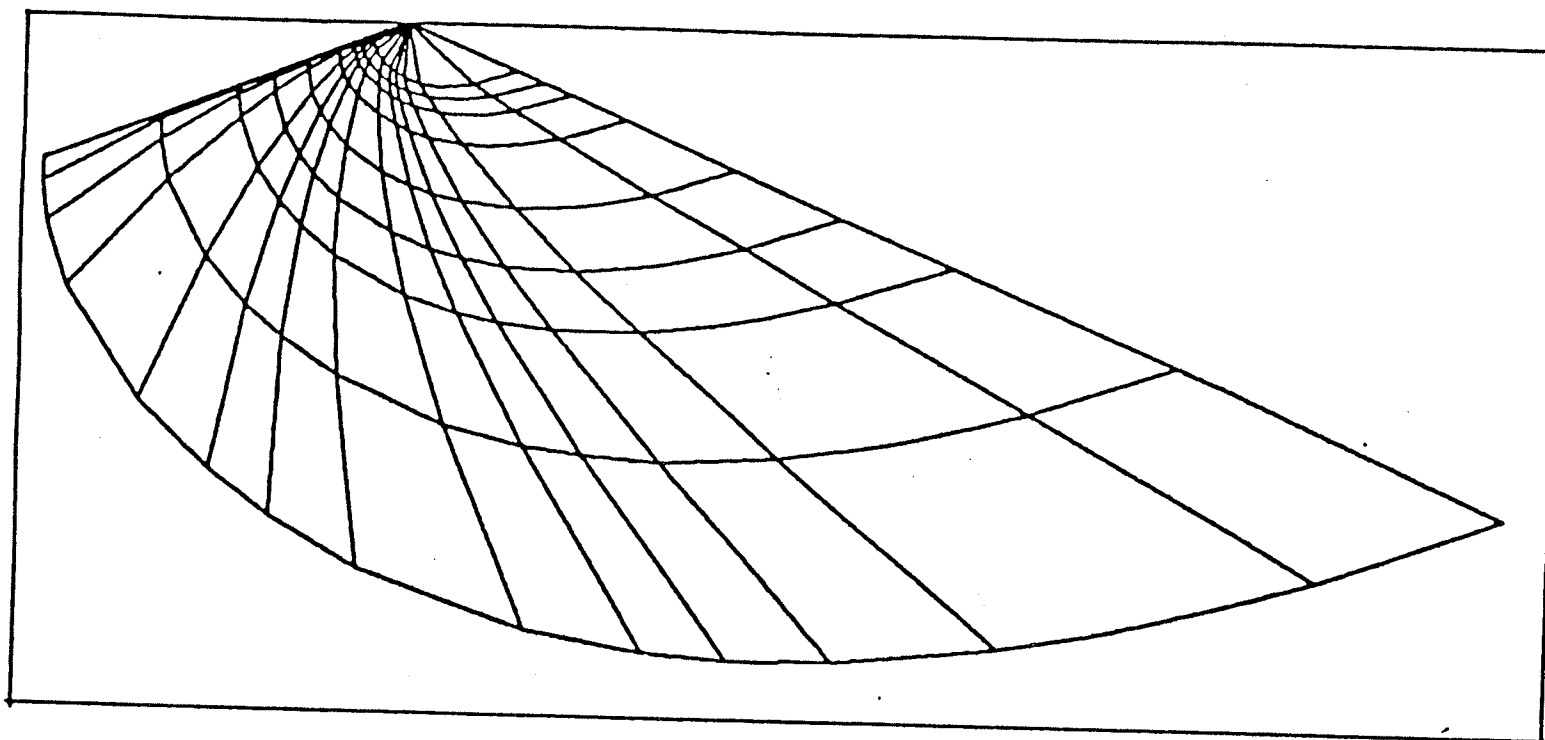


FIG 5.1 - SLIP-LINE FIELD (STRESS-CHARACTERISTICS)  
BENEATH SURFACE FOOTING

## SLIP LINES BEHIND A PASSIVE WALL

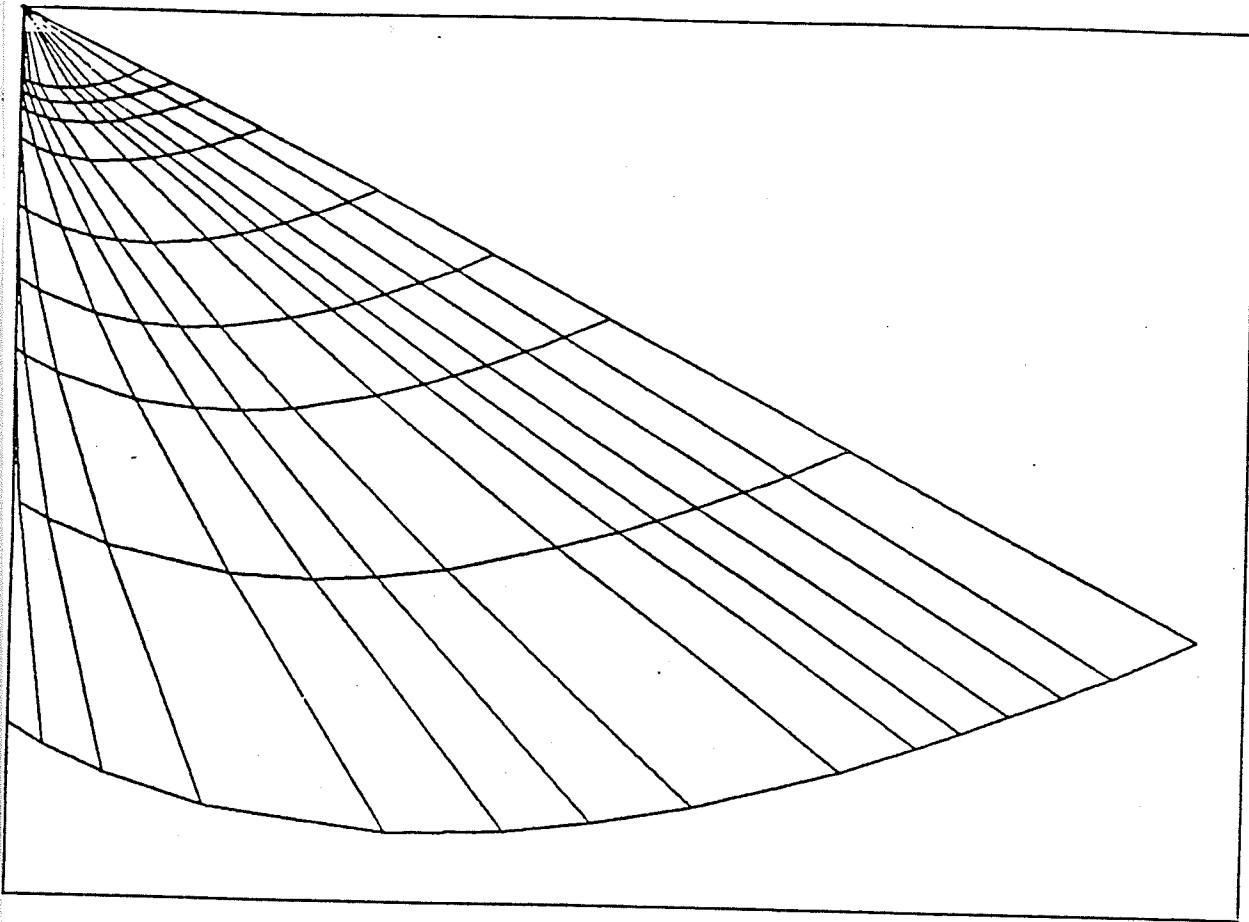
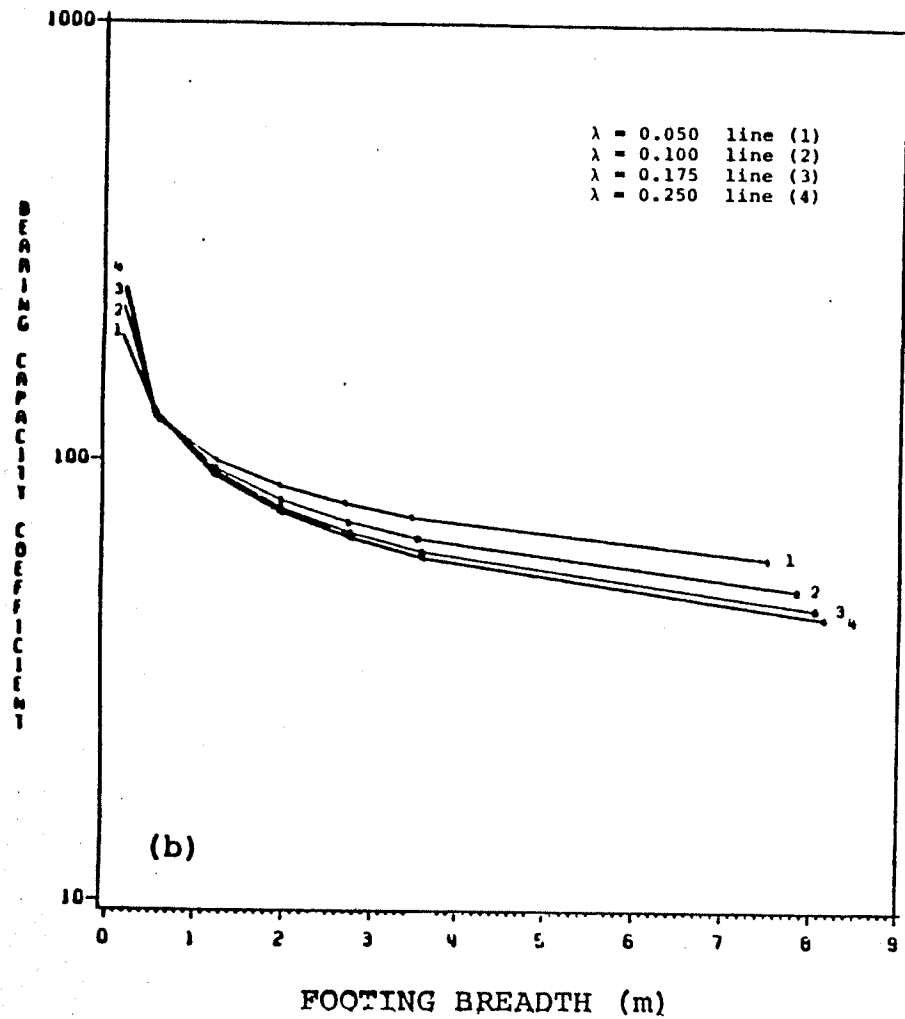


FIG 5.2 - SLIP-LINE FIELD (STRESS-CHARACTERISTICS) BEHIND PASSIVE WALL

# BROAD RANGE CASE



# LIMITED RANGE CASE

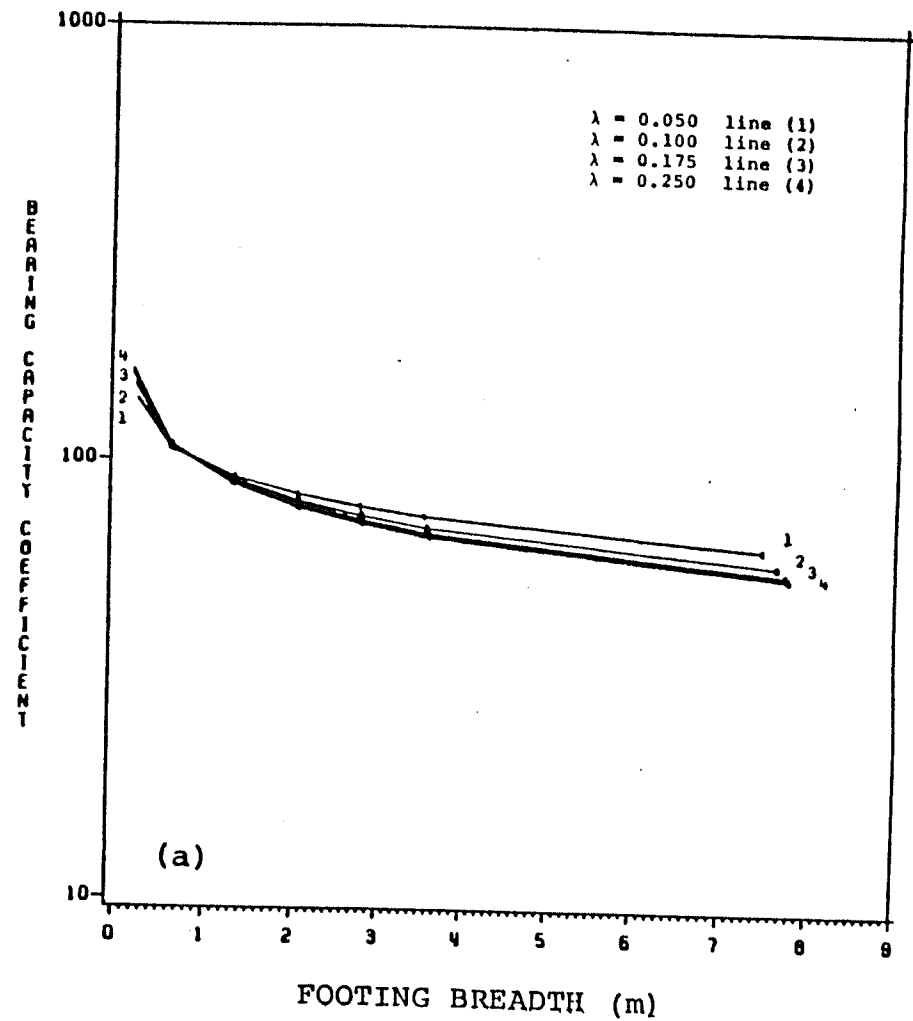
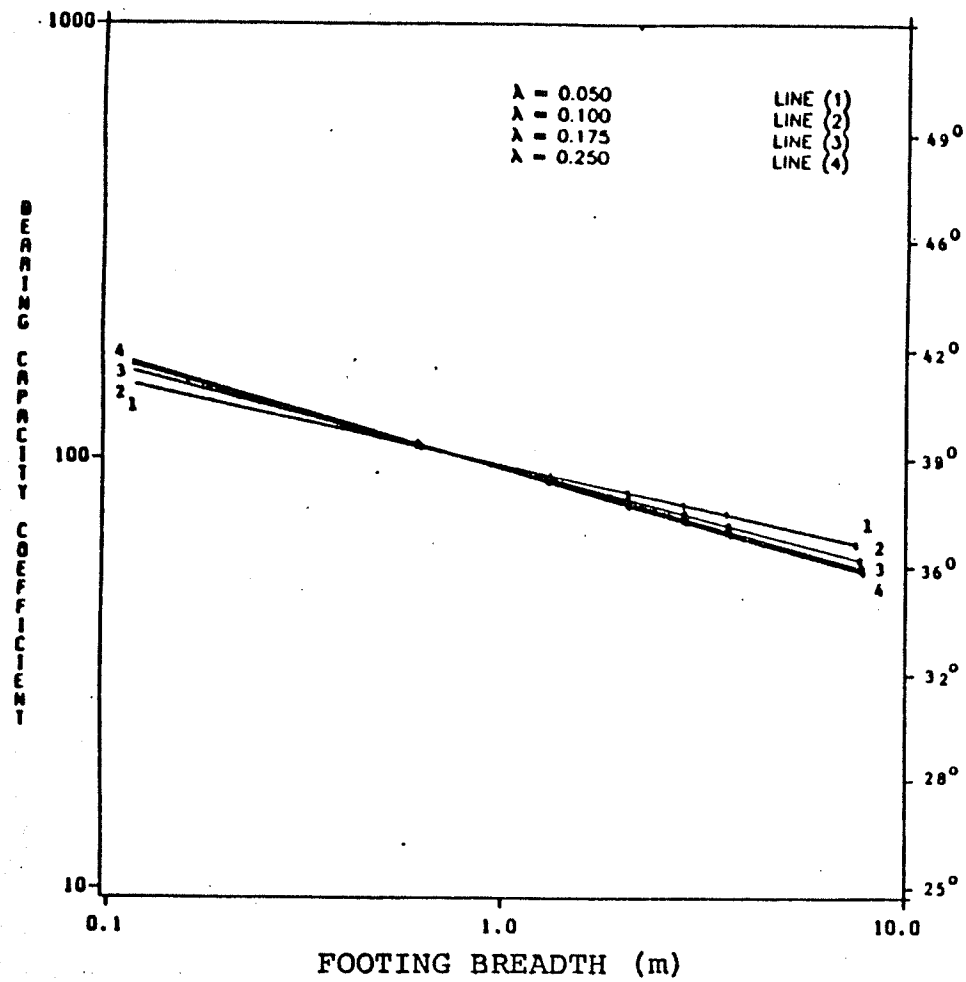


FIG 5.3 — BEARING CAPACITY COEFFICIENTS , (a) LIMITED RANGE, (b) BROAD RANGE (SEMI LOG SCALE)

# LIMITED RANGE CASE



# BROAD RANGE CASE

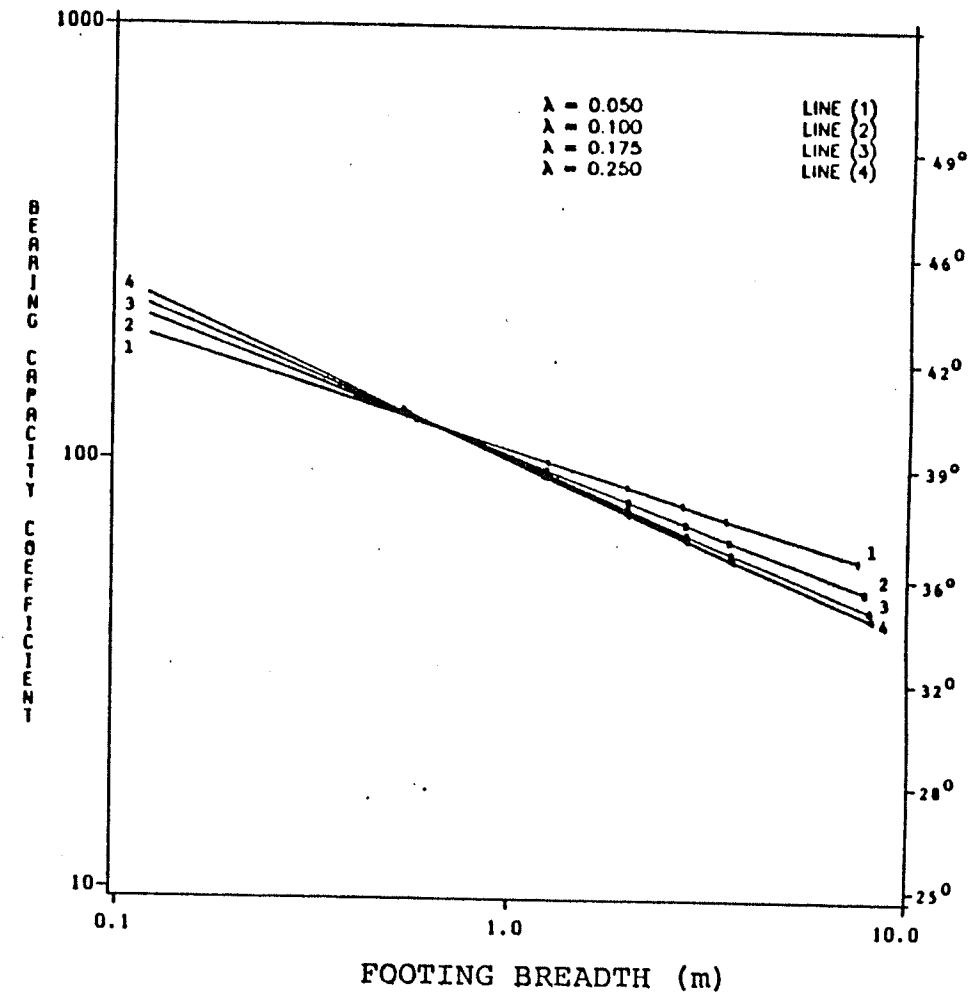


FIG 5.4 — BEARING CAPACITY COEFFICIENTS , (a) LIMITED RANGE, (b) BROAD RANGE (LOG-LOG SCALES)

# BEARING CAPACITY COEFFICIENT LIMITED RANGE CASE

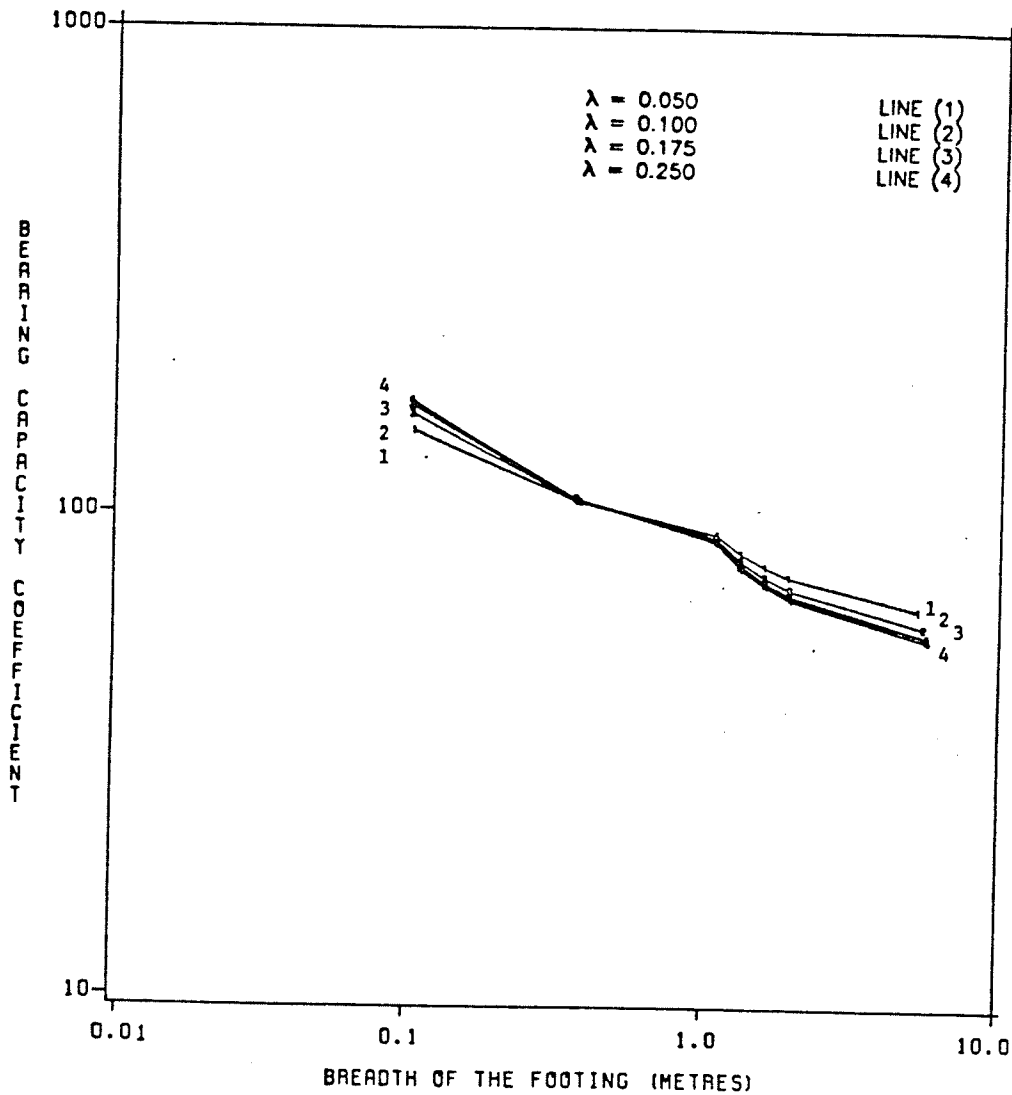


FIG 5.5a - BEARING CAPACITY COEFFICIENTS,  
(a) LIMITED RANGE



BEARING CAPACITY COEFFICIENT  
BROAD RANGE CASE

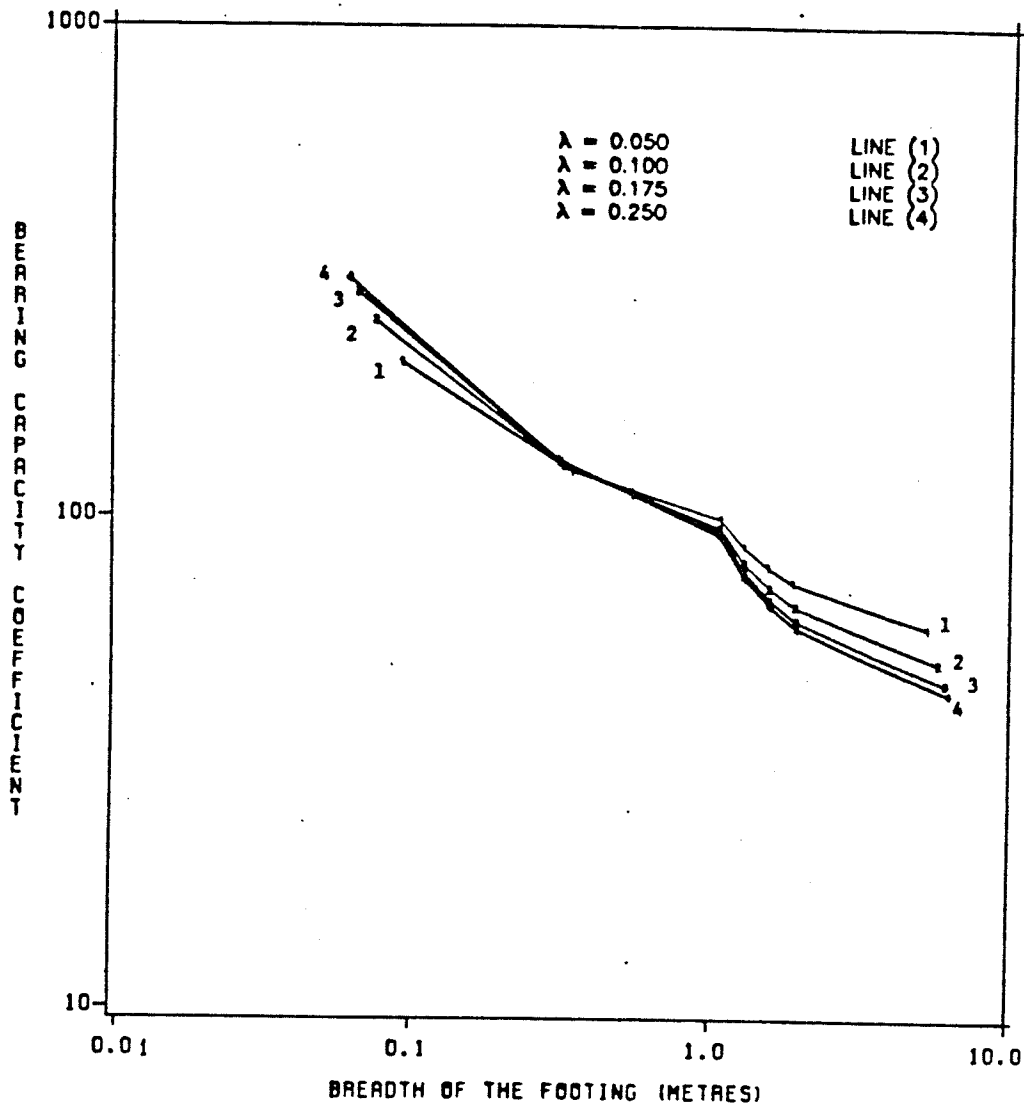
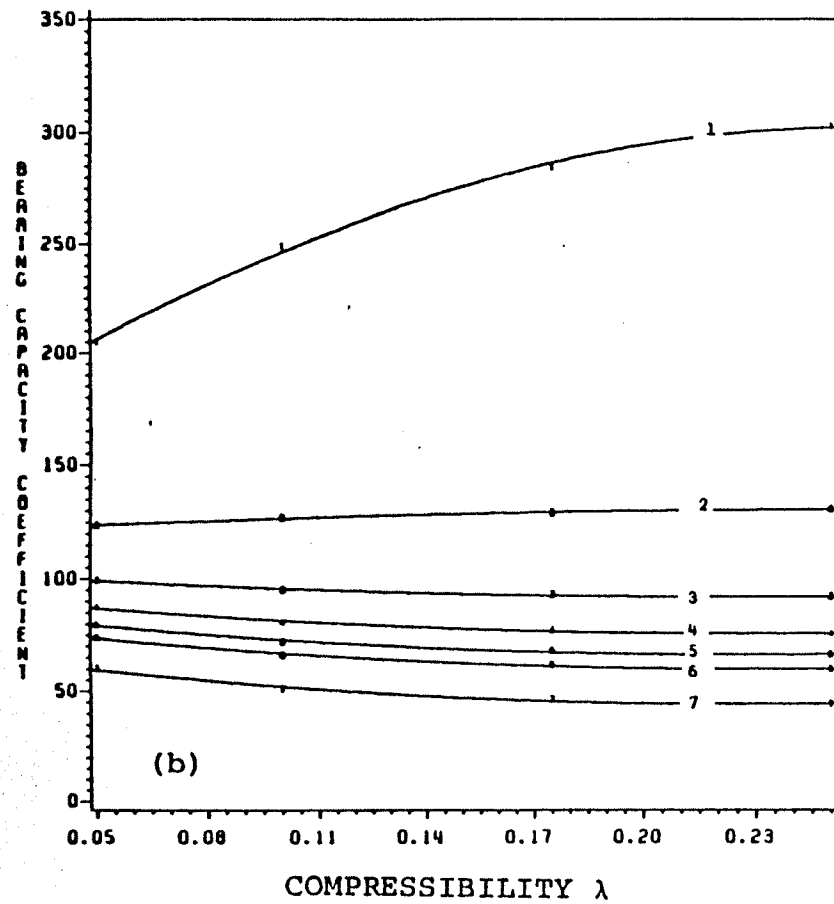


FIG 5.5b - BEARING CAPACITY COEFFICIENTS,  
(b) BROAD RANGE

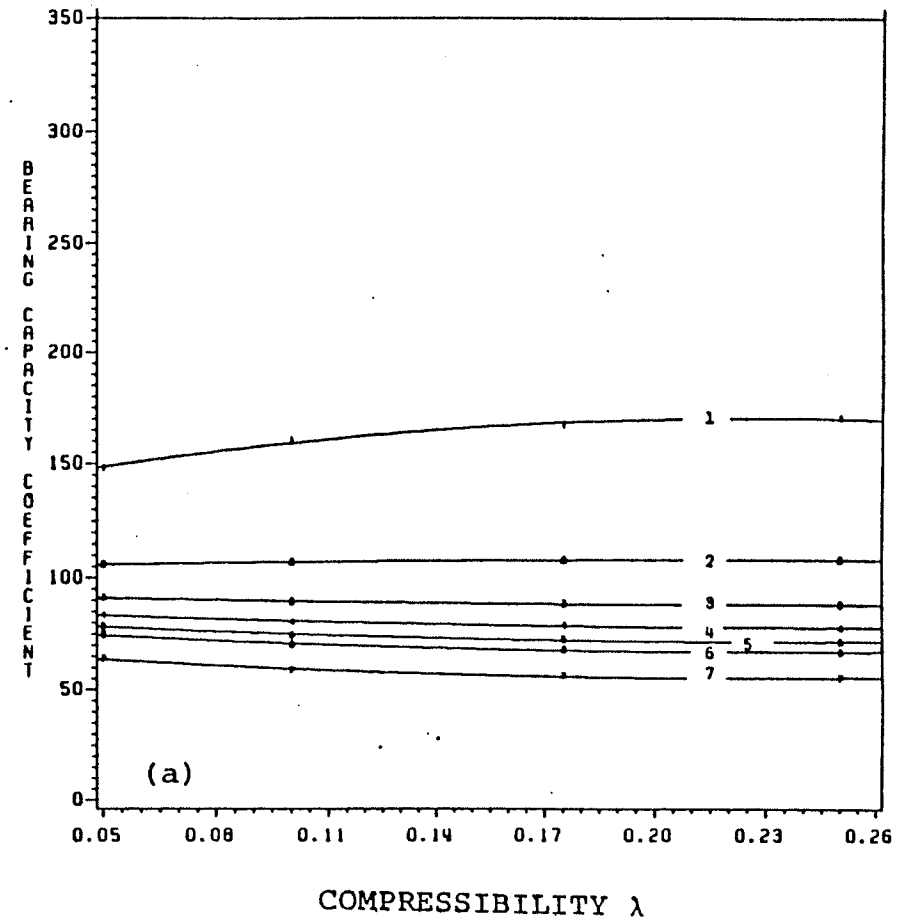
# BEARING CAPACITY COEFFICIENT BROAD RANGE CASE



COMPRESSIBILITY  $\lambda$

Line number : 1 2 3 4 5 6 7  
Footing breadth : 0.10 0.55 1.25 1.95 2.80 3.60 8.10

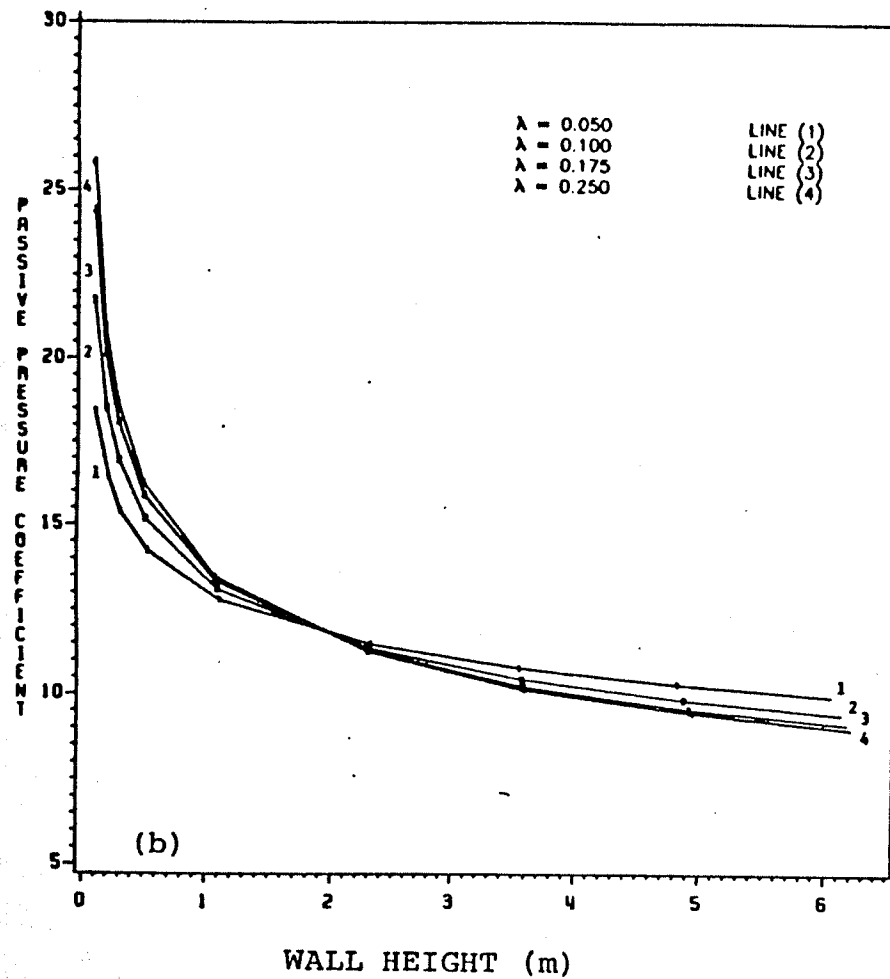
# BEARING CAPACITY COEFFICIENT LIMITED RANGE CASE



COMPRESSIBILITY  $\lambda$

FIG 5.6 - RELATIONSHIPS BETWEEN BEARING CAPACITY COEFFICIENTS AND COMPRESSIBILITY,  
(a) LIMITED RANGE, (b) BROAD RANGE.

# BROAD RANGE CASE



# LIMITED RANGE CASE

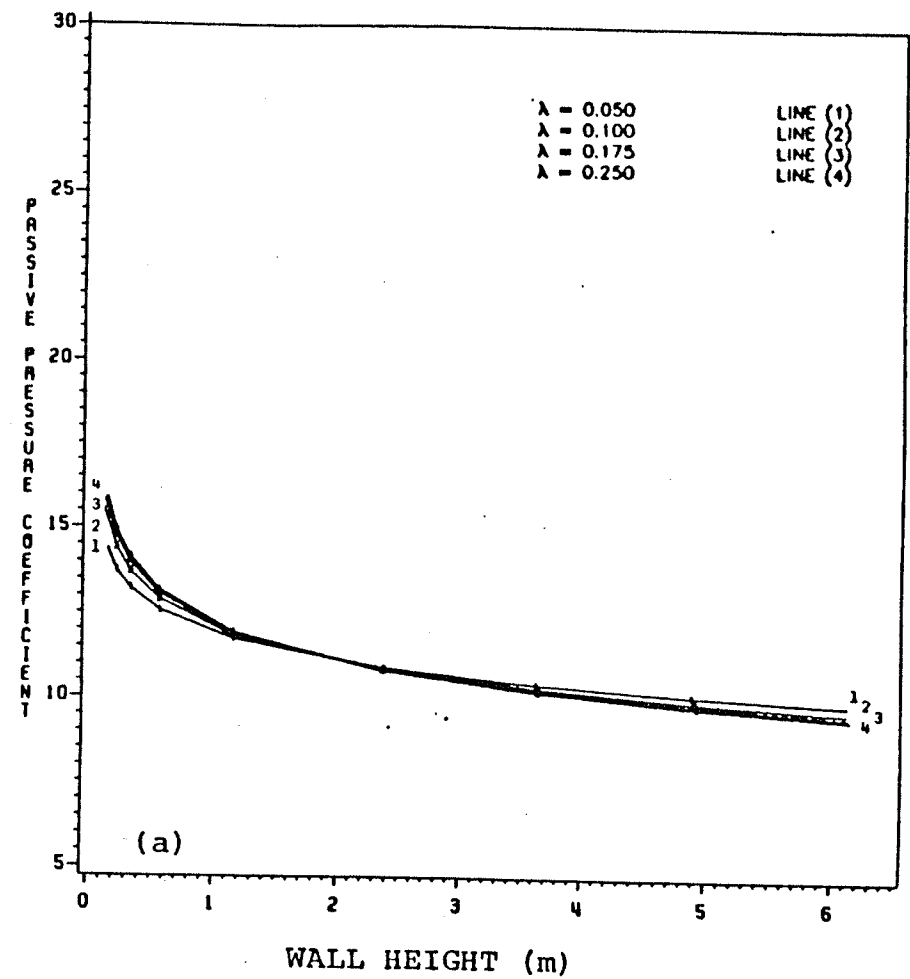
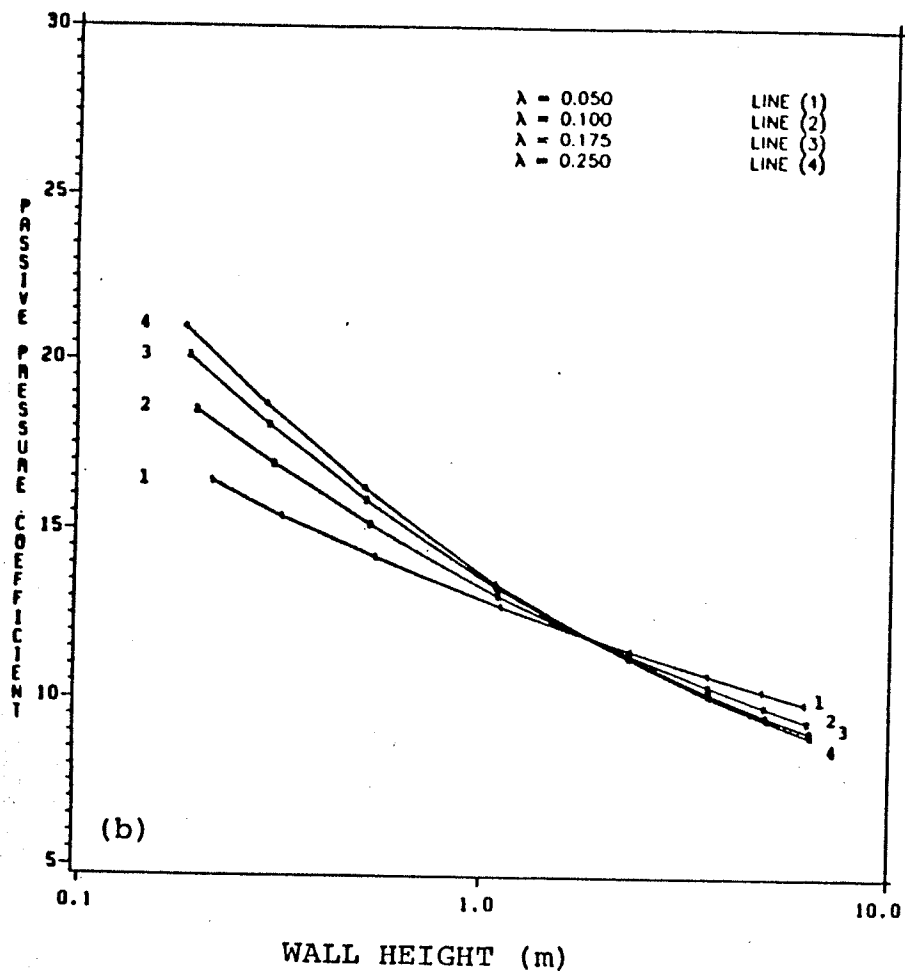


FIG 5.7 - PASSIVE PRESSURE COEFFICIENTS, (a) LIMITED RANGE, (b) BROAD RANGE (SEMI LOG SCALE)

COEFFICIENT OF PASSIVE PRESSURE  
BROAD RANGE CASE



COEFFICIENT OF PASSIVE PRESSURE  
LIMITED RANGE CASE

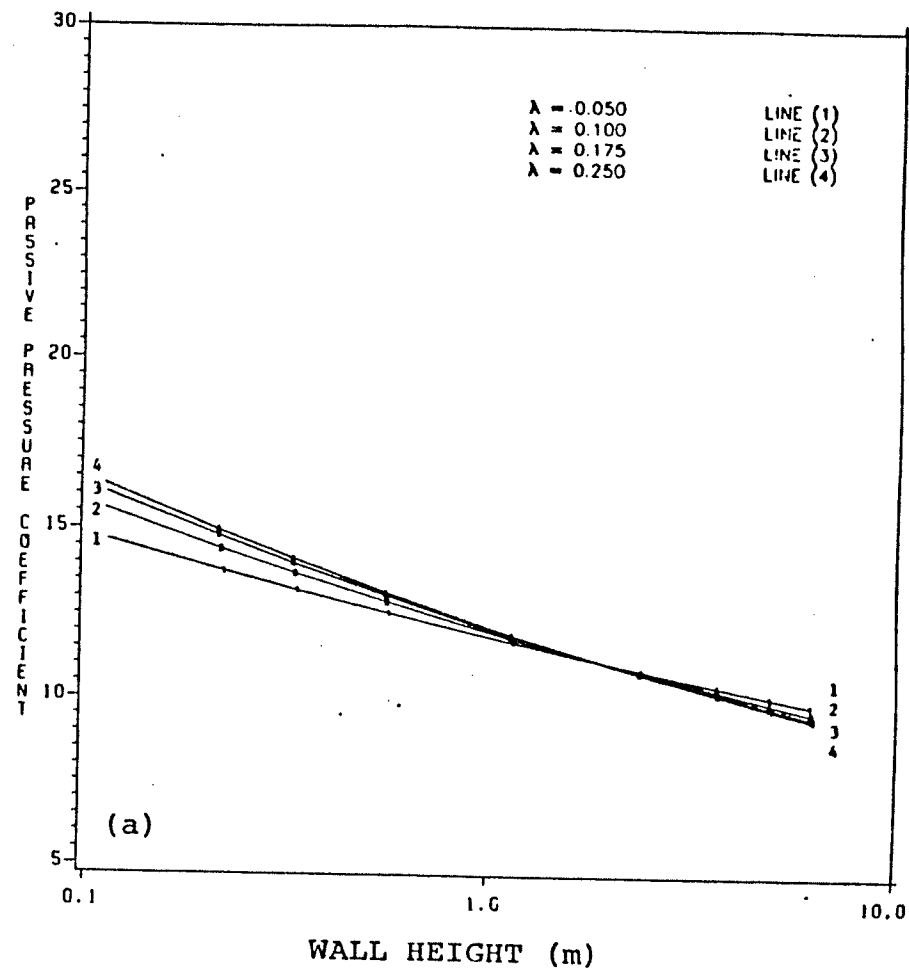


FIG 5.8 - PASSIVE PRESSURE COEFFICIENTS, (a) LIMITED RANGE, (b) BROAD RANGE (LOG-LOG SCALES)

COEFFICIENT OF PASSIVE PRESSURE  
LIMITED RANGE CASE

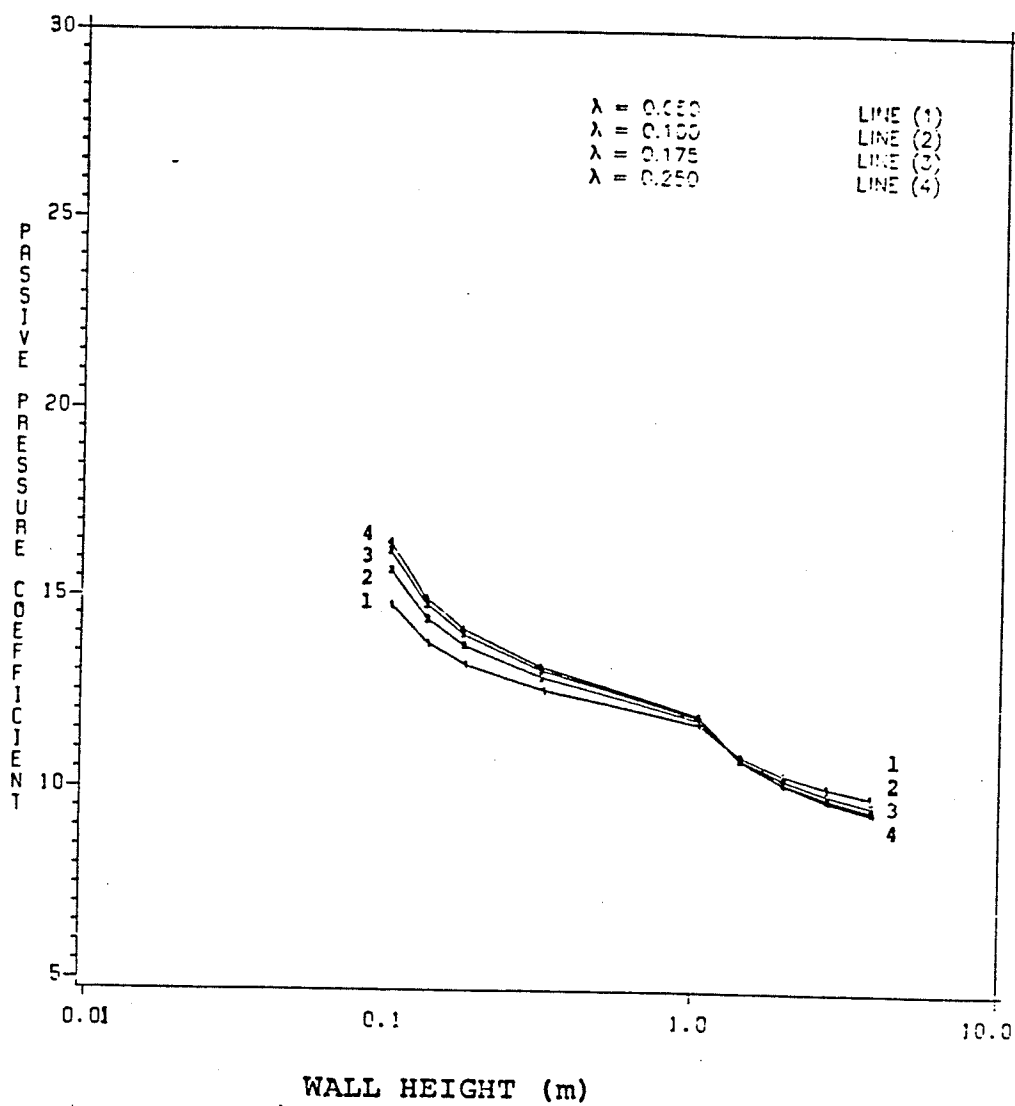


FIG 5.9a - PASSIVE PRESSURE COEFFICIENTS,  
LIMITED RANGE.

COEFFICIENT OF PASSIVE PRESSURE  
BROAD RANGE CASE

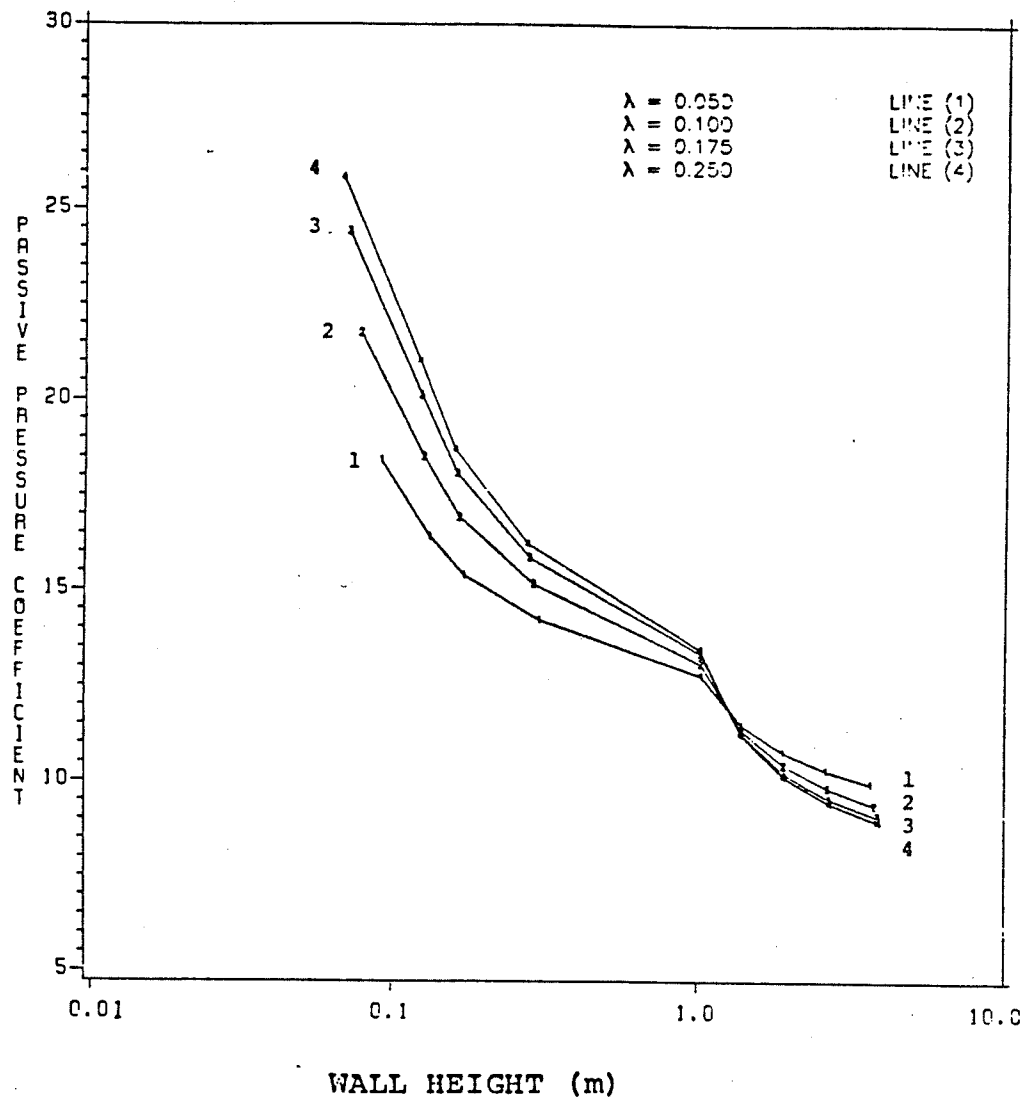
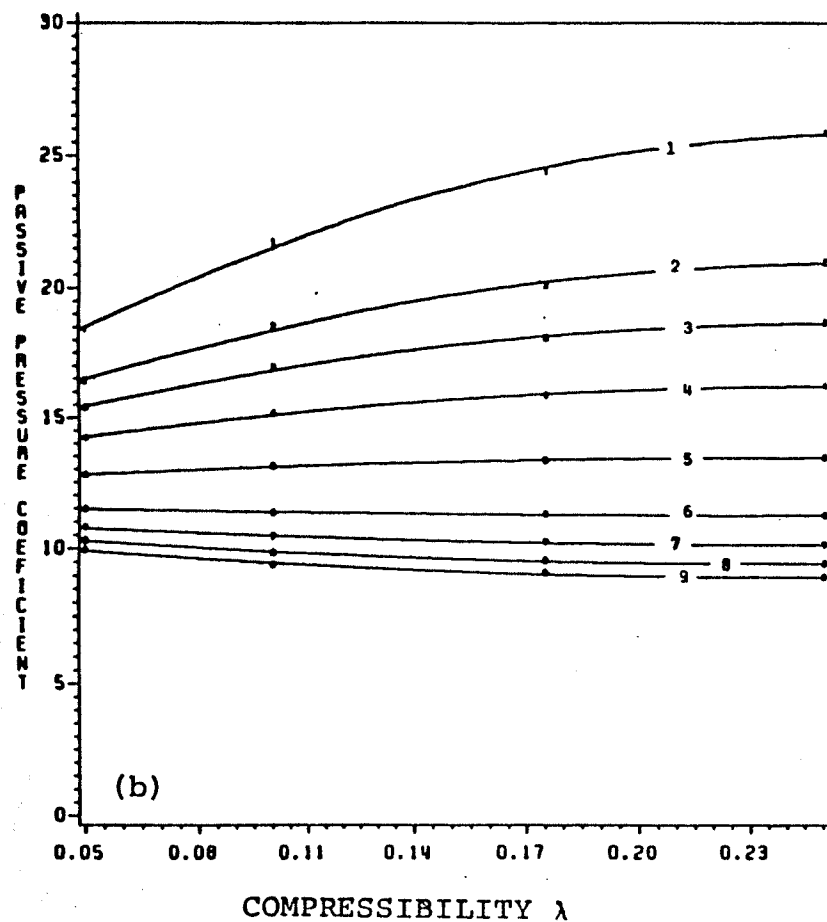
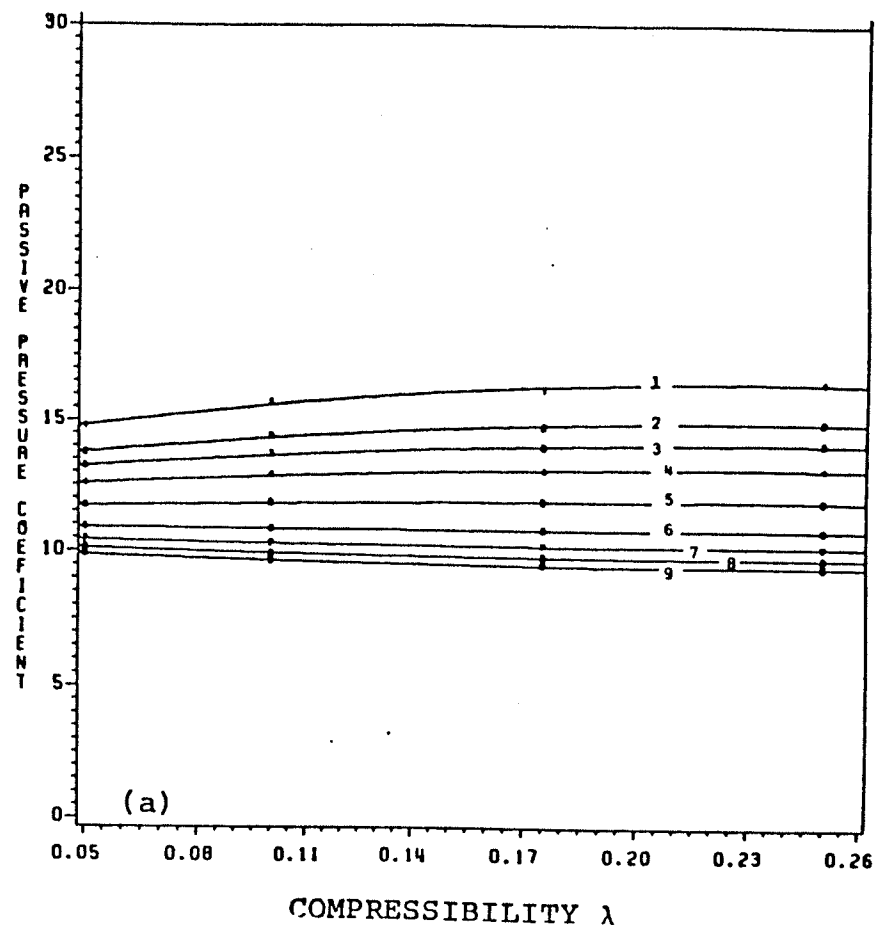


FIG 5.9b - PASSIVE PRESSURE COEFFICIENTS,  
BROAD RANGE.

# COEFFICIENT OF PASSIVE PRESSURE BROAD RANGE CASE

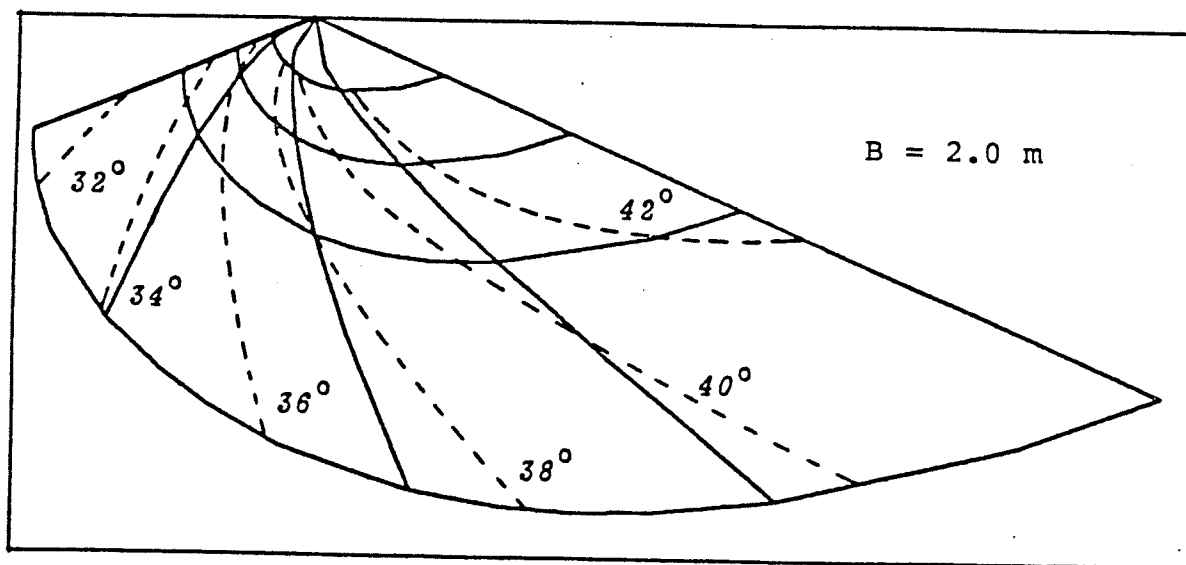
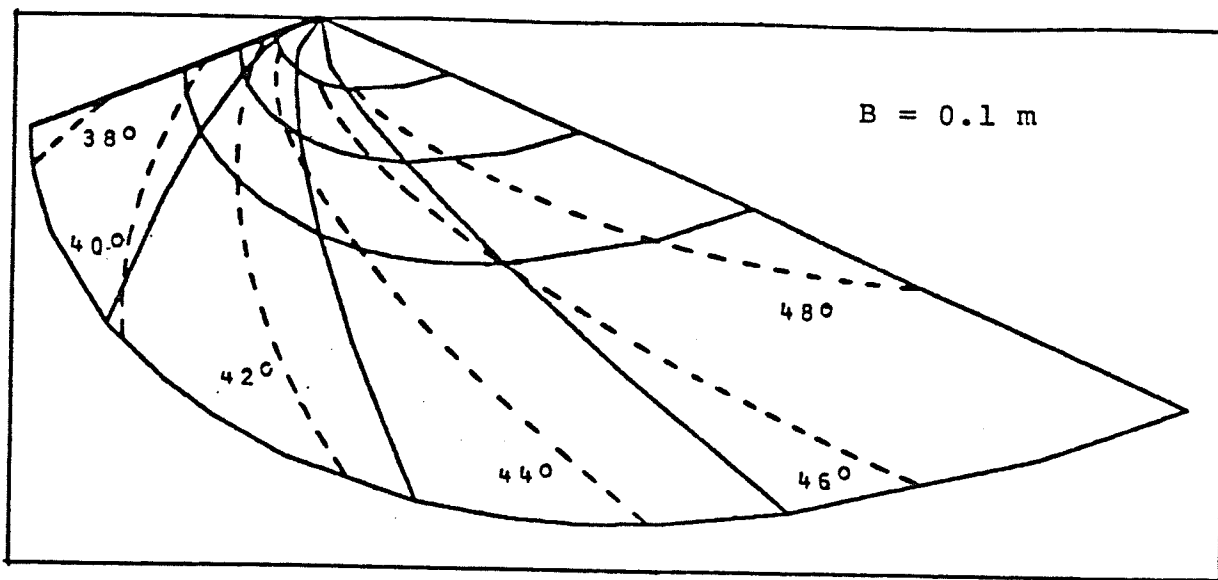


# COEFFICIENT OF PASSIVE PRESSURE LIMITED RANGE CASE



Line number	:	1	2	3	4	5	6	7	8	9
Wall height (m)	:	0.10	0.18	0.28	0.50	1.10	2.30	3.60	4.90	6.40

FIG 5.10 - RELATIONSHIPS BETWEEN PASSIVE PRESSURE COEFFICIENTS AND COMPRESSIBILITY,  
(a) LIMITED RANGE, (b) BROAD RANGE.



6.1 - CONTOURS OF MOBILIZED ANGLE OF SHEARING RESISTANCE (ISOPHIS) BENEATH SURFACE FOOTING



# CONTOURS OF THE ANGLE OF SHEARING RESISTANCE

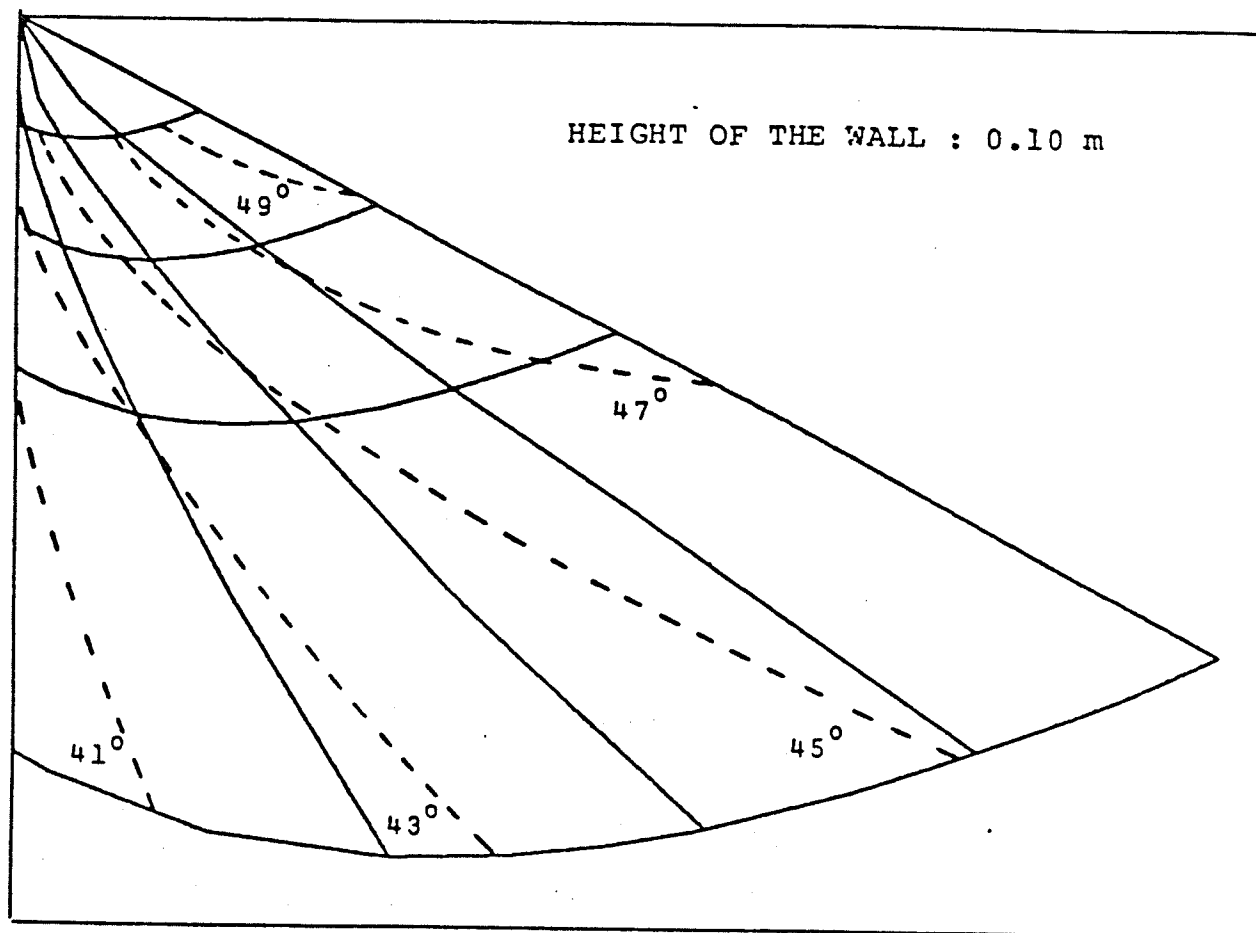


FIG 6.2a : CONTOURS OF MOBILIZED ANGLE OF SHEARING  
RESISTANCE BEHIND PASSIVE WALL

## CONTOURS OF THE ANGLE OF SHEARING RESISTANCE

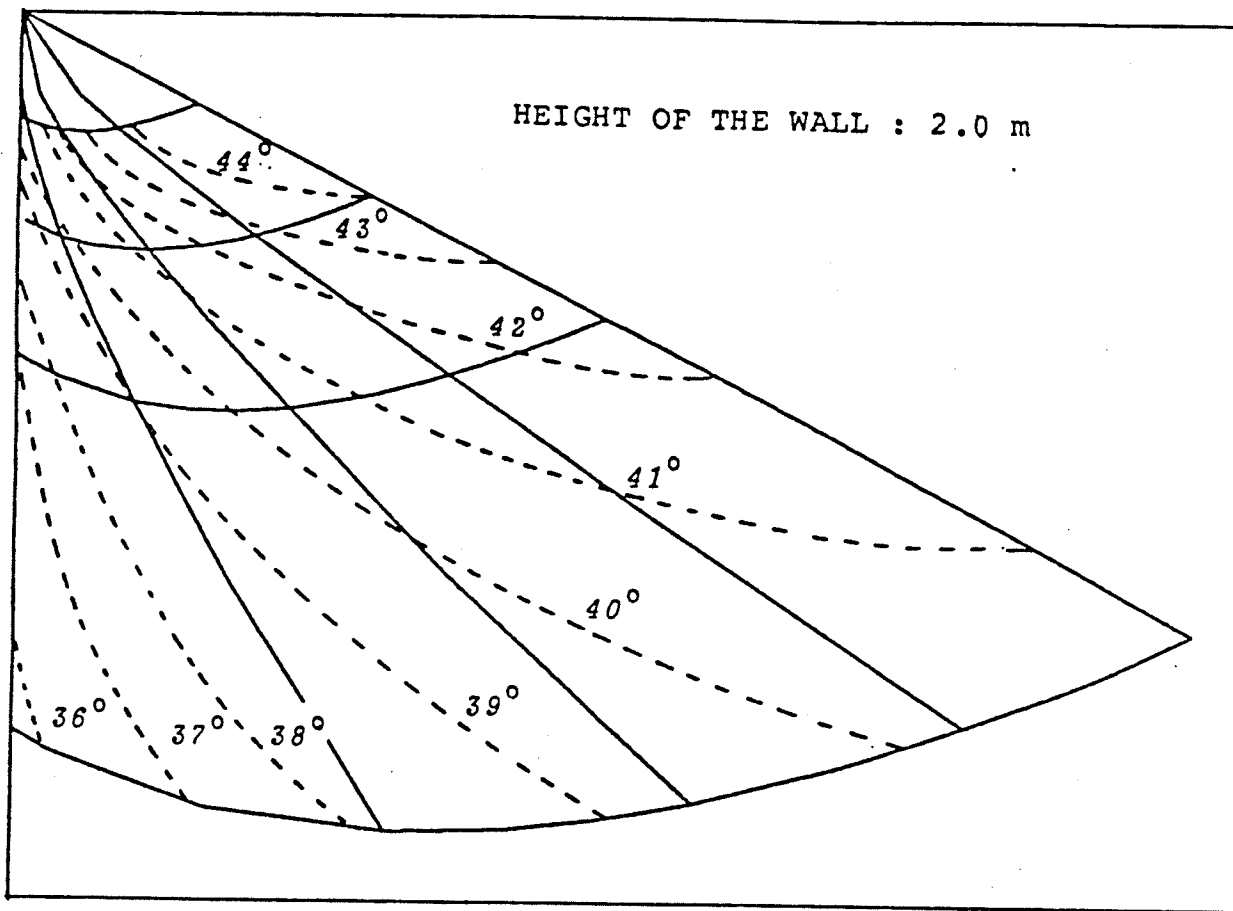


FIG 6.2b : CONTOURS OF MOBILIZED ANGLE OF SHEARING  
RESISTANCE BEHIND PASSIVE WALL

# BEARING CAPACITY COEFFICIENT COMPARISON WITH GRAHAM RESULTS

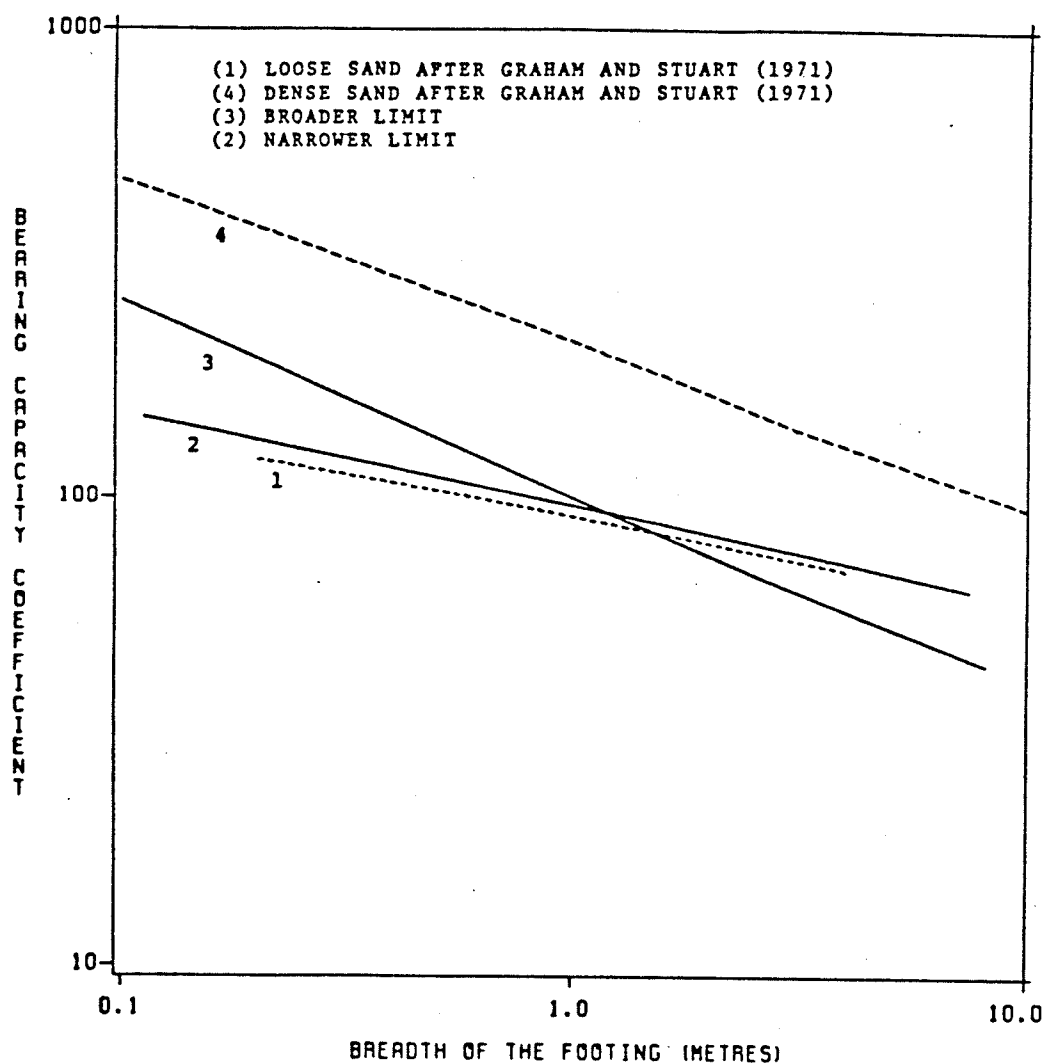


FIG 6.3 - VARIATION OF  $N_\gamma$  WITH FOOTING BREADTH,  
COMPARISON WITH THEORETICAL RESULTS FROM  
GRAHAM AND STUART (1971).

# COEFFICIENT OF PASSIVE PRESSURE COMPARISON WITH GRAHAM RESULTS

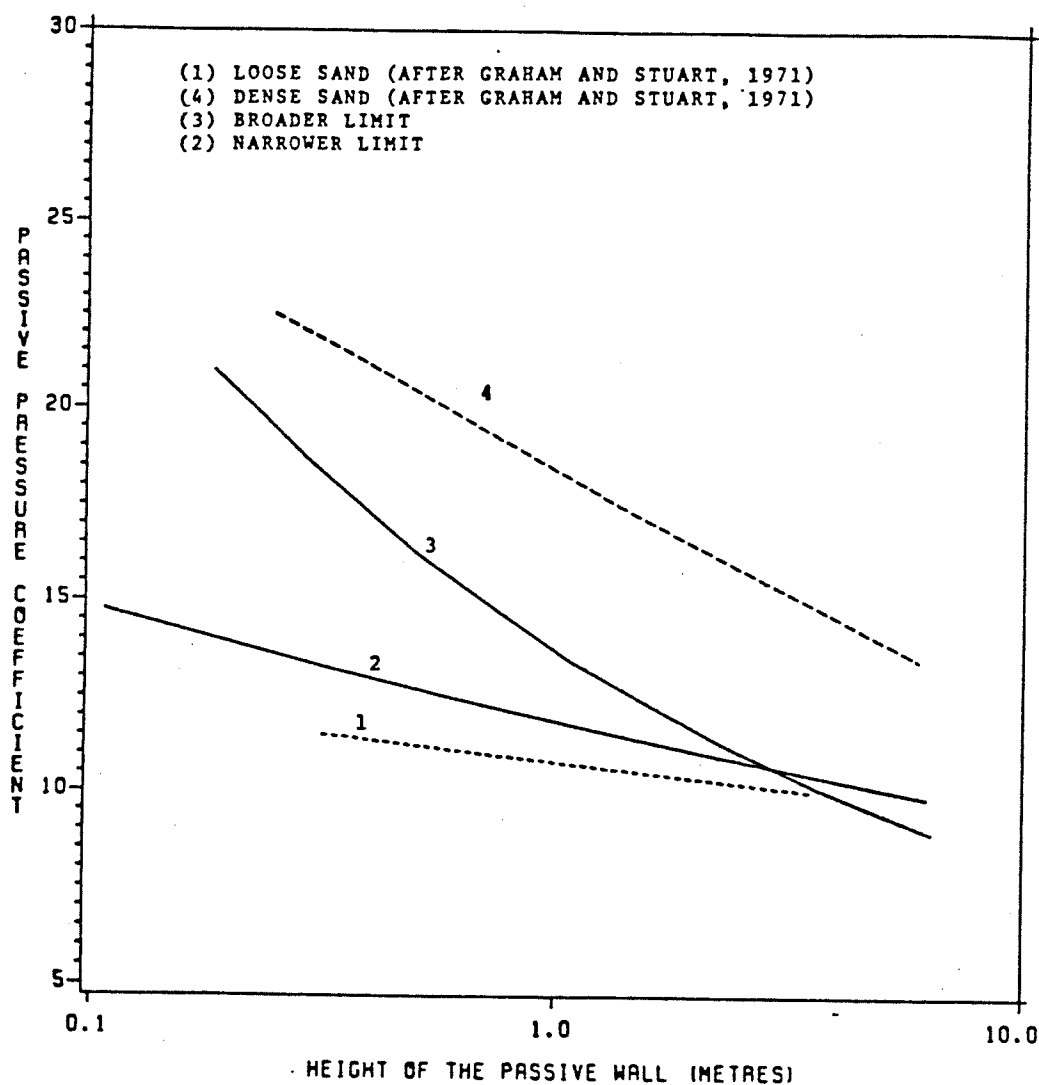


FIG 6.4 - VARIATION OF  $K_p$  WITH WALL HEIGHT,  
COMPARISON WITH THEORETICAL RESULTS  
FROM GRAHAM AND STUART (1971).

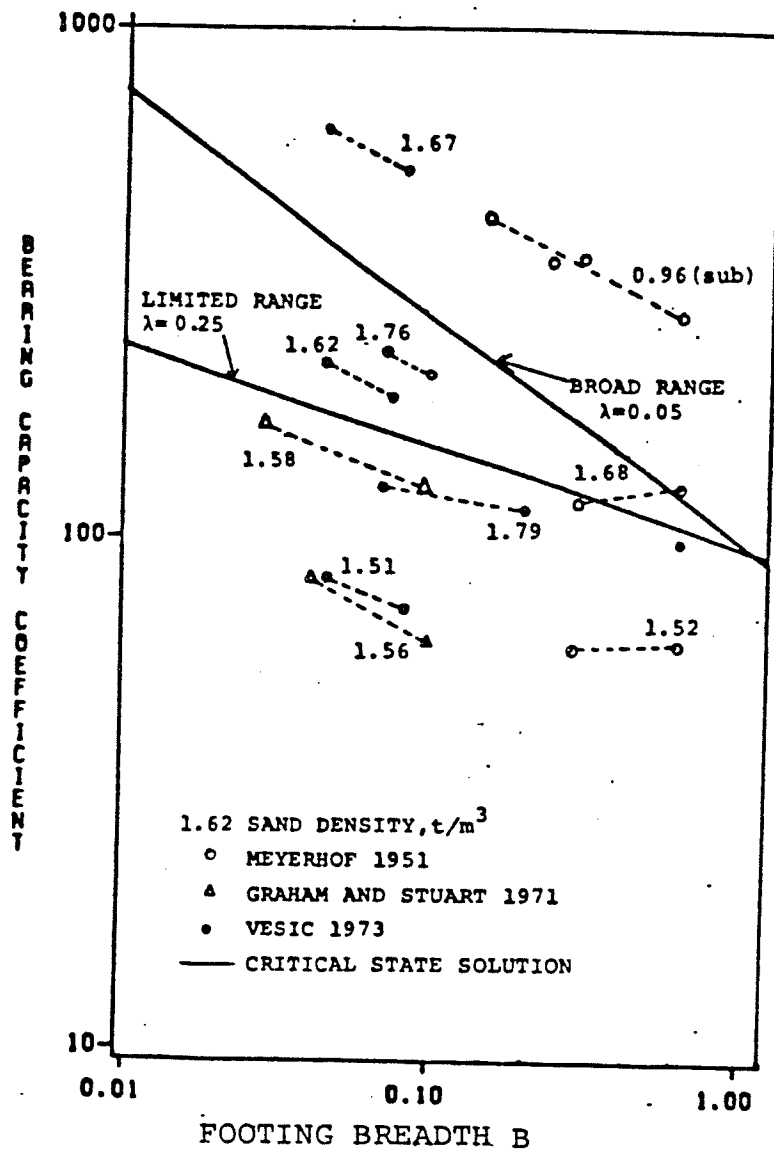
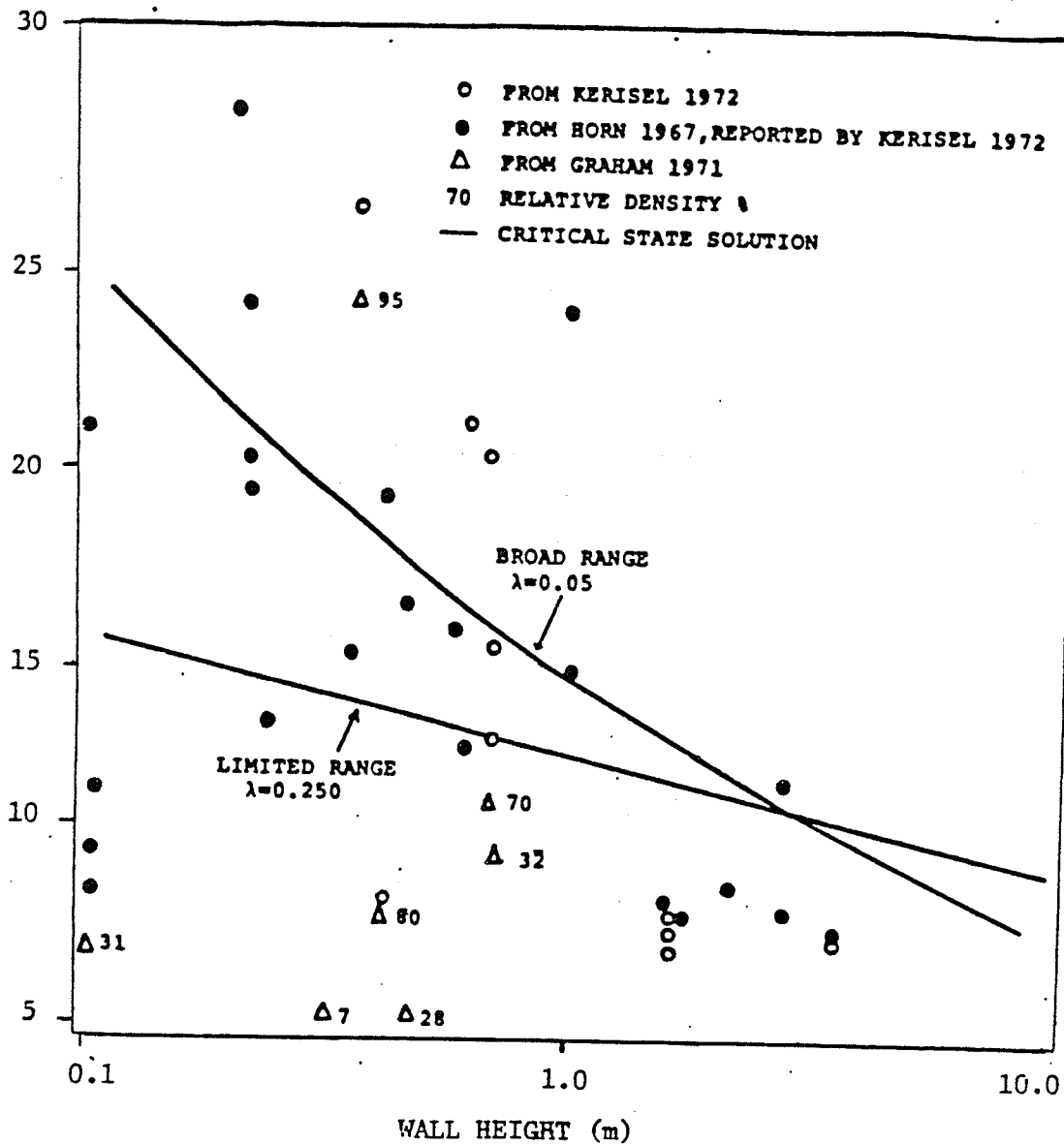


FIG 6.5 : VARIATION OF N WITH FOOTING BREADTH B,  
COMPARISON WITH EXPERIMENTAL RESULTS.



3 6.6 — VARIATION OF  $K_p$  WITH WALL HEIGHT  $H$ , COMPARISON WITH EXPERIMENTAL RESULTS.

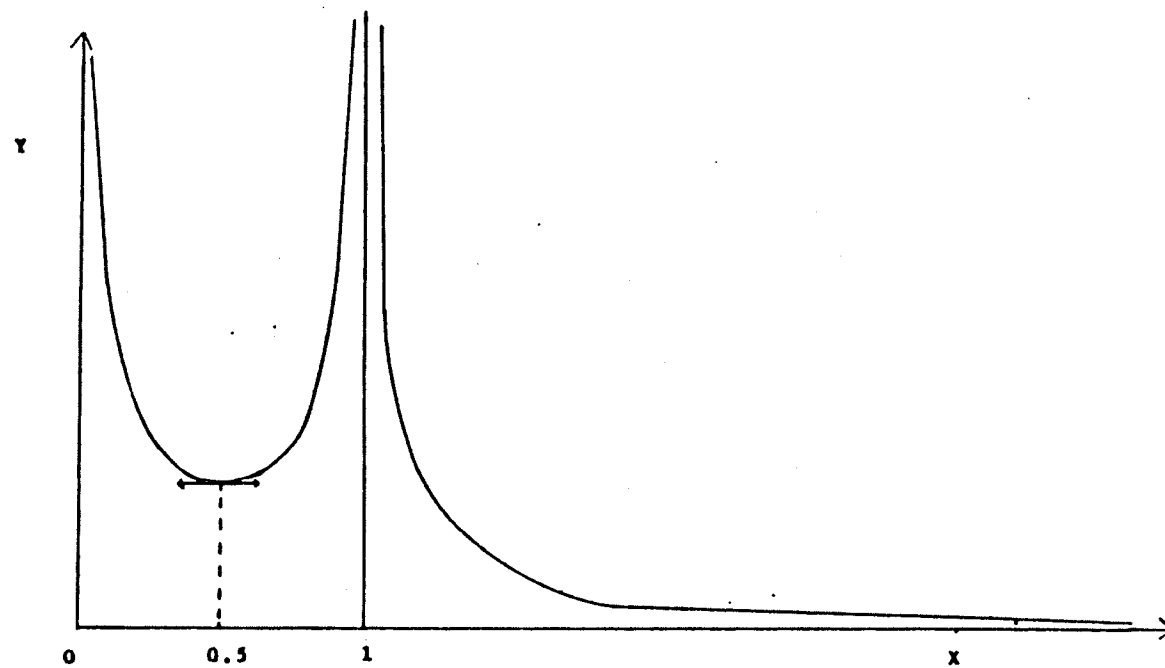


FIG 6.7 - PLOT OF THE FUNCTION  $Y=1/|X(X-1)|$

## Appendix A

### FACTORS INFLUENCING THE CHOICE OF PARAMETERS FOR COMPUTATION

Bearing in mind that the computer program has been developed for a sand failing in a rigid-plastic mode, it is relatively easy to make use of the program to analyse practical applications. This appendix summarizes the necessary input parameters and proposes guidelines for their estimation.

#### A.1 NECESSARY INPUT PARAMETERS

The input is fed into the computer in the following order :

1.  $e_{\min}, e_{\max}$  : minimum, maximum void ratio
2.  $p'_{\min}, p'_{\max}$  : minimum, maximum pressure
3.  $\phi'_{\min}, \phi'_{\max}$  : minimum, maximum angle of shearing resistance
4.  $\lambda$  : compressibility ( $\lambda = C_c/2.3$ )
5.  $l$  : scale parameter
6.  $n$  : number of spiral lines
7.  $m$  : number of radial lines
8.  $x(n)$  : abscissa of the intersections of the spirals and of the edge of the passive zone



## A.2 PRACTICAL ESTIMATION

### A.2.1 Void ratio

Table A.1 indicates for various soils the :

1. maximum void ratio
2. minimum void ratio

The initial condition of a sand can be described by its relative density  $I_D = (e_{\max} - e) / (e_{\max} - e_{\min})$ . The medium dense range of behaviour can be taken to be cases where the  $I_D$  is greater than 35 and less than 85 (Lambe and Whitman, 1969).

### A.2.2 Pressures

Minimum as well as maximum pressures can be estimated as the unloaded state of the sand and from the anticipated maximum stress intensity under the footing or behind a passive wall. Typical values are respectively 10 kPa and  $10^4$  kPa.

### A.2.3 Angle of shearing resistance

Usually, the engineer who calculates the size of a footing or that of a wall has also the task of estimating the angle of shearing resistance of the soil. The choice lies in the results of laboratory tests combined with experience. Minimum and maximum angles of shearing resistance can be taken as the extreme values that the engineer considers can be mobilized in the sand by the pressure levels in the problem. The values are affected by the density, mineralogy, and by the particle size distribution, particle shape of the sand.

#### A.2.4 Compressibility $\lambda$

Since  $\lambda = C_c/2.304$ , the determination of  $\lambda$  is related to that of the compression index. Compression tests into the normally consolidated (linear  $V, \log p'$ ) range can supply such information. Such tests are rare. In the absence of reliable data, a typical value of  $\lambda$  is 0.1 (Atkinson and Bransby, 1978).

#### A.2.5 Scale parameter

Note that this is not equal to the height of the wall or the footing breadth and the structure size results from this choice. Any figure can be taken as scale parameter. Confusion might appear for the user who is not familiar with the computer program. The choice of the horizontal length of the edge of the passive zone is recommended. This figure corresponds to the abscissa of the intersection on the extreme spiral and of the edge of the passive zone (Figs 4.1 and 4.2).

#### A.2.6 Radial and spiral lines

Augmenting the number of spiral and radials lines leads to greater accuracy. Nonetheless, the user should keep it mind that it will also increase the computing time. Ten spirals and twenty radials in the present work have been used.

### A.2.7 Edge of the passive zone

If ten spirals are chosen, the horizontal length of the edge of the passive zone is divided in ten equal segments. A spiral line will depart from each of these ten points. The abscissa  $x(n)$  are the required input parameters. The dimensionless figures  $n\pi/10$  are convenient ( $n = 1, 10$ ).

Appendix B  
LISTING OF THE COMPUTER PROGRAM

PHI -  $\phi$   
FIRPHI - First  $\phi$   
PHILO - Local  $\phi$   
XX,ZZ - Physical coordinates  
SIMA, SIGMA1, SIGMA2 -  $\sigma$   
PSSI, PSI1, PSI2 -  $\psi$   
P - Pressure  
E - voids ratio  
L - Number of spirals  
K - Number of radials  
C - Control  
LANDA -  $\lambda$   
GS - Specific gravity  
GW - Water density  
PP, QQ - Iteration counts

SUBSCRIPTS :

MIN - Minimum  
MAX - Maximum

## COMPUTER PROGRAM

```

REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
INTEGER L, K, C
COMMON L, K, C
REAL XX, ZZ, SIMA, PSS1, PHH1
COMMON /TATA/ X1, X2, XX, Z1, Z2, ZZ, PHH1
COMMON /TITI/ SIGMA1, SIGMA2, SIMA
COMMON /TETE/ PS11, PS12, PSS1
COMMON /TOTO/ P1, SNPHI, CSPHI
INTEGER PP, QQ, CONTRO
COMMON /OUI/ PP, QQ, CONTRO
REAL MU
COMMON /AH/ MU, PHILO, SIGLOC, PHI1, PHI2, PHILOP, PHILOT
REAL
X(40,40), Z(40,40), SIGMA(40,40), PSI(40,40), PHI(40,40)
COMMON /LALA/ X, Z, SIGMA, PSI
COMMON /ALO/ PHI
INTEGER P(40), Q(40)
COMMON /LOLO/ P, Q
REAL EO, LANDA
COMMON
EO, LANDA, PMIN, PMAX, PHMIN, PHIMAX, PHI1, EMIN, EMAX /CRIT/
INTEGER SPI, RAD, W, H, H1, KS, L1, LK, I, J, RAD1
REAL DIF, DI, TNPHI
REAL KP, NY, FIRSIG, GS, GW
C
C
C
PI = ARCOS(-1.)
READ, FIRPHI, DELTA, BETA, TOL, CONVG, SIZE, GS, GW
READ, CONTRO
IF (CONTRO .EQ. 1) THEN DO
PRINT 43
43
FORMAT('1', '//////////20X, '*****',
1 /20X, '*'
2 /20X, '*' CONTROL= 1
3 /20X, '*'
4 /20X, '*' PARTIAL RESULTS PRINTOUT
5 /20X, '*'

```

- i -

```

/20X, '*****')
ELSE
GO TO 44
ENDIF
44 CONTINUE
READ, L, K
READ, EMIN, EMAX, LANDA, PMIN, PMAX, PHMIN, PHIMAX
EO = (EMIN + EMAX) / 2.
GAMMA = (GW * GS) / (1. + EO)
PHMIN = PHMIN * PI / 180.
PHIMAX = PHIMAX * PI / 180.
IF (ABS(BETA) .LE. 0.0001) BETA = 0.0001
C
INITIALISATION
CALL HEADIN
PRINT 825
825 FORMAT('1', '//////////20X)
PRINT, '
*****
PRINT, '
* THE PROGRAM COMPUTES THE FAILURE
LINES OF
PRINT, '
* A SAND HAVING A PLASTIC BEHAVIOUR.
THIS IS
PRINT, '
* THE CASE OF A WALL (ROUGH OR SMOOTH),
OR
PRINT, '
* A FOOTING. THE COMPUTATION PRESENTS A
FEATURE
PRINT, '
* WHICH CONSISTS OF TRAPPING THE
FAILURE LINES
PRINT, '
* AROUND A POLE BY IMPLEMENTING VARIOUS
SCALE
PRINT, '
* REDUCTIONS WITH A DIMINISHING
SURCHARGE EFFECT.
PRINT, '
PRINT, '
*****
L1=40
LK=40
DO 1 I=1, L1

```

- ii -

```

DO 2 J=1,LK
  X(1,J)=0.
  Z(1,J)=0.
  SIGMA(1,J)=0.
  PSI(1,J)=0.
  PHI(1,J)=0.
2 CONTINUE
1 CONTINUE C C
DO 3 I=1,L1
  P(I)=0
  Q(I)=0
3 CONTINUE
  BETA= BETA * PI/180.
  DELTA= DELTA * PI/180.
  FIRPHI=FIRPHI/57.2957795
  CSPHI= COS(FIRPHI)
  SNPHI= SIN(FIRPHI)
  TNPHI= TAN(FIRPHI)
  X(1,1)=0.0
  Z(1,1)=0.0
21 CONTINUE
  FIRSIG = 1.0/(1 - SIN(FIRPHI))
  GAMMA=GAMMA * 6.369
  SIZE= SIZE * 3.2808
  SIGLO C= FIRSIG * SIZE * GAMMA
  GAMMA=GAMMA/6.369
  SIZE=SIZE/3.2808
  CALL PHILOC
  DI=(FIRPHI-PHILO)/PHILO
  IF(ABS(DI) .GT. 0.001) THEN
    FIRPHI=PHILO
    GO TO 21
  ELSE
    GO TO 24
  ENDIF
24 CONTINUE
  SIGMA(1,1)=SIGLOC/(SIZE*GAMMA*3.2808*6.369)
  PSI(1,1)=0.5*3.1415927
  PHI(1,1)=PHILO
  PHI(1,1)=PHI(1,1)*180./PI
  PRINT,PHI(1,1)
  PHI(1,1)=PHI(1,1)*PI/180.
  DO 5 I=2,L
    READ, X(1,1)
5 CONTINUE
  C=0
  CALL EDGEPA
  PRINT 12
12 FORMAT('1',////////12X,' EDGE OF RECTILINEAR PASSIVE
ZONE')
  PRINT 19
19 FORMAT(/17X,'X',13X,'Z',12X,'SIGMA',12X,'PSI',
1 15X,'P',4X,'Q',8X,'PHI'/)
  DO 15 I=1,L

```

- iii -

```

CALL RESOUT(1,1)
15 CONTINUE
  C=1
  CALL POLCOM
  SPI=L-1
  C FIRST TIME COMPUTATION
  C=0
450 C=C+1
  PRINT 22,C
22 FORMAT('1',////////25X,'*****',
1 1 25X,'*
2 1 25X,'* SCALE ITERATION C=',13,
3 1 25X,'*
4 1 25X,'* REGULAR DOMAIN
5 1 25X,'*
6 1 25X,'*****')
  IF( C.EQ.1 ) THEN DO
    RAD= K+1
    W=0
  ELSE DO
    RAD= K+ SPI
    W= SPI-1
  ENDIF
  DO 100 J=1,RAD
    IF(CONTRO .EQ. 1) GO TO 53
    PRINT ,
    PRINT ,
    PRINT ,
    PRINT 37,J
    PRINT 39
39 FORMAT(10X,' *****')
37 FORMAT(10X,' J=',13)
    PRINT 19
53 CONTINUE
    IF(J .EQ. 1) GO TO 50
    DO 75 I= 3,L
      X1= X(I-1,J)
      Z1= Z(I-1,J)
      SIGMA1= SIGMA(I-1,J)
      PSI1= PSI(I-1,J)
      X2= X(I,J-1)
      Z2= Z(I,J-1)
      SIGMA2= SIGMA(I,J-1)
      PSI2= PSI(I,J-1)
      PHI1=PHI(I-1,J)
      PHI2=PHI(I,J-1)
      CALL NUPT C
      X(1,J)= XX
      Z(1,J)= ZZ
      SIGMA(1,J)= SIMA
      PSI(1,J)= PSS1
      P(I)= PP
      Q(I)= QQ

```

- iv -

```

      PHI(I,J)=PHHI
75  CONTINUE
50  GO TO 51
51  IF(CTRO .EQ. 1) GO TO 100
      IF(J.EQ.1) THEN
        PRINT 19
        DO 80 I=2,L
          CALL RESOUT(I,1)
80    CONTINUE
        ELSE
          GO TO 100
        ENDIF
100 CONTINUE
C   END OF RADIAL ZONE COMPUTATION
      IF(CTRO .EQ. 1) GO TO 67
      PRINT, ' '
      PRINT, ' '
      PRINT, ' '
      PRINT 25
25  FORMAT(11X,'RADIAL ZONE COMPUTATION COMPLETED')
      PRINT 26
26  FORMAT(11X,'*****')
      PRINT 67B
67B FORMAT('1',///25X,' COMPUTATION OF THE EXTRA DOMAIN',
1      //25X,' *****')
67  CONTINUE
      H=2
      J=RAD
300 H= H+1
      J= J+1
      H1=H+1
      IF(CTRO .EQ. 1) GO TO 68
      PRINT, ' '
      PRINT, ' '
      PRINT, ' '
      PRINT, ' '
      PRINT 47,J,H
47  FORMAT(10X,'J=',I3,5X,'FIRST SPIRAL=',I3)
      PRINT 59
59  FORMAT(10X,'*****',5X,'*****')
      PRINT 19
68  CONTINUE
      X1= X(H-1,J-1)
      Z1= Z(H-1,J-1)
      SIGMA1= SIGMA(H-1,J-1)
      PSI1= PSI(H-1,J-1)
      X2= X(H,J-1)
      Z2= Z(H,J-1)
      SIGMA2= SIGMA(H,J-1)
      PSI2= PSI(H,J-1)
      PHI1=PHI(H-1,J-1)
      PHI2=PHI(H,J-1)
      CALL ENDPT
      X(H,J)= XX

```

- v -

```

      Z(H,J)= Z2
      SIGMA(H,J)= SIMA
      PSI(H,J)= PSS1
      PHI(H,J)=PHHI
      IF(H1 .GT. L) GO TO 117
      DO 250 I=H1,L
        X1= X(I-1,J)
        Z1= Z(I-1,J)
        SIGMA1= SIGMA(I-1,J)
        PSI1= PSI(I-1,J)
        X2= X(I,J-1)
        Z2= Z(I,J-1)
        SIGMA2= SIGMA(I,J-1)
        PSI2= PSI(I,J-1)
        PHI1=PHI(I-1,J)
        PHI2=PHI(I,J-1)
        CALL NUPT
        X(I,J)= XX
        Z(I,J)= Z2
        SIGMA(I,J)= SIMA
        PSI(I,J)= PSS1
        P(I)= PP
        Q(I)= QQ
        PHI(I,J)=PHHI
250  CONTINUE
117  IF(H .LT. L) GO TO 300
      RAD1=RAD+L-2
      PRINT,RAD1
      IF(ABS(BETA) .GE. 0.1) THEN
        CALL NGAMA(RAD)
      ELSE
        CALL KAPE(RAD)
      ENDIF
      DIF=(SIGMA(L,RAD1)-10*SIGMA(2,RAD))
      DIF=DIF/SIGMA(L,RAD1)
      DIF= ABS(DIF)
      PRINT 200, SIGMA(L,RAD1), SIGMA(2,RAD), DIF, CONVG
200  FORMAT(////////20X,'*****')
1    FORMAT(////////20X,'*****')
      SIGMA(L,RAD+L-2)=',F12.6,'*', //20X,'*
2    //20X,'*      SIGMA(2, RAD)=',F12.6,'
3    //20X,'*      DIF=',F8.5,'
4    //20X,'*      CONVG=',F8.5,'
5    //20X,'*****')
      IF(DIF .GT. CONVG) THEN
        DO 500 I=2,L
          SIGMA(I,1)= (Z(I,1)-Z(2,1))/(1.0-SNPHI)
          SIGMA(I,1)= SIGMA(I,1) + (SIGMA(L,1)/10.0)
500  CONTINUE

```

- vi -

```

      KS=X*SPI
      DO 510 J= 2, KS
        X(2,J)= (X(L,W*J))/10.0
        Z(2,J)= (Z(L,W*J))/10.0
        SIGMA(2,J)= (SIGMA(L,W*J))/10.0
        PSI(2,J)= PSI(L,W*J)
        PHI(2,J)=PHI(L,W*J)
510    CONTINUE
      PRINT 22,C
      PRINT
27, X(2,RAD), Z(2,RAD), SIGMA(2,RAD), PSI(2,RAD), PHI(2,RAD)
27    FORMAT(////10X, 'X(2,RAD)=' ,F5.2,
1      //10X, 'Z(2,RAD)=' ,F5.2,
2      //10X, 'SIGMA(2,RAD)=' ,F12.8,
3      //10X, 'PSI(2,RAD)=' ,F12.8,
4      //10X, 'PHI(2,RAD)=' ,F12.8//)
      PRINT 30,C
30    FORMAT('1',////////14X, 'RESULTS OF FIRST SPIRAL
AFTER',13, 'SCALE REDUCTION')
      PRINT 19
      DO 415 J=1,RAD
        CALL RESOUT(2,J)
415    CONTINUE
      GO TO 450
      ELSE
      GO TO 222
      ENDIF
222    PRINT 333
333    FORMAT('1',////////20X, '*****',
1      //20X, '*'
2      //20X, '*' END OF FINAL COMPUTATION
3      //20X, '*'
4      //20X, '*' ZERO SURCHARGE CASE
5      //20X, '*'
6      //20X, '*' FIRST: REGULAR RADIAL ZONE
7      //20X, '*'
8      //20X, '*****'//)
      DO 555 J=2,RAD
        PRINT 890,J
        PRINT 891
890    FORMAT(////10X, 'J=' ,I3)
891    FORMAT(10X, '*****')
        PRINT 19
        DO 666 I=1,L

```

- vii -

```

      CALL RESOUT(1,J)
666    CONTINUE
555    CONTINUE
      DO 557 I=1,L
        PRINT 893,I
        PRINT 896
893    FORMAT('1',////10X, 'I=' ,I3)
896    FORMAT(10X, '*****')
        PRINT 19
        DO 669 J=2,RAD
          CALL RESOUT(1,J)
669    CONTINUE
557    CONTINUE
      DO 92 I=1,L
        DO 93 J=1,RAD
          CALL OUTPUT(1,J,RAD)
93    CONTINUE
92    CONTINUE
103    PRINT 104
104    FORMAT( 20X, '*****')
        PRINT 756
756    FORMAT('1',////45X, '*****',
1      //45X, '*'
2      //45X, '*' EXTRA DOMAIN CREATED BY
3      //45X, '*'
4      //45X, '*' SCALE REDUCTIONS
5      //45X, '*'
6      //45X, '*****'//)
      KP=0.0
      DO 777 H=3,L
        J=RAD+H-2
        PRINT 890,J
        PRINT 891
        PRINT 19
        DO 888 I=H,L
          CALL RESOUT(1,J)
          CALL OUTPUT(1,J,RAD)
888    CONTINUE
        KP=(1+SIN(PHI(H,J)))/(1-SIN(PHI(H,J)))*KP
777    CONTINUE
      KP=KP/(L-2)
      PRINT 112,KP
C234567
112    FORMAT('1',////////20X, 'KP=' ,F9.4) C
      PRINT 104
      STOP
      END

```

- viii -



```

C *****
C
SUBROUTINE HEADIN
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
INTEGER L, K, C
COMMON L, K, C
REAL XX, ZZ, SIMA, PSS1, PHH1
COMMON /TATA/ X1, X2, XX, Z1, Z2, ZZ, PHH1
COMMON /TITI/ SIGMA1, SIGMA2, SIMA
COMMON /TETE/ PSI1, PSI2, PSS1
COMMON /TOTO/ PI, SNPHI, CSPHI
INTEGER PP, QQ, CONTRO
COMMON /OUI/ PP, QQ, CONTRO
REAL MU
COMMON /AH/ MU, PHILO, SIGLOC, PHI1, PHI2, PHILOP, PHILOT
REAL
X(40,40), Z(40,40), SIGMA(40,40), PSI(40,40), PHI(40,40)
COMMON /LALA/ X, Z, SIGMA, PSI
COMMON /ALO/ PHI
INTEGER P(40), Q(40)
COMMON /LOLO/ P, Q
REAL EO, LANDA
COMMON
EO, LANDA, PMIN, PMAX, PHIMIN, PHIMAX, PHI1, EMIN, EMAX /CRIT/
C
PRINT 10, FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
10 FORMAT('1', //10X, 'FIRPHI=', P5.1,
1 //10X, 'BETA=', P5.1,
2 //10X, 'DELTA=', P5.1,
3 //10X, 'TOL=', P7.5,
4 //10X, 'CONVG=', P7.5,
5 //10X, 'SIZE=', P6.3,
6 //10X, 'GAMMA=', P6.3)
PHIMIN=PHIMIN*180./PI
PHIMAX=PHIMAX*180./PI
PRINT 11, L,
11 FORMAT(//10X, 'NUMBER OF SPIRAL LINES=', I3,
1 //10X, 'NUMBER OF RADIAL LINES=', I3,
2 //10X, 'VOID RATIO EO =', P6.3,
3 //10X, 'SLOPE LANDA=', P6.3,
4 //10X, 'E-MINIMUM=', P6.3,
5 //10X, 'E-MAXIMUM=', P6.3,
6 //10X, 'P-MINIMUM=', P7.1,
7 //10X, 'P-MAXIMUM=', P7.1,
8 //10X, 'PHI-MIN=', P5.2,
9 //10X, 'PHI-MAX=', P5.2)
PHIMIN=PHIMIN*PI/180.
PHIMAX=PHIMAX*PI/180.
RETURN
END
C

```

- ix -

```

C *****
C
SUBROUTINE RESOUT(A,B)
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
INTEGER L, K, C
COMMON L, K, C
REAL XX, ZZ, SIMA, PSS1, PHH1
COMMON /TATA/ X1, X2, XX, Z1, Z2, ZZ, PHH1
COMMON /TITI/ SIGMA1, SIGMA2, SIMA
COMMON /TETE/ PSI1, PSI2, PSS1
COMMON /TOTO/ PI, SNPHI, CSPHI
INTEGER PP, QQ, CONTRO
COMMON /OUI/ PP, QQ, CONTRO
REAL MU
COMMON /AH/ MU, PHILO, SIGLOC, PHI1, PHI2, PHILOP, PHILOT
REAL
X(40,40), Z(40,40), SIGMA(40,40), PSI(40,40), PHI(40,40)
COMMON /LALA/ X, Z, SIGMA, PSI
COMMON /ALO/ PHI
INTEGER P(40), Q(40)
COMMON /LOLO/ P, Q
REAL EO, LANDA
COMMON
EO, LANDA, PMIN, PMAX, PHIMIN, PHIMAX, PHI1, EMIN, EMAX /CRIT/
C
PHI(A,B)=PHI(A,B)*180./PI
PRINT
20, X(A,B), Z(A,B), SIGMA(A,B), PSI(A,B), P(A), Q(A),
1 PHI(A,B)
20
FORMAT(11X, P12.6, 3X, P12.6, E15.5, 3X, P12.6, 8X, I4, 1X, I4, 6X,
1 P6.3)
PHI(A,B)=PHI(A,B)*PI/180.
RETURN
END
C
C
C
C *****
C
SUBROUTINE OUTPUT(A,B,RAD)
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
INTEGER L, K, C
COMMON L, K, C
REAL XX, ZZ, SIMA, PSS1, PHH1
COMMON /TATA/ X1, X2, XX, Z1, Z2, ZZ, PHH1
COMMON /TITI/ SIGMA1, SIGMA2, SIMA
COMMON /TETE/ PSI1, PSI2, PSS1
COMMON /TOTO/ PI, SNPHI, CSPHI
INTEGER PP, QQ, CONTRO
COMMON /OUI/ PP, QQ, CONTRO
REAL MU

```

- x -

```

COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
COMMON /CRIT/
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX
INTEGER A,B,THI(40,40),AN,BN,LHI(40,40),MI,RAD1,RAD
REAL T,R,XP(40,40),XN(40,40),TN,RN,MN,M C
RAD1=PI/2
MI=(B-1)/2
MI=INT(MI)
MC=MN-MI
IF(M C.GT. 0.001) THEN
GO TO 44
ELSE
GO TO 33
ENDIF
33 PHI(A,B)=PHI(A,B)*1800./PI
THI(A,B)=INT(PHI(A,B))
PHI(A,B)=PHI(A,B)*PI/1800.
WRITE(8,20) X(A,B),Z(A,B),A,B,THI(A,B)
20 FORMAT(P8.4,3X,P8.4,8X,I8,8X,I8,3X,I3)
21 FORMAT(P8.4,3X,P8.4,3X,P9.0,3X,P9.0,3X,I3)
T=A*100.
R=B*100.
IF(R.EQ. 500) GO TO 44
IF(R.EQ. 700) GO TO 44
IF(R.EQ. 900) GO TO 44
IF(R.EQ. 1100) GO TO 44
WRITE(8,21) X(A,B),Z(A,B),R,T,THI(A,B)
44 CONTINUE
RETURN
END

```

```

C
C
C
C345678
SUBROUTINE POLCOM
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX, ZZ, SIMA, PSS1,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSS1
COMMON /TOTO/ PI ,SNPHI,CSPHI
INTEGER PP, QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO

```

- xi -

```

REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
COMMON /CRIT/
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX
INTEGER KK, K1, I, J
REAL PSI1,PSIF,INT,FIRSIG,FINSIG,TNPHI,FINPHI
C
PSI1= 3.1415927 / 2.
FINPHI=39./57.2957795
MU=PI/4.-FINPHI/2.
400 PSIF=MU* BETA
405 INT= (PSI1-PSIF)/ K
KK=K-1
PRINT, ' '
PRINT, '*****'
1*****
PRINT, ' '
PRINT, ' POLE COMPT < > FIRST
SPIRAL:(1,J) '
PRINT, ' '
PRINT, ' '
PRINT, 14
PRINT, ' '
DO 350 J=2,K
X(1,J)=X(1,1)
Z(1,J)=Z(1,1)
FIRPHI=PHI(1,J-1)
34 TNPHI=TAN(FIRPHI)
FIRSIG=EXP(ALOG(SIGMA(1,J-1))+INT*2*TNPHI)
SIGLOC=FIRSIG*SIZE*GAMMA*3.2808*6.369
CALL PHILOC
DI=(PHILO-FIRPHI)/PHILO
IF(ABS(DI) .GT. 0.001) THEN
FIRPHI=PHILO
GO TO 34
ELSE
GO TO 37
ENDIF
37 SIGMA(1,J)=FIRSIG
PHI(1,J)=PHILO
PSI(1,J)= PSI(1,J-1) - INT
PRINT, 12,X(1,J),Z(1,J),SIGMA(1,J),PSI(1,J)
12 FORMAT(1X,P8.5,3X,P8.5,215.5,3X,P8.5)
350 CONTINUE
PRINT, ' '
PRINT, '*****'
1*****

```

- xii -

```

IF(CONTRO.EQ. 1) GO TO 64
PRINT, ' '
PRINT, ' '
PRINT, ' '
PRINT, ' '
PRINT, ' '
19 1 FORMAT(////17X,'X',13X,'Z',12X,'SIGMA',12X,'PSI',
15X,'P',4X,'Q',8X,'PHI'//)
64 PRINT, ' '
CONTINUE
PRINT, ' '
DO 355 J=2,K
X1= X(1,J)
Z1= Z(1,J)
SIGMA1= SIGMA(1,J)
PSI1= PSI(1,J)
X2= X(2,J-1)
Z2= Z(2,J-1)
SIGMA2= SIGMA(2,J-1)
PSI2= PSI(2,J-1)
PHI1=PHI(1,J)
PHI2=PHI(2,J-1)
CALL NUPT
X(2,J)= XX
Z(2,J)= ZZ
SIGMA(2,J)= SIMA
PSI(2,J)= PSSI
P(J)=PP
Q(J)= QQ
PHI(2,J)=PHHI
355 CONTINUE
C CALCULATION OF A WALL POINT
X1= X(1,K)
Z1= Z(1,K)
SIGMA1= SIGMA(1,K)
PSI1= PSI(1,K)
X2= X(2,K)
Z2= Z(2,K)
SIGMA2= SIGMA(2,K)
PSI2= PSI(2,K)
PHI1=PHI(1,K)
PHI2=PHI(2,K)
C
CALL ENDPT
K1=K+1
X(2,K1)=XX
Z(2,K1)=ZZ
SIGMA(2,K1)=SIMA
PSI(2,K1)=PSSI
P(K1)=PP
Q(K1)=QQ
PHI(2,K1)=PHHI
PRINT, ' '
PRINT, '*****'
1*****

```

- xiii -

```

PRINT, ' '
PRINT, ' '
PRINT, ' '
PRINT, ' '
PRINT, ' '
LAST POINT:
1 (2,K+1)
PRINT, ' '
PRINT, ' '
PRINT, ' '
C2345678
11
FORMAT(14X,'X',10X,'Z',10X,'SIGMA',9X,'PSI',15X,'P',4X,'Q',
18X,'PHI')
14 FORMAT(14X,'X',10X,'Z',10X,'SIGMA',10X,'PSI')
PRINT, ' '
PRINT, ' '
DO 65 J=1,K1
PHI(2,J)=PHI(2,J)*180./PI
PRINT
16,X(2,J),Z(2,J),SIGMA(2,J),PSI(2,J),P(J),Q(J),
1PHI(2,J)
16
FORMAT(11X,P8.5,3X,P8.5,E15.5,3X,P8.5,9X,I4,1X,I4,7X,P6.3)
PHI(2,J)=PHI(2,J)*PI/180.
65 CONTINUE
RETURN
END

```

```

C
C
C*****
C
C
SUBROUTINE NUPT
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX, ZZ, SIMA, PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ PI ,SNPHI,CSPHI
INTEGER PP, QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
C2345678

```

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```

COMMON
EO, LANDA, PMIN, PMAX, PHIMIN, PHIMAX, PHI1, EMIN, EMAX
REAL ZIGMA1, ZIGMA2, PZ11, PZ12, SUM1, SUM2
REAL DIFF1, DIFF2, CSDIF1, CSSUM2
REAL A, D, ETA, X1, F, G, U, V
INTEGER I, J
REAL TNPH1, PHHH1

C
C
C
SIMA = (SIGMA1+SIGMA2)/2.
PSSI = (PSI1+PSI2)/2.
PP=1
QQ=1
C2345678
PHI1 = (PHI1+PHI2)/2.0
PHILOP = PHI1
PHI1 = PHILOP
CALL RPLANE
CALL CRITIC
MU = PI/4. * PHILOP/2.
ZIGMA1 = (SIMA+SIGMA1)/2.
ZIGMA2 = (SIMA+SIGMA2)/2.
PZ11 = (PSSI+PSI1)/2.
PZ12 = (PSSI+PSI2)/2.
SUM1 = SIN(PZ11+MU)
SUM2 = SIN(PZ12+MU)
DIFF1 = SIN(PZ11-MU)
DIFF2 = SIN(PZ12-MU)
CSDIF1 = COS(PZ11-MU)
CSSUM2 = COS(PZ12-MU)
A = SUM1/(2*ZIGMA1*SNPH1*CSDIF1)
B = -1. * (DIFF2)/(2.*ZIGMA2*SNPH1*CSSUM2)
ZZ = Z1*DIFF1/CSDIF1 - Z2*SUM2/CSSUM2 - X1*X2
ZZ = ZZ/(DIFF1/CSDIF1 - SUM2/CSSUM2)
TNPH1 = TAN(PHILOP)
ETA = ALOG(SIGMA1) / (2.*TNPH1) - PSI1 + A*(ZZ - Z1)
XI = ALOG(SIGMA2) / (2.*TNPH1) + PSI2 + D*(ZZ - Z2)
P = EXP((XI-ETA)*TNPH1)
G = (X1-ETA)/2.
U = ABS((P-SIMA)/P)
V = ABS((G-PSSI)/G)

C
IF(V.GT.U) GO TO 202
IF(U.LE.TOL) GO TO 203
SIMA = F
PP = PP+1
IF(PP-100) 201,201,203
C2345678
202 IF(V.LE.TOL) GO TO 203
PSSI = G
QQ = QQ+1
IF(QQ-100) 201,201,203

```

- xv -

```

203 XX = X1 + (ZZ - Z1)*DIFF1/CSDIF1
XX = (XX+X2 + (ZZ-Z2)*SUM2/CSSUM2)/2.
PHHI = PHILOP
PHHH1 = PHHI*180./PI

C
IF(CONTRO.EQ. 1) GO TO 16
PRINT 13, XX, ZZ, SIMA, PSSI, PP, QQ, PHHH1
13
FORMAT(11X, F12.8, 3X, F12.8, E15.5, 3X, F12.8, 8X, 14, 1X, 14,
16X, F6.3)
16 CONTINUE
RETURN
END

C
C*****
C
SUBROUTINE ENDPT
REAL FIRPHI
COMMON FIRPHI, BETA, DELTA, TOL, CONVG, SIZE, GAMMA
INTEGER L, K, C
COMMON L, K, C
REAL XX, ZZ, SIMA, PSSI, PHHI
COMMON /TATA/ X1, X2, XX, Z1, Z2, ZZ, PHHI
COMMON /TIT1/ SIGMA1, SIGMA2, SIMA
COMMON /TETE/ PSI1, PSI2, PSSI
COMMON /TOTO/ P1, SNPH1, CSPHI
INTEGER PP, QQ, CONTRO
COMMON /OUI/ PP, QQ, CONTRO
REAL MU
COMMON /AH/ MU, PHILO, SIGLOC, PHI1, PHI2, PHILOP, PHILOT
REAL
X(40,40), Z(40,40), SIGMA(40,40), PSI(40,40), PHI(40,40)
COMMON /LALA/ X, Z, SIGMA, PSI
COMMON /ALO/ PHI
INTEGER P(40), Q(40)
COMMON /LOLO/ P, Q
REAL EO, LANDA
C2345678
COMMON
EO, LANDA, PMIN, PMAX, PHIMIN, PHIMAX, PHI1, EMIN, EMAX
REAL ENDPS1
REAL ZIGMA2, PZ12, SUM2, DIFF2, CSSUM2
REAL B, X1, F, U, PHHH1
INTEGER I, J

```

/CRIT/

```

C
C
PP=1
QQ=1
SIMA = (SIGMA1+SIGMA2)/2.
PHI1 = (PHI1+PHI2)/2.0
PHILOP = PHI1
PHI1 = PHILOP
CALL RPLANE
CALL CRITIC
302

```

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```

      MU=PI/4.-PHILOP/2.
      DELTA=PHILOP
      IF(BETA .LE. 0.001) THEN
        GO TO 301
      ELSE
        GO TO 309
      ENDIF
C2345678
301  ENDPST=MU+BETA
      GO TO 305
309  ENDPST=MU*(1.732-DELTA)
305  PSS1=ENDPST
      ZIGMA2=(SIMA+SIGMA2)/2.
      PZ12=(PSS1+PSI2)/2.
      SUM2=SIN(PZ12*MU)
      DIPF2=SIN(PZ12-MU)
      CSSUM2= COS(PZ12*MU)
      B=-1.*(DIPF2)/(2.*ZIGMA2*SNPHI*CSSUM2)
      IF(ABS(BETA) .LE. 0.0001) THEN
        BETA=0.0001
      ELSE
        BETA=-1.*(1.732-DELTA)
      ENDIF
      ZZ=Z1*TAN(BETA)-X1-Z2*SUM2/CSSUM2+X2
      ZZ=ZZ/(TAN(BETA)-(SUM2/CSSUM2))
      XI=ALOG(SIGMA2)/(2.*TAN(PHILOP))+PSI2+B*(ZZ-Z2)
      F=EXP(2.*TAN(PHILOP)*(X1-PSS1))
      U=ABS((F-SIMA)/F)
      IF(U .LE. TOL) GO TO 303
      SIMA=F
      PP=PP+1
      IF(PP-100) 302,302,303
C2345678
303  XX=X1+(ZZ-Z1)*TAN(BETA)
      PHH1=PHILOP
      PHHH1=PHH1*180./PI
      IF (CONTRO .EQ. 1) GO TO 26
      PRINT 13, XX,ZZ,SIMA,PSS1,PP,QQ,PHHH1
13  FORMAT(11X,F12.8,3X,F12.8,E15.5,3X,F12.8,8X,14,1X,14,
16X,F6.3)
      PRINT 306
306  FORMAT(' ',
1  'BOUNDARY')
26  CONTINUE
      RETURN
      END
C
C
C
C
SUBROUTINE EDGEPA
REAL FIRPHI

```

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```

COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX, ZZ, SIMA, PSS1,PHH1
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHH1
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSS1
COMMON /TOTO/ P1,SNPH1,CSPH1
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHILOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
C2345678
COMMON
EO,LANDA,PHIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX /CRIT/
INTEGER CC,1,J
REAL DI,NUZ,NSIMA
C
C
IF (C.EQ. 0) THEN
  CC=2
  C=1
  ELSE
  CC=3
  ENDIF
DO 23 I=CC,L
  FIRPHI=PHI(I-1,1)
  MU=PI/4.0-FIRPHI/2.0
  PHILOP=FIRPHI
  CALL RPLANE
  PHI1=PHILOT
  MU=PI/4.-PHI1/2.
  Z(I,1)=X(I,1)*TAN(MU)
  SIMA=(1.0+Z(I,1))/(1.0-SIN(PHI1))
  CALL CRITIC
  MU=PI/4.0-PHILOP/2.0
  NUZ=X(I,1)*TAN(MU)
  NSIMA=(1.0+NUZ)/(1.0-SIN(PHILOP))
  SIMA=NSIMA
C234567
  Z(I,1)=NUZ
  PHI(I,1)=PHILOP
  SIGMA(I,1)=NSIMA
  PSI(I,1)=3.1415927/2.0
23  CONTINUE
      RETURN
      END

```

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```

C*****
C
C
SUBROUTINE PHILOC
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ P1,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
COMMON
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX /CRIT/
REAL TT
INTEGER DD
C
C*** SIGLO CIS DIMENSIONAL ****
C
DD=2
IF(DD.EQ. 1) THEN
GO TO 10
ELSE
GO TO 11
ENDIF
11 IF(ABS(SIGLOC).LT. 1000.0) THEN
C234567
PHILO=(47.8-4.7*(ALOG(SIGLOC)-2.30258)/4.60517)/57.295779
ELSE
IF(ABS(SIGLOC).LT. 200000.0) THEN
PHILO=(43.2-11.2*(ALOG(SIGLOC)-6.90776)/5.29832)/57.295779
ELSE
PHILO=32.0/57.2957795
ENDIF
ENDIF
GO TO 30
C234567
10 IF(ABS(SIGLOC).LT. 100.0) THEN
PHILO=(47.8-4.7*(ALOG(SIGLOC)-2.30258)/4.60517)/57.295779
ELSE

```

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```

**0.5) TT=(47.8-4.7*(ALOG(100.0)-2.30258)/4.60517)/(100.0
PHILO=(TT*(SIGLOC**0.5))/57.295779
ENDIF
30 SNPHI=SIN(PHILO)
CSPHI=COS(PHILO)
RETURN
END

```

```

C
C
C*****
C2345678
SUBROUTINE CRITIC
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ P1,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
COMMON
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX /CRIT/
REAL V,VLANDA,KK,PHICI,QOP,DI,FPR,M
REAL MMAX,MMIN,VLAMAX,VLAMIN,AA,BB
C
V=1.0+EO
PHI1=PHILOT
5 KK=(1.0+SIN(PHI1))/(1.0-SIN(PHI1))
SIMA=SIMA*GAMMA*SIZE
PFR=(2*(1.0+2*KK)*SIMA)/(3*(1+KK))
SIMA=SIMA/(SIZE*GAMMA)
VLANDA=V + LANDA*ALOG(PFR)
VLAMAX=(1.+EMAX) + LANDA*ALOG(PMAX)
VLAMIN=(1.+EMIN) + LANDA*ALOG(PMIN)
IF(VLANDA.GT. VLAMAX) THEN
VLANDA=VLAMAX
PHICI=PHIMIN
GO TO 2
ELSE
GO TO 7
ENDIF

```

- xx -

```

7 IF(VLANDA .LT. VLAMIN) THEN
  VLANDA=VLAMIN
  PHICI=PHIMAX
  GO TO 2
ELSE
  GO TO 3
ENDIF
3 MMAX=(6.*SIN(PHIMIN))/(3-SIN(PHIMIN))
MMIN=(6.*SIN(PHIMAX))/(3-SIN(PHIMAX))
AA=(MMAX-MMIN)/(VLAMAX-VLAMIN)
BB=MMAX AA*VLAMAX
QOP=AA*VLANDA + BB
M=QOP
TT=(3*M)/(6+M)
PHICI=ARSIN(TT)
2 PHILOT=PHICI
DI=(PHILOT-PHII)/PHILOT
IF(ABS(DI).GT.0.11) THEN
  PHII=PHILOT
  GO TO 5
ELSE
  GO TO 9
ENDIF
9 CALL PPLANE
PHII=PHILOT
SNPHI=SIN(PHILOT)
CSPHI=COS(PHILOT)
RETURN
END
*****
SUBROUTINE PPLANE
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ PI,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
C234567

```

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```

COMMON /CRIT/
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHII,EMIN,EMAX
REAL PHIPL1,PHIPL2
C
C
C
PHIPL1=32.8328*PI/180.
PHIPL2=36.*PI/180.
IF(PHILOT .LT. PHIPL1) GO TO 10
IF(PHILOT .GT. PHIPL2) GO TO 20
PHILOT=PHILOT*180./PI
PHIOP=1.6667*ALOG(PHILOT) - 2.3362
PHIOP=EXP(PHIOP)
PHIOP=PHIOP*PI/180.
GO TO 35 C
C234567
10 PHIOP=PHILOT
GO TO 35
20 PHILOT=PHILOT*180./PI
PHIOP=1.2944*ALOG(PHILOT)-1.002
PHIOP=EXP(PHIOP)
PHIOP=PHIOP*PI/180.
GO TO 35
35 RETURN
END
C
C*****
SUBROUTINE RPLANE
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TITI/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ PI,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
C234567
COMMON /CRIT/
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHII,EMIN,EMAX
REAL PHIPL1,PHIPL2
C
C

```

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```

PHIPL1=32.8328*PI/180.
PHIPL2=37.9569*PI/180.
IF(PHILOP.LT. PHIPL1) GO TO 10
IF(PHILOP.GT. PHIPL2) GO TO 20
PHILOP=PHILOP*180./PI
PHILOT=(ALOG(PHILOP)/2.3362)/1.6667
PHILOT=EXP(PHILOT)
PHILOT=PHILOT*PI/180.
GO TO 35
C234567
10 PHILOT=PHILOP
GO TO 35
20 PHILOP=PHILOP*180./PI
PHILOT=(ALOG(PHILOP)+1.002)/1.2944
PHILOT=EXP(PHILOT)
PHILOT=PHILOT*PI/180.
GO TO 35
35 RETURN
END
C
C-----
C
SUBROUTINE NGAMA(RAD)
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TIT1/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ PI,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHILOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA
C2345678
COMMON /CRIT/
EO,LANDA,PHIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX
REAL COSPSI,SIMAV(30),NGAMMA,F,E,SOMME,RATIO,CENTRE
REAL*8 B
INTEGER L1,RAD,I,J,H,RAD1,RAD3
C
RAD1=RAD+L-2
RAD3=RAD+L-3
I=0
H=0
- xxiii

```

```

J=RAD-2
4 H=H+1
J=J+1
C2345678
I=I+1
SNPHI=SIN(PHI(H,J))
PSI(H,J)=2.*PSI(H,J)
COSPSI=COS(PHI(H,J))
PSI(H,J)=PSI(H,J)/2
SIMAV(I)=SIGMA(H,J)*(1+SNPHI*COSPSI)
IF(H.LT. L) GO TO 4
I=2
H=2
J=RAD
NGAMMA=0.0
SOMME=0.0
6 I=I+1
H=H+1
J=J+1
SOMME=(SIMAV(I-1)+SIMAV(I))/2.
SOMME=SOMME*(X(H,J)-X(H-1,J-1))
NGAMMA=NGAMMA + SOMME
L1=L
IF(I.LT. L1) GO TO 6
C234567
NGAMMA=-2.*(NGAMMA)
PRINT 10,NGAMMA
B=X(L,RAD1)*2.
RATIO= 1./B*2
NGAMMA=NGAMMA*RATIO*2.
PRINT 10,NGAMMA
10 FORMAT('1.//////20X,'NGAMMA=',F9.4)
B=B*SIZE/1.732
PRINT 12,B
12 FORMAT('//////20X,'BREADTH=',F9.4)
B=B*1.732/SIZE
B=(1./B)*(10.**((1-C)))
PRINT 11,B
11 FORMAT('//////20X,'D/B=',F9.4)
RETURN
END
C-----
C
SUBROUTINE KAPE(RAD)
REAL FIRPHI
COMMON FIRPHI,BETA,DELTA,TOL,CONVG,SIZE,GAMMA
INTEGER L,K,C
COMMON L,K,C
REAL XX,ZZ,SIMA,PSSI,PHHI
COMMON /TATA/ X1,X2,XX,Z1,Z2,ZZ,PHHI
COMMON /TIT1/ SIGMA1,SIGMA2,SIMA
COMMON /TETE/ PSI1,PSI2,PSSI
COMMON /TOTO/ PI,SNPHI,CSPHI
INTEGER PP,QQ,CONTRO
- xxiv -

```



```

COMMON /OUI/ PP,QQ,CONTRO
REAL MU
COMMON /AH/ MU,PHILO,SIGLOC,PHI1,PHI2,PHIOP,PHILOT
REAL
X(40,40),Z(40,40),SIGMA(40,40),PSI(40,40),PHI(40,40)
COMMON /LALA/ X,Z,SIGMA,PSI
COMMON /ALO/ PHI
INTEGER P(40),Q(40)
COMMON /LOLO/ P,Q
REAL EO,LANDA

C234567
COMMON /CRIT/
EO,LANDA,PMIN,PMAX,PHIMIN,PHIMAX,PHI1,EMIN,EMAX
REAL COSPSI,SIMAH(30),KP,F,E,SOMME,D,RATIO,CENTRE
INTEGER L1,RAD,I,J,H,RAD1,RAD3

C
RAD1=RAD*L-2
I=0
H=0
J=RAD-2
RAD3=RAD*L-3
4 H=H+1
J=J+1
C2345678
I=I+1
SINPHI=SIN(PHI(H,J))
PSI(H,J)=2.*PSI(H,J)
COSPSI=CCS(PSI(H,J))
PSI(H,J)=PSI(H,J)/2.
SIMAH(I)=SIGMA(H,J)*(1-SINPHI*COSPSI)
IF(H.LT.L) GO TO 4

C
I=1
H=1
J=RAD-1
KP=0.0
SOMME=0.0
6 I=I+1
H=H+1
J=J+1
SOMME=(SIMAH(I-1)+SIMAH(I))/2.
SOMME=SOMME*(Z(H,J)-Z(H-1,J-1))
KP=KP+SOMME
L1=L
IF(I.LT.L1) GO TO 6

C234567
KP=KP
PRINT 10,KP
D=Z(L,RAD1)
D=D*SIZE/1.732
PRINT 14,D
14 FORMAT(/////////20X,'HEIGHT=',F9.4)
D=D*1.732/SIZE
RATIO=1./D**2

```

- xxv -

```

C
KP=KP*RATIO*2.

10 PRINT 10,KP
FORMAT(/////////20X,'KP=',F9.4)
RETURN
END

C*****
|C
C L'ORGANISATION DE L'ENTREE DES DONNEES
C EST LA SUIVANTE:
|C
C FIRPHI,DELTA,BETA,TOL,CONVG,SIZE,GS,GW
C CONTROL
C L,K
C EMIN,EMAX,LANDA,PMIN,PMAX,PHIMIN,PHIMAX
C X(1,1)
|C
C*****
|C
ENTRY
30.0 30.0 -30.0 0.0001 0.001 0.2 2.65 9.81
2
10 20
0.40 0.80 0.175 10 10000 28 45
0.1732
0.215
0.260
0.346
0.519
0.693
0.866
1.216
1.732
SSTOP
//FT08F001 DD DSN=HOVAN.SAS.OUTPUT2,DISP=OLD
/

```

The following pages (115-117) show a sample of the computer print-out. Each point is characterized by :

- X, Z, Physical coordinates
- SIGMA, Stress
- PSI, Direction of the major principal stress
- P,Q, Iteration counts
- PHI, Angle of shearing resistance

The radial number is J. For each radial, the print-out shows the characteristics (X,Z,SIGMA,PSI,P,Q,PHI) of the intersections of the spirals with the radial.

# COMPUTATION OF THE EXTRA DOMAIN

\*\*\*\*\*

J: 30 FIRST SPIRAL: 3  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI
0.00001288	0.12878980	0.25534E 01	0.45423710	5	1	37.980
0.00047808	0.15258320	0.27553E 01	0.49758420	4	3	37.895
0.00401488	0.20177190	0.31138E 01	0.55628030	4	3	37.274
0.02040798	0.30987940	0.37430E 01	0.66647570	4	3	36.845
0.04783745	0.42581180	0.43071E 01	0.73838130	4	3	36.157
0.08400088	0.54478350	0.48148E 01	0.79455700	4	3	35.791
0.17771180	0.79511840	0.57533E 01	0.87837630	4	3	35.181
0.35387220	1.17343100	0.69785E 01	0.98158120	4	3	34.548

BOUNDARY

J: 31 FIRST SPIRAL: 4  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI
0.00001522	0.15224380	0.28283E 01	0.45835840	4	1	37.487
0.00150884	0.20019380	0.33010E 01	0.52765750	3	3	37.074
0.01333430	0.30821200	0.38574E 01	0.62904450	3	3	36.455
0.03593770	0.42084380	0.45443E 01	0.70306010	3	3	35.986
0.06696808	0.53813190	0.50711E 01	0.75923530	3	3	35.615
0.15031390	0.78871870	0.60430E 01	0.84482040	3	3	35.027
0.31143700	1.16745100	0.73055E 01	0.93018770	4	3	34.393

BOUNDARY

J: 32 FIRST SPIRAL: 5  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI
0.00001992	0.19915340	0.36282E 01	0.46477110	4	1	36.753
0.00478881	0.30150820	0.43285E 01	0.56783390	3	3	36.150
0.01950893	0.41347970	0.49549E 01	0.64346880	3	3	35.694
0.04282328	0.52888910	0.55147E 01	0.70112080	3	2	35.333
0.10928480	0.77728600	0.65437E 01	0.78934620	3	3	34.781
0.24575180	1.15554900	0.78781E 01	0.87802170	4	3	34.144

BOUNDARY

J: 33 FIRST SPIRAL: 8  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI	
0.00002884	0.28842800	0.48843E 01	0.47400880	4	1	35.894	BOUNDARY
0.00378835	0.40481920	0.56440E 01	0.55164050	3	2	35.256	
0.01510308	0.51711880	0.62583E 01	0.61124480	3	3	34.810	
0.08889202	0.75884880	0.73820E 01	0.70304010	4	2	34.380	
0.15617370	1.13334500	0.88287E 01	0.78819880	4	3	33.787	

J: 34 FIRST SPIRAL: 7  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI	
0.00004025	0.40250870	0.82318E 01	0.48072880	4	1	34.924	BOUNDARY
0.00301403	0.51034300	0.88813E 01	0.54152290	3	2	34.588	
0.02870457	0.74556500	0.80928E 01	0.53586250	3	3	34.055	
0.08785414	1.11443600	0.96322E 01	0.73188870	3	3	33.481	

J: 35 FIRST SPIRAL: 8  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI	
0.00005084	0.50843820	0.74380E 01	0.48588450	4	1	34.335	BOUNDARY
0.00888252	0.73703140	0.87031E 01	0.58159830	3	3	33.815	
0.05843602	1.08954900	0.10321E 02	0.67897550	3	3	33.254	

J: 36 FIRST SPIRAL: 9  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI	
0.00007311	0.73108820	0.97745E 01	0.49374380	4	1	33.433	BOUNDARY
0.01613783	1.07843400	0.11525E 02	0.59513710	3	3	32.893	

J: 37 FIRST SPIRAL: 10  
\*\*\*\*\*

X	Z	SIGMA	PSI	P	Q	PHI	
0.00010702	1.07016200	0.13041E 02	0.48932180	4	1	32.793	BOUNDARY

HEIGHT: 4.9430

KP: 9.4443

```
*****  
* SIGMA(L,RAD+L-2): 13.040830*  
* SIGMA(2, RAD): 1.305003 *  
* DIF: 0.00070 *  
* CONV: 0.00100 *  
*****
```

## Appendix C

### SAMPLE OF A PLOTTING COMPUTER PROGRAM

1 SAS LOG OS SAS 82.4 VS2/MVS JOB PRINT STEP 00 PROC 18:08 SATURDAY, NOVEMBER 17, 1984

NOTE: THE JOB PRINT HAS BEEN RUN UNDER RELEASE 82.4 OF SAS AT THE UNIVERSITY OF MANITOBA (02248001).

NOTE: CPUID VERSION = 01 SERIAL = 000103 MODEL = 0980

NOTE: SAS OPTIONS SPECIFIED ARE:  
SORT=4

1 OPTIONS DEVICE=VERSATIC;  
2 DATA OYO;  
3 INFILE SASDATA;  
4 INPUT X 2 A B PHI;  
5

NOTE: INFILE SASDATA IS:  
DSNAME=NOVAN.SAS.DUPUTW,  
UNIT=DISK,VOL=SER=USER04,DISP=SHR,  
DCB=(BLKSIZE=8080,LRECL=80,RECFM=FB)

NOTE: 292 LINES WERE READ FROM INFILE SASDATA.  
NOTE: DATA SET WORK.OYO HAS 292 OBSERVATIONS AND 5 VARIABLES. 433 OBS/TRK.  
NOTE: THE DATA STATEMENT USED 0.08 SECONDS AND 388K.

5 PROC PRINT;  
NOTE: THE PROCEDURE PRINT USED 0.18 SECONDS AND 380K AND PRINTED PAGES 1 TO 6.  
6

7 PROC Gplot DATA=OYO UNIFORM;  
8 LABEL Z=HEIGHT OF THE PASSIVE WALL  
9 X=HORIZONTAL;  
10 TITLE .H=2 .P=DUPLX SLIP LINES BEHIND A PASSIVE WALL;  
11 FOOTNOTES .H=1 .P=DUPLX POUR RIE;  
12 PLOT Z=X=A / NOLEGEND VREVERSE VAXIS=0 TO 1.8 BY 0.2  
13 HAXIS=0 TO 1.8 BY 0.2 VREF=1.8 HREF=1.8;  
14 SYMBOL1 V=NONE L=1 I=JOIN;  
15 SYMBOL2 V=NONE L=1 I=JOIN;  
16 SYMBOL3 V=NONE L=1 I=JOIN;  
17 SYMBOL4 V=NONE L=1 I=JOIN;  
18 SYMBOL5 V=NONE L=1 I=JOIN;  
19 SYMBOL6 V=NONE L=1 I=JOIN;  
20 SYMBOL7 V=NONE L=1 I=JOIN;  
21 SYMBOL8 V=NONE L=1 I=JOIN;  
22 SYMBOL9 V=NONE L=1 I=JOIN;  
23 SYMBOL10 V=NONE L=1 I=JOIN;  
24 SYMBOL11 V=NONE L=1 I=JOIN;  
25 SYMBOL12 V=NONE L=1 I=JOIN;  
26 SYMBOL13 V=NONE L=1 I=JOIN;  
27 SYMBOL14 V=NONE L=1 I=JOIN;  
28 SYMBOL15 V=NONE L=1 I=JOIN;  
29 SYMBOL16 V=NONE L=1 I=JOIN;  
30 SYMBOL17 V=NONE L=1 I=JOIN;  
31 SYMBOL18 V=NONE L=1 I=JOIN;  
32 SYMBOL19 V=NONE L=1 I=JOIN;  
33 SYMBOL20 V=NONE L=1 I=JOIN;  
34 SYMBOL21 V=NONE L=1 I=JOIN;  
35 SYMBOL22 V=NONE L=1 I=JOIN;  
36 SYMBOL23 V=NONE L=1 I=JOIN;  
37 SYMBOL24 V=NONE L=1 I=JOIN;  
38 SYMBOL25 V=NONE L=1 I=JOIN;  
39 SYMBOL26 V=NONE L=1 I=JOIN;  
40 SYMBOL27 V=NONE L=1 I=JOIN;  
41 SYMBOL28 V=NONE L=1 I=JOIN;  
42 SYMBOL29 V=NONE L=1 I=JOIN;  
43 SYMBOL30 V=NONE L=1 I=JOIN;  
44  
45  
46

NOTE: THE PROCEDURE Gplot USED 0.80 SECONDS AND 882K.  
NOTE: SAS USED 882K MEMORY.

NOTE: SAS INSTITUTE INC.  
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